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High-power and High Efficiency Yb:YAG Ceramic Laser at Room Temperature

Shinki Nakamura
Ibaraki University
Japan

1. Introduction

Ceramic laser media fabricated by vacuum sintering (Ikesue et al., 1995), (Ikesue et al., 1996), and nanocrystalline (Lu et al., 2000) technology are very attractive materials because they have several remarkable advantages compared with single crystal laser materials. Ceramic samples with a large size can be easily fabricated, whereas this is extremely difficult for single crystals; multiplayer and multifunctional ceramic laser materials are possible because of the polycrystallinity of ceramics (Taira, 2007). Potentially, because of their short fabrication period and because they can be mass-produced, the cost of ceramic laser materials could be much lower than that of single crystals. Furthermore, no complex facilities and critical techniques are required for the growth of ceramics. Since 1995, Ikesue and coworkers have been developing several types of ceramic laser material (Ikesue et al., 1995), (Ikesue et al., 1996), and they found in 2000 that the output power from a 3.4 at.% Nd:YAG ceramic microchip laser is twice that from a Nd:YAG crystal microchip laser of the same size (Shoji et al., 2000). At a low doping concentration, it was found that the efficiency of a diode-end-pumped Nd:YAG ceramic laser is even higher than that of a Nd:YAG single crystal laser. Since 1998, Yanagitani and coworkers have been developing several types of ceramic lasers, and Lu et al. reported the Nd:YAG ceramic laser as one of them in 2001 (Lu et al., 2001). The mechanical properties of YAG ceramics were reported by Kaminskii et al. in ref. (Kaminskii et al., 2003). YAG ceramics had a 10% higher hardness than a YAG single crystal, and the fracture toughness of the YAG ceramics was more than threefold that of the YAG single crystal. Therefore, the ceramics had a higher resistance to thermal shock than the single crystal. Ytterbium (Yb^{3+})-doped materials are very attractive for diode-pumped solid-state lasers (DPSSLs) (Krupke, 2000). The Yb^{3+} -doped materials have high quantum efficiency and exhibit no concentration quenching simply because the Yb^{3+} ion has only two manifolds, namely, the ground state $^2\text{F}_{7/2}$ and the upper level $^2\text{F}_{5/2}$. Thus far, many articles about Yb:YAG crystal lasers have been published (Lacovara et al., 1991), (Hönninger et al., 1995), (Sumida et al., 1998). Yb:YAG has broad absorption and emission bands. The broad absorption band in the near-IR region is suitable for laser-diode (LD) pumping, and the broad emission band enables the generation of ultrashort pulses (Hönninger et al., 1995). However, an Yb:YAG laser is known as a quasi-three-level laser or a quasi-four-level laser, and a finite population exists at the Stark level of the lower manifold $^2\text{F}_{7/2}$, where laser transition terminates, which requires high-intensity pumping, a high-brightness pump

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source, and an efficient heat removal technique (Matsubara et al., 2007), (Sato & Taira, 2004), (Shoji et al., 2004) to prevent reabsorption from the lower level of the laser. Takaichi et al. reported the absorption and emission spectra of a Yb:YAG ceramic ($C_{Yb} = 1$ at.%) and demonstrated laser oscillation, which was the first diode-end-pumped Yb:YAG ceramic laser (not Nd:YAG) with a 345 mW cw output power and a slope efficiency of 26% (Takaichi et al., 2003). Recently, Tsunekane and Taira have demonstrated a high-power diode-edge-pumped single-crystal Yb:YAG / ceramic undoped YAG composite microchip laser (Tsunekane&Taira, 2005), (Tsunekane&Taira, 2006). Early in 2007, a diode-edge-pumped, composite all-ceramic Yb:YAG ($C_{Yb} = 10$ at.%) microchip laser was demonstrated by Tsunekane and Taira, and a 414 W cw output power was obtained with a slope efficiency of 47% (Tsunekane&Taira, 2007). Very recently, Dong et al. have demonstrated a 2.7 W heavily doped (20 at.%) Yb:YAG ceramic laser with a slope efficiency of 52% (Dong et al., 2007); however, its two-pass-pumping miniature laser configuration was more complex than a simple conventional end-pumping configuration and its output power was not markedly high. Nakamura et al. demonstrated a 5.5 W cw Yb:YAG (9.8 at.%) ceramic laser with a slope efficiency of 52% using a simple end-pumping scheme (Nakamura et al., 2008a) with a 400 μm fiber-coupled LD. Dong et al. demonstrated a highly efficient (a slope efficiency of 79%) Yb:YAG ceramic laser (Dong et al., 2006) with a 100 μm fiber-coupled LD using an end-pumping scheme, but its output power was 1.7 W.

In section 2, we report a high-power (6.8 W) and high-efficiency tunable Yb:YAG ceramic laser demonstrated using an end-pumping scheme with a slope efficiency of 72% (Nakamura et al., 2008b) at room temperature (20°C).

Additionally, the tunability of the Yb:YAG ceramic laser for spectroscopy and femtosecond laser application. The broadest tunability is demonstrated in section 3. Finally, A diode-pumped passively mode-locked Yb:YAG ceramic laser is demonstrated in section 4. 417 and 286 fs pulses with average powers of 250 and 25 mW were obtained at 1030 nm using 1 and 0.1% output couplers, respectively. 233 fs pulses with an average power of 20 mW were also obtained at a center wavelength of 1048.3 nm using a 0.1% output coupler. To the best of our knowledge, this is the first demonstration of a diode-pumped mode-locked Yb:YAG ceramic laser.

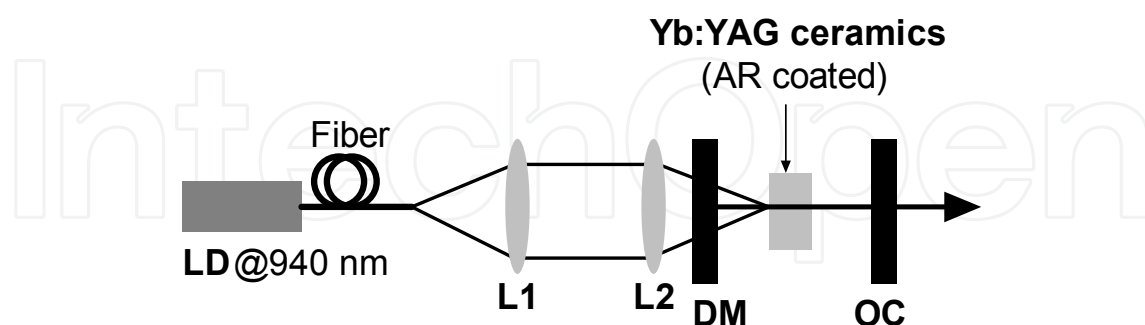


Fig. 1. Experimental setup for the high-power Yb:YAG ceramic laser (Nakamura et al., 2008a)

2. High-power and high efficiency Yb:YAG

2.1 Experimental setup of linear cavity

The experimental setup for the Yb:YAG ceramic laser is shown in Fig. 1. A 940 nm fiber-coupled LD (JENOPTIK Laserdiode, JOLD-30-FC-12) was used as a pumping source, the

core diameter of the fiber was 200 μm , and the numerical aperture (NA) of the fiber was 0.22. The pumping beam was focused onto the ceramic with a ratio of 1:1 using the lenses L1 ($f = 25\text{ mm}$) and L2 ($f = 25\text{ mm}$). The diameter of the focused spot on the ceramic was $\sim 200\text{ }\mu\text{m}$. To obtain high efficiency and high power, a laser cavity consisting of a flat dichroic mirror (DM) and a flat output coupler (OC) as a linear resonator was used. The DM was antireflection (AR)-coated at 940 nm and had a high reflectivity at 1030 nm. The OC was partially-reflection-coated with a transmittance of $T = 1, 5$, and 10% at 1030 nm. An AR-coated ceramic Yb:YAG ($C_{\text{Yb}} = 9.8\text{ at.}\%$, Konoshima Chemical) with dimensions of $5 \times 10 \times 1\text{ mm}^3$ was used. A 1-mm-thick Yb:YAG ceramic plate was wrapped with indium foil and mounted in a water-cooled copper block that acted as a heat sink. Water was maintained at a room temperature of 20°C during laser oscillation. The cavity length was 20 mm, which was optimized, as shown in the later part of this section.

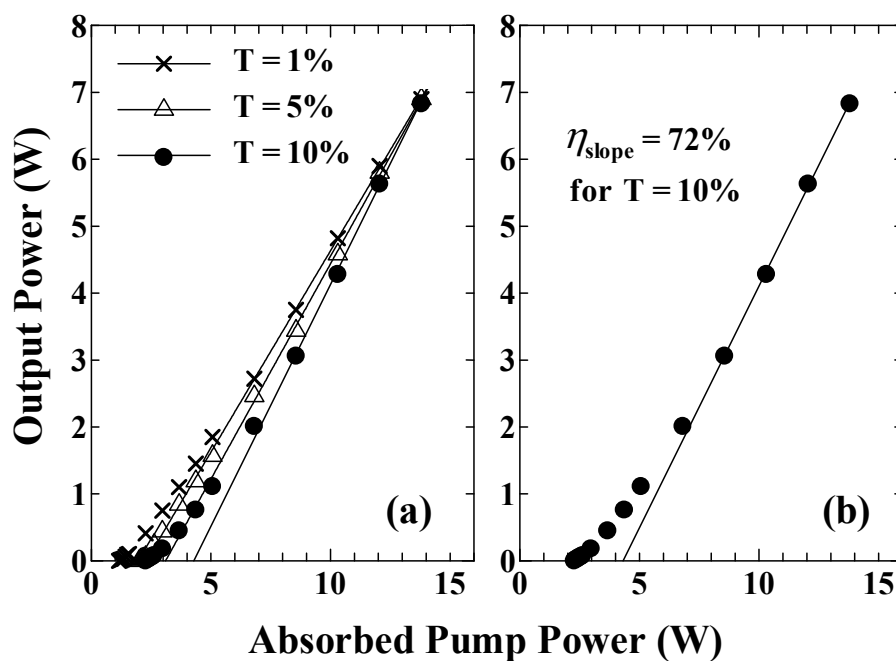


Fig. 2. Input-output power dependence of efficient Yb:YAG ($C_{\text{Yb}} = 9.8\text{ at.}\%$) ceramic laser: (a) with output couplers of $T = 1, 5$, and 10%, and (b) only $T = 10\%$ (Nakamura et al., 2008b).

2.2 Experimental results of linear cavity

Figure 2(a) shows the output power as function of the absorbed pump power in the cases for the three transmittances of the output couplers $T = 1, 5$, and 10%, (Nakamura et al., 2008b) and Fig. 2(b) shows the output power as function of the absorbed pump power only for the case of $T = 10\%$ (Nakamura et al., 2008b). The absorbed pump powers at the lasing threshold were 1.2, 2.0, and 2.3 W, and the maximum output powers of 6.9, 6.9, and 6.8 W for $T = 1, 5$, and 10%, respectively, were obtained at the absorbed pump power of 13.8 W. The round trip loss L in the resonator was estimated to be 0.09 by the lasing thresholds and the reflectivity of the output couplers (Koechner, 1992), which resulted in a small signal gain g_0 of 2.0 cm^{-1} , and a single pass gain G of 1.2 with the 1 mm thick gain medium. Each linear line was fit in Fig. 2(a) for $T = 1, 5$, and 10%. The slope efficiencies η_{slope} were 60, 64, and 72% for $T = 1, 5$, and 10%, respectively. Since we considered that $T = 10\%$ is best for obtaining the highest slope efficiency of 72%, we filled the data for the $T = 10\%$ case to Fig. 2(b). The maximum

output power of 6.8 W for $T = 10\%$ was obtained at the absorbed pump power of 13.8 W, indicating that the efficiency of converting pumping optical power to output optical power, $\eta_{\text{opt-opt}}$, was 49%. The line of the best fit is shown in Fig. 2(b). The slope efficiency η_{slope} was 72% for $T = 10\%$. The maximum output power of 6.8 W was determined to be fourfold higher than 1.7 W and the slope efficiency was determined to be 7% lower than 79% using the 100 μm fiber-coupled LD reported by Dong et al. (Dong et al., 2006). Our 6.8 W laser with the slope efficiency of 72% is expected to have a higher slope efficiency than the present result if the pumping source is replaced with a 100 μm fiber-coupled 25 W LD, for example, LIMO25-F100-DL940 (Lissotschenko Mikrooptick) while maintaining the high output power, because the pumping intensity would increase to a value of fourfold higher than that of a 200 μm fiber-coupled LD. In comparing our laser with the edge-pumped composite Yb:YAG ceramic laser (Tsunekane&Taira, 2007) developed by Tsunekane and Taira, we limit our discussion to the cw case; the laser power of 414 W obtained by Tsunekane and Taira is much higher than our result, but their slope and optical-optical conversion efficiency were 47 and 44%, which were 25 and 5% lower than our slope and optical-optical conversion efficiency of 72 and 49%, respectively. The transverse intensity profile of the Yb:YAG ceramic laser beam is shown in Fig. 3. The intensity distribution indicates that the beam is a Gaussian beam (a TEM_{00} mode beam). The beam image in Fig. 3 was as stable as the pumping LD and we found no amplitude instability.

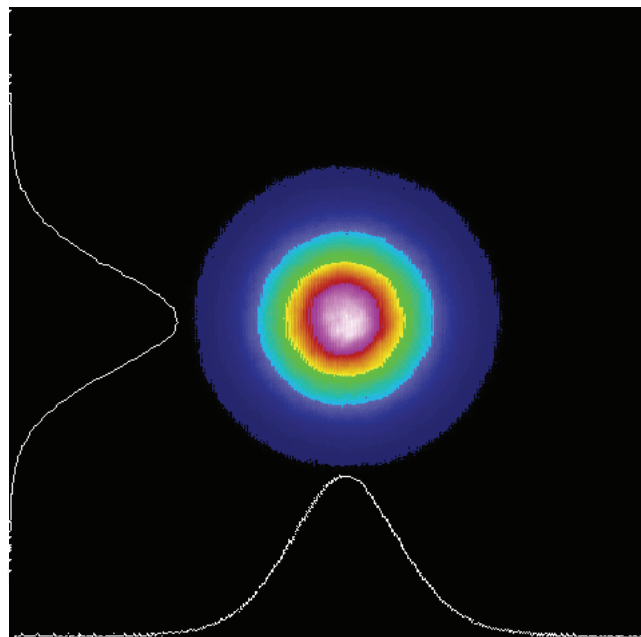


Fig. 3. Transverse intensity profile of the Yb:YAG ceramic laser beam (Nakamura et al., 2008b). The intensity distribution indicates that the beam is a Gaussian beam (a TEM_{00} mode beam).

These results of high output power, high efficiency, and good beam quality were obtained after the optimization of the cavity length. The cavity length was varied to obtain an optimum value for the highest efficiency and highest output power, and the focal length of the thermal lens for designing a tunable laser cavity configuration. Figure 4 shows the maximum output power as a function of the cavity length. Figure 4 shows that the optimum cavity length is less than 20 mm. This value is the appropriate cavity length for our laser,

because there is no space to reduce the cavity length less than 20 mm. When we used the 400 μm fiber-coupled LD (Nakamura et al., 2008a), the optimum cavity length with the highest output power and highest slope efficiency was 25 mm, and reducing the length less than 25 mm yielded a worse result.

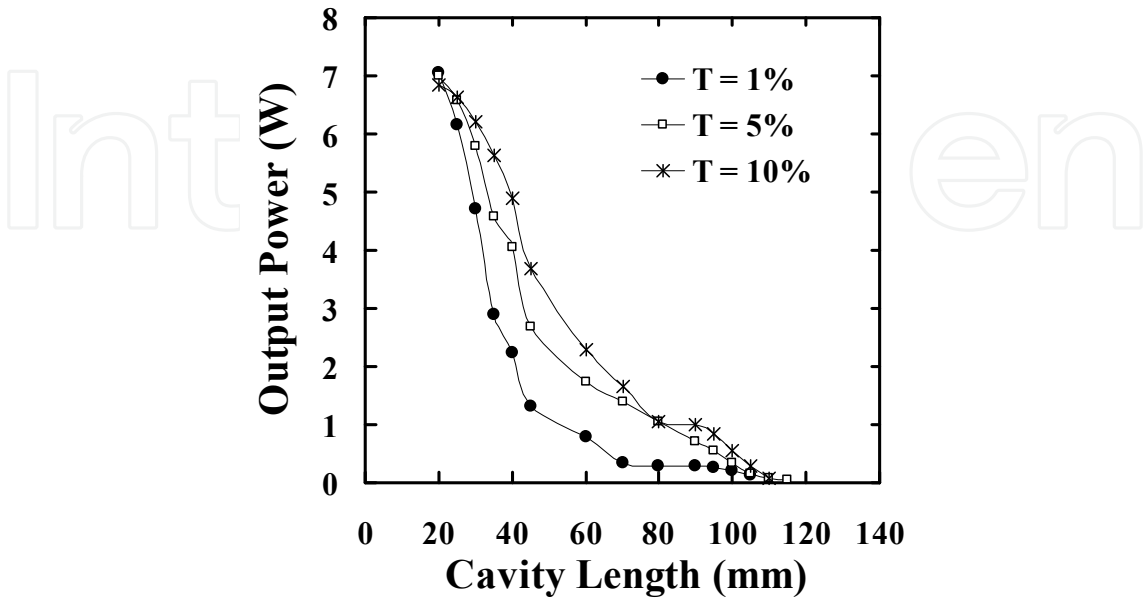


Fig. 4. Maximum output power as function of cavity length of efficient Yb:YAG ceramic laser (Nakamura et al., 2008b).

The focal length of the thermal lens in the ceramic Yb:YAG plate was considered for designing a tunable laser cavity configuration. Figure 4 also shows that the focal length of the thermal lens is 109 mm [120 mm (the cavity length) minus 11 mm (the distance of the ceramic Yb:YAG and the DM)], because the cavity becomes unstable, terminating the laser oscillation when the Fabry-Perot cavity length exceeds the thermal lens focal length. By considering this thermal lens, a tunable laser with a v-shape cavity including a concave mirror M (radius of curvature, ROC = 250 mm) and an SF10 dispersive prism was obtained, as shown in Fig. 5 in section 3.

2.3 Summary of linear cavity

A diode-end-pumped high-efficiency high-power Yb:YAG ceramic laser was demonstrated at a room temperature of 20°C with an Yb concentration of 9.8 at.%, a gain medium thickness of 1 mm, a pumping power of 13.8 W, an output coupler of T = 10%, and a cavity length of 20 mm. A 6.8 W cw output power was obtained with a slope efficiency of 72%. This is the relatively high efficiency of ceramic Yb:YAG lasers at room temperature. The beam quality was shown as a transverse intensity distribution indicating a Gaussian beam (a TEM₀₀ mode beam).

3. Tunability of Yb:YAG ceramic laser

3.1 Tunable Yb:YAG laser

In previous reports on Yb:YAG ceramic lasers, no descriptions of the tunability of the lasers are given. However, there are some reports on the tunability of Yb:YAG crystal lasers.

(Brauch et al., 1995), (Saikawa et al., 2000), (Saikawa et al., 2002) In 2000, the widest tunability range from 1024.1 to 1108.6 nm was demonstrated for a 160 mW Yb:YAG crystal laser using a 0.1% output coupler and a birefringent filter by Saikawa et al. (Saikawa et al., 2000) Subsequently, Saikawa et al. reported a 180 mW Yb:YAG crystal laser with a tunability in the spectral range of 59 nm from 1021.9 to 1081.2 nm in 2002. (Saikawa et al., 2002) Very recently, we have reported the first study of a tunable Yb:YAG ceramic laser (Nakamura et al., 2008b) that exhibits continuous tunability with a spectral range of 63.5 nm from 1020.1 to 1083.6 nm for $T = 1\%$ at a maximum output power of 1.6 W. In this section, we investigated the broader tunability range of a 1.4 W Yb:YAG ceramic laser using a 1% output coupler and a prism; this range was twofold broader than that in our previous work. (Nakamura et al., 2008b)

Comparing the emission spectra of a Yb:YAG crystal (Dong et al., 2003) doped with 10 at.% Yb^{3+} and the ceramic Yb:YAG (Dong et al., 2006) with 9.8 at.% Yb^{3+} reported by Dong et al., we found that the ceramic Yb:YAG had a slightly higher emission intensity and the possibility of producing shorter-wavelength oscillation than the crystal Yb:YAG. To consider the difference in the emission spectra in the short-wavelength region, we employed a dichroic mirror (DM; Layertec No. 102927) with a high-reflectivity (HR) wavelength range from 1010 to 1200 nm and an antireflective range from 800 to 970 nm as a replacement for the previously used DM (Layertec No. 103542) (Nakamura et al., 2008b) with an HR range from 1020 to 1200 nm and an antireflective range from 808 to 980 nm. To the best of our knowledge, we obtained the broadest reported tunable range of 118.31 nm from 992.52 to 1110.83 nm for a Yb:YAG laser owing to this modification.

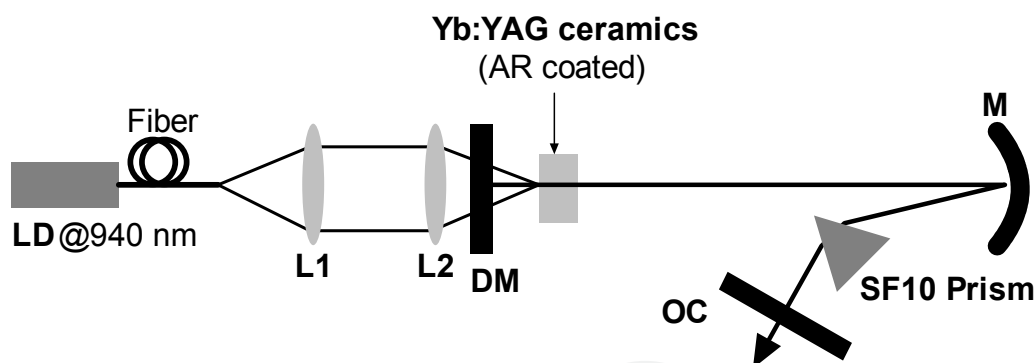


Fig. 5. Experimental setup for tunable Yb:YAG ceramic laser (Nakamura et al., 2008b). LD: fiber-coupled diode laser; L1, L2: focusing lenses; DM: flat dichroic mirror; M: a concave mirror (ROC = 250 mm); OC: output coupler.

3.2 Experimental setup of tunable Yb:YAG ceramic laser

The experimental setup for the Yb:YAG ceramic laser is shown in Fig. 5. A 940 nm fiber-coupled LD (JENOPTIK Laserdiode JOLD-30-FC-12) was used as a pumping source; the fiber has a core diameter of 200 μm and a numerical aperture (NA) of 0.22. The pumping beam was focused onto the ceramic at a ratio of 1:1 using the lenses L1 ($f = 25$ mm) and L2 ($f = 25$ mm). The diameter of the spot focused on the ceramic was ~ 200 μm . A laser cavity consisting of a flat DM and a flat output coupler (OC) with a mirror M and an SF10 prism was used. The DM was antireflection (AR)-coated at 940 nm and had an HR at 1030 nm. In detail, we employed a DM (Layertec No. 102927) with an HR wavelength range from 1010 to 1200 nm and an AR range from 800 to 970 nm, which was a replacement for the previously

used DM (Layertec, No. 103542) (Nakamura et al, 2008b) with an HR range from 1020 to 1200 nm and an AR range from 808 to 980 nm. The OC was partially reflection-coated with transmittances of $T = 0.1, 1, 5$, and 10% at 1030 nm. An AR-coated ceramic Yb:YAG ($C_{Yb} = 9.8$ at.%, Konoshima Chemical) with dimensions of $5 \times 10 \times 1$ mm³ was used. The dimensions of the ceramics were selected to enable Brewster angled alignment for a mode-locked laser performance in future to reduce cost. A 1-mm-thick Yb:YAG ceramic plate was wrapped with indium foil and mounted in a water-cooled copper block that acted as a heat sink. The water was maintained at room temperature (20°C) during laser oscillation. By considering the thermal lens effect (Nakamura et al, 2008b), a tunable laser with a V-shape cavity including a concave mirror M (radius of curvature, $ROC = 250$ mm) and an SF10 dispersive prism was obtained, as shown in Fig. 5. The SF10 dispersive prism was inserted into a V-shape resonator as the tuning element between the folded mirror M and the output coupler OC at the Brewster angle. The cavity length was 315 mm.

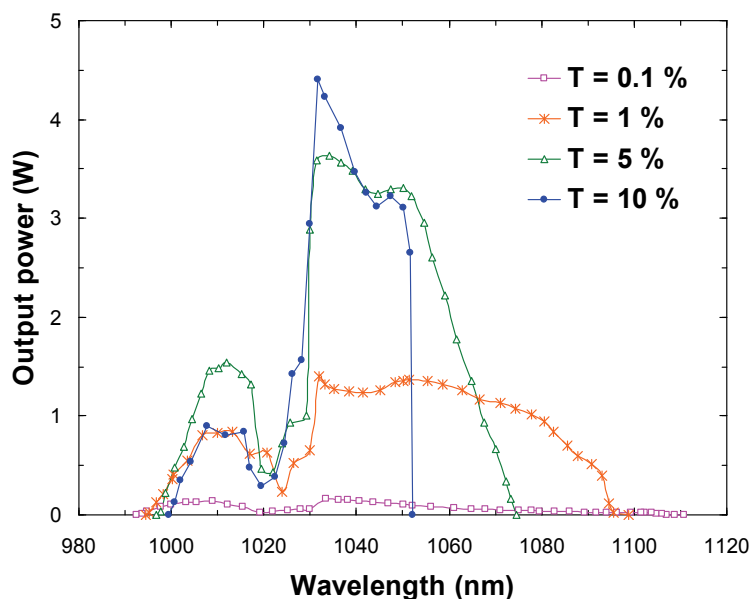


Fig. 6. Output power of ceramic Yb:YAG tunable laser as a function of oscillation wavelength for various output couplers (Nakamura et al., 2009).

3.3 Experimental results of tunable Yb:YAG ceramic laser

Figure 6 shows the dependence of output power on laser oscillation wavelength for output couplers of $T = 0.1, 1, 5$, and 10% when the absorbed pump power was 13.8 W. The averaged spectral linewidth of the output was 3.1 nm in these measurements. We achieved smoothly continuous tunability. A tunable range of 52.64 nm from 999.48 to 1052.12 nm was obtained with a maximum output power of 4.41 W at 1031.78 nm for $T = 10\%$. A tunable range of 77.83 nm from 996.75 to 1074.58 nm and a maximum power of 3.63 W at 1034.14 nm were obtained for $T = 5\%$. A maximum power of 1.40 W was obtained at 1033.42 nm and a tunable range of 104.52 nm from 994.35 to 1098.87 nm was obtained for $T = 1\%$, which is broader than the tunability of 84.5 nm reported by Saikawa et al. (Saikawa et al., 2000). Finally, as shown in Fig. 6, $T = 0.1\%$ results in the largest tunable range of 118.31 nm from 992.52 to 1110.83 nm, which is broader than the tunability for $T = 1\%$, with a maximum output power of 163 mW at 1033.42 nm. To the best of our knowledge, this is the broadest

tunability (Nakamura et al, 2009) reported for an Yb:YAG ceramic laser. The shortest wavelength of the tuning range in Fig. 6 is not limited to 1020 nm (Nakamura et al, 2008b) because of the broad dichroic coating range of the pumping mirror (Layertec No. 102927) with an HR wavelength range from 1010 to 1200 nm and an AR range from 800 to 970 nm. This tuning range of 118.31 nm for the Yb:YAG ceramic laser at 20 °C or 293 K is 3.29-fold broader than the 36.0 nm tuning range, from 1018 to 1054 nm, which was produced from a Yb:YAG crystal laser with a three-plate birefringent filter at 218 and 245 K. (Brauch et al., 1995). Furthermore, our tuning range of 118.31 nm from 992.52 to 1110.83 nm for the high-power ceramic Yb:YAG laser at 20 °C is twofold broader than that for the crystal Yb:YAG laser reported by Saikawa et al. (Saikawa et al., 2002), which has a tuning range of 59 nm from 1022 to 1081 nm at 18 °C. The widely tunable Yb:YAG crystal laser with birefringent filters reported by Saikawa et al. (Saikawa et al., 2000) had a tuning range of 84.5 nm from 1024.1 to 1108.6 nm; however, the highest output power was 180 mW, which is much lower than the maximum output power of 1.4 W of our ceramic laser with an SF10 prism for $T = 1\%$, and our tuning range of 104.52 nm for $T = 1\%$ is 1.24-fold broader than their 84.5 nm tunable range.

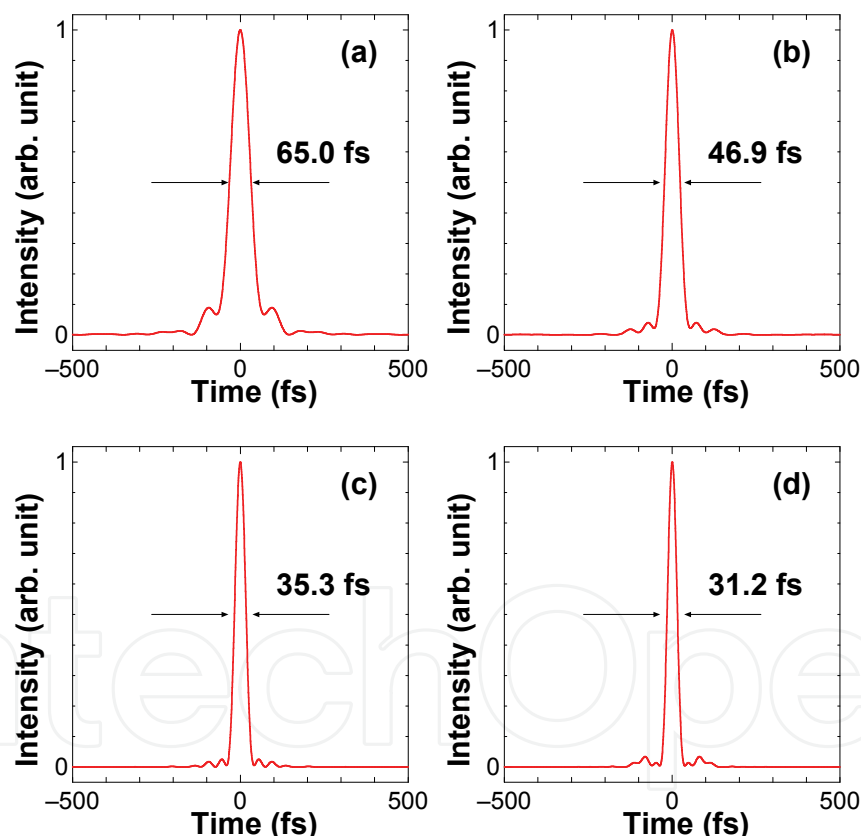


Fig. 7. Expected transform limited pulses numerically obtained by inverse Fourier transformation of Fig. 6 for $T =$ (a) 10, (b) 5, (c) 1, and (d) 0.1% (Nakamura et al., 2009).

Figure 7 shows the expected transform-limited pulses obtained by an inverse Fourier transformation of Fig. 6 for $T =$ (a) 10, (b) 5, (c) 1, and (d) 0.1%. The full width at half maximum (FWHM) of the numerically obtained pulse duration was $t_p = 65.0, 46.9, 35.3,$ and 31.2 fs for Figs. 7(a)-7(d), respectively. These numerical results suggest that the ceramic Yb:YAG laser medium has the potential for a laser performance of less than 65 fs.

3.4 Summary of tunable Yb:YAG ceramic laser

A diode-end-pumped broadly tunable Yb:YAG ceramic laser was demonstrated at room temperature (20 °C) with Yb concentration of 9.8 at.%, a gain medium thickness of 1 mm, a pumping power of 13.8 W. Continuous tunability with a spectral range of 104.52 nm from 994.35 to 1098.87 nm for $T = 1\%$ at a maximum output power of 1.4 W and with a spectral range of 118.31 nm from 992.52 to 1110.83 nm for $T = 0.1\%$ at a maximum output power of 0.16 W was also obtained at room temperature. To the best of our knowledge, this is the broadest tunability reported for a Yb:YAG laser. This tunability is expected to be very attractive for femtosecond laser applications. Numerical results suggest that the ceramic Yb:YAG laser medium has the potential for a laser performance of less than 65 fs. The cost of Yb:YAG ceramic laser materials is much lower than that of single crystals because of their high-speed, large production, and mass production potential, is making them tremendously attractive for industrial applications.

4. Mode-locked Yb:YAG ceramic laser

4.1 Mode-locked Yb-doped lasers

Femtosecond mode-locked lasers are applied in various fields of physics, engineering, chemistry, biology and medicine, with applications including ultrafast spectroscopy, metrology, superfine material processing and microscopy. Specific and interesting properties of the femtosecond laser pulses have been used in these applications. For example, femtosecond pulses have a very precise time resolution, and their strong electric field induces important and unusual nonlinear effects. For those applications, high-power, high-efficiency and compact femtosecond lasers are required. Ceramic materials are attractive for satisfying these requirements. YAG ceramics have 10% higher hardness than a YAG single crystal, and the fracture toughness of YAG ceramics is more than threefold that of the YAG single crystal. Therefore, the ceramics have a higher resistance to thermal shock than the single crystal. Ytterbium (Yb^{3+}) also has interesting properties satisfying the above requirements. Its broad absorption and emission spectra allow the realization of a directly laser-diode (LD)-pumped femtosecond laser. Moreover, its small quantum defect, absence of excited-state absorption, upconversion and cross-relaxation reduce the thermal load and enable highly efficient operation. The emission and absorption spectra and thermal conductivity strongly depend on the host material.

Recently, various ceramic materials have been progressively investigated for use in ultrashort-pulse lasers (Shirakawa et al., 2003), (Tokurakawa et al., 2006), (Shirakawa et al., 2003), (Tokurakawa et al., 2007), (Tokurakawa et al., 2008). A diode-pumped femtosecond Yb:Y₂O₃ ceramic laser was demonstrated, and 615 fs pulses at a center wavelength of 1076.5 nm were obtained with a 420 mW average power (Shirakawa et al., 2003). A diode-pumped passively mode-locked Yb:Lu₂O₃ ceramic laser was demonstrated, for which 357 fs pulses at a center wavelength of 1033.5 nm with a 352 mW average power were obtained (Tokurakawa et al., 2006). A passively mode-locked femtosecond Yb³⁺-doped Y₃(Sc_{0.5}Al_{0.5})₂O₁₂ (Yb:YSAG) ceramic laser pumped by a Ti:sapphire laser was also demonstrated, and 280 fs pulses at a center wavelength of 1035.8 nm with a 62 mW average power were obtained (Saikawa et al., 2007), but the laser was not diode-pumped. A diode-pumped Kerr lens mode-locked laser of Yb³⁺:Sc₂O₃ ceramics was demonstrated with 92 fs pulses at a center wavelength of 1042 nm and an 850 mW average power, and 90 fs pulses at a center wavelength of 1092 nm and a 160 mW average power (Tokurakawa et al., 2007). A

was 0.22. The maximum pump power was 26.6 W. The pumping beam was imaged by relay to the ceramics using lens L1 ($f = 50$ mm) and lens L2 ($f = 70$ mm). The 1-mm-thick Yb:YAG ($C_{Yb} = 9.8$ at.%) ceramic plate was arranged at the Brewster's angle. The Yb:YAG plate was wrapped with indium foil and mounted in a water-cooled copper heat sink block. The copper block was cooled by flowing water at 20°C. The ceramic was placed between two high-reflectivity mirrors (M1, M2) that were anti-reflection (AR)-coated at 940 nm and had high reflectivity at 1030 nm with a 100 mm radius of curvature (ROC). The reference laser operated in cw mode with a 10% output coupler (OC1) and a high-reflectance end mirror (M3) along the dashed lines in Fig. 8. For passive mode locking, a 1 or 0.1% output coupler (OC2) and a semiconductor saturable absorber mirror (SESAM, BATOP) with 2% saturable absorption at 1030 nm, 70 $\mu\text{J}/\text{cm}^2$ saturation fluence and 500 fs relaxation time constant were used in the respective arms. The total cavity length was 1620 mm. The laser beam was focused onto the SESAM by a concave mirror (M4, ROC = 250 mm). The distance among the mirrors and the folded angle of the mirrors are shown in Fig. 8. The astigmatism compensation was not considered. The spot sizes of laser mode in the laser crystal and on SESAM were estimated $\sim 61 \times 53$ μm and $\sim 450 \times 330$ μm , respectively. An SF10 Brewster prism pair (P1, P2) with 465 mm separation was inserted in the other arm to compensate for the dispersion. The total negative GDD of this cavity was about -2670 fs^2 per a round trip.

4.3 Experimental results of mode-locked Yb:YAG ceramic laser

We demonstrated mode locking for three cases. The first case was mode locking at 1030 nm using a 1% output coupler. The second case was mode locking at 1030 nm using a 0.1% output coupler. The third case was mode locking at 1050 nm region with cw oscillation at 1030 nm region using a 0.1% output coupler.

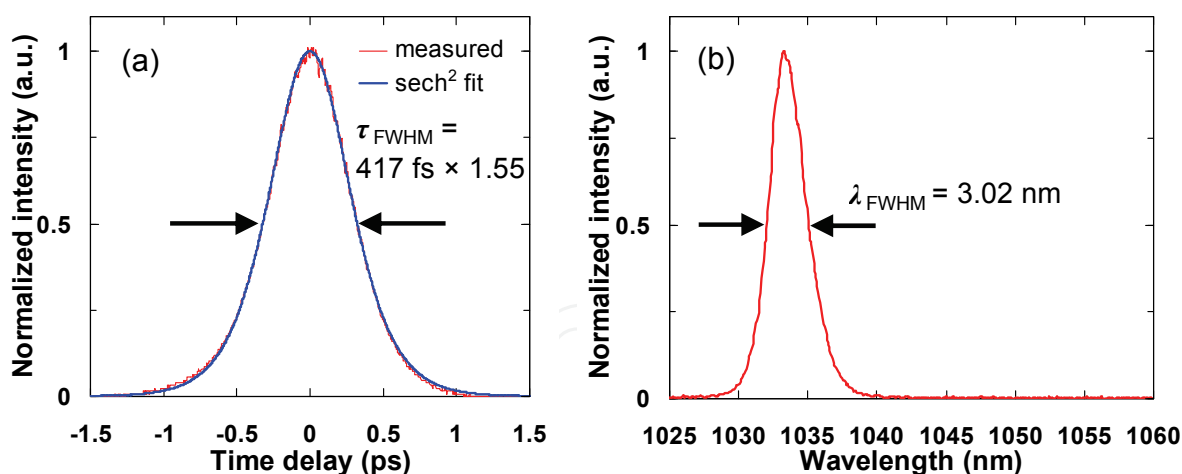


Fig. 9. (a) Measured autocorrelation trace and the sech² fitting, and (b) spectrum of the mode-locked pulses using a 1% output coupler (Yoshioka et al., 2009).

4.3.1 Mode locking at 1030 nm region using a 1% output coupler

Figure 9 shows the intensity autocorrelation trace and the spectrum of mode-locked pulses using a 1% output coupler. The average output power was 250 mW at a pump power of 26.6 W. The sech²-fitted pulse width was 417 fs and the spectral width was 3.02 nm at the center wavelength of 1033.3 nm, which results in a time-bandwidth product of 0.353, slightly above

the Fourier limit for a sech^2 pulse (0.315). The repetition rate was 91 MHz. The pulse energy and peak power were 2.75 nJ and 6.60 kW, respectively. When the 1% output coupler was used, mode locking at 1050 nm region was not observed but cw oscillation at 1050 nm region or simultaneous cw dual-wavelength oscillation at 1030 and 1050 nm was observed.

4.3.2 Mode locking at 1030 nm region using a 0.1% output coupler

Figure 10 shows the intensity autocorrelation trace and the spectrum of mode-locked pulses using a 0.1% output coupler. The average output power was 25 mW at a pump power of 26.6 W. The sech^2 -fitted pulse width was 286 fs and the spectral width was 4.51 nm, and centered at 1033.5 nm. This results in a time-bandwidth product of 0.363, which is 15% above the Fourier limit for a sech^2 pulse (0.315). This indicates that the pulse has potential to be further shortened by extracavity compression. The repetition rate was 91 MHz. The pulse energy and peak power were 0.275 nJ and 0.960 kW, respectively. The pulse width of 286 fs obtained using a 0.1% output coupler (Fig. 10) was shorter than that of 417 fs obtained using a 1% output coupler (Fig. 9) because the intracavity power is increased using a 0.1% output coupler, which makes it possible for the laser to oscillate with a broad spectral range. However, the average output power decreased to 25 mW owing to the low transmission of the 0.1% output coupler.

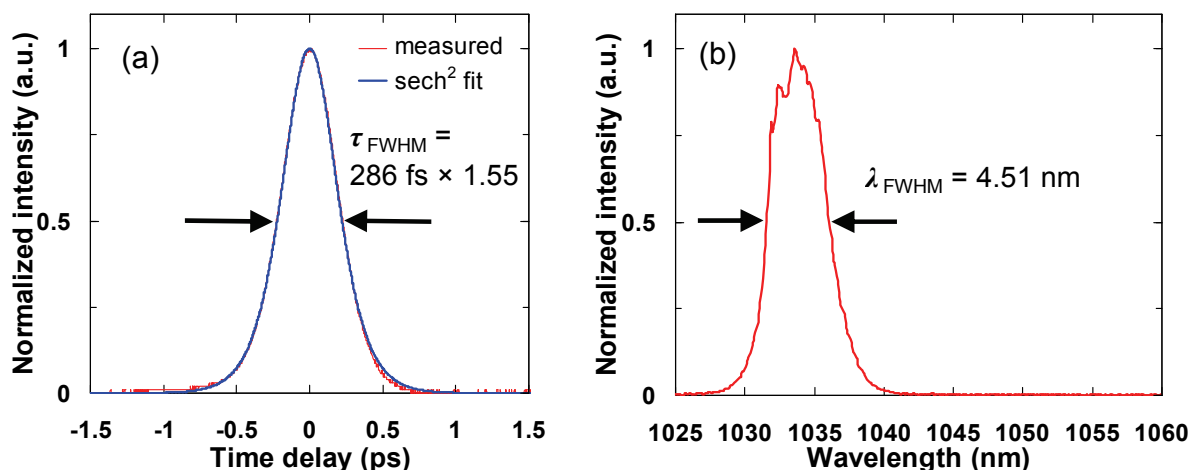


Fig. 10. (a) Measured autocorrelation trace and the sech^2 fitting, and (b) spectrum of the mode-locked pulses at 1030 nm region using a 0.1% output coupler (Yoshioka et al., 2009).

4.3.3 Mode locking at 1050 nm region using a 0.1% output coupler

When the SESAM angle was changed slightly from the conditions described in Sec. 4.3.2, mode locking at 1033.5 nm suddenly stopped and simultaneous cw dual-wavelength oscillation at 1030 and 1050 nm occurred. When the intensity of the cw spectral component at 1050 nm was increased and the laser mode radius focused on the SESAM was optimized, mode-locked pulses were generated again. Figure 11 shows (a) the intensity autocorrelation trace, (b) the spectrum of mode-locked pulses. The output spectrum in Fig. 11(b) was measured by injection into the fiber-coupled spectrometer directly from OC2. The average output power was 20 mW at a pump power of 26.6 W. The sech^2 -fitted pulse width was 233 fs and the spectral width was 5.20 nm and centered at 1048.3 nm, which results in a time-bandwidth product of 0.330, slightly above the Fourier limit for a sech^2 pulse (0.315). The

repetition rate was 91 MHz. The pulse energy and peak power were 0.220 nJ and 0.946 kW, respectively. Figure 12 shows a pulse train of cw mode-locking in millisecond time scale. This indicates that the mode-locked pulses had a long-term stability. The pulse width of 233 fs at 1048.3 nm in Fig. 11 obtained using a 0.1% output coupler was the shortest among these results. This indicates that the laser has the potential to generate shorter pulses at 1050 nm region rather than 1030 nm region. However, the cw spectral component at 1032.4 nm was not quenched in our cavity during mode locking at 1048.3 nm. This occurred owing to the strong emission at 1030 nm.

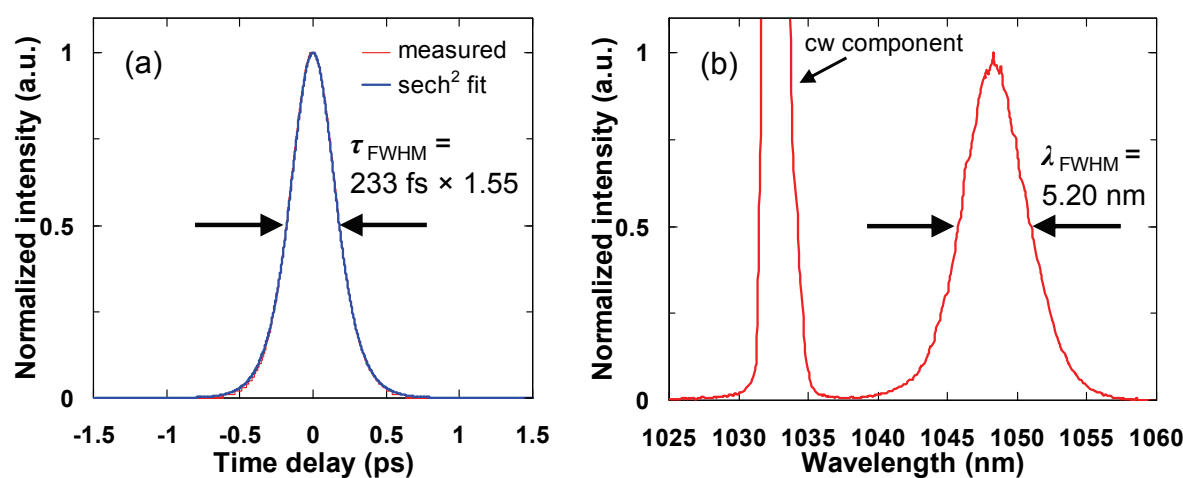


Fig. 11. (a) Measured autocorrelation trace and the sech² fitting, (b) spectrum of the mode-locked pulses at 1050 nm region using a 0.1% output coupler (Yoshioka et al., 2009).

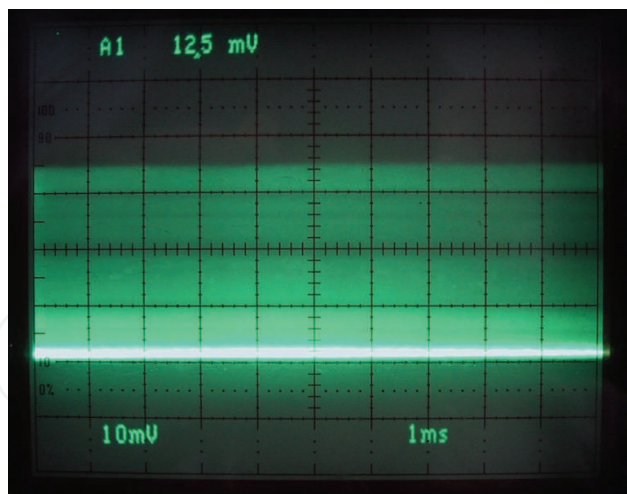


Fig. 12. Pulse train of cw mode-locking in millisecond time scale (Yoshioka et al., 2009).

4.4 Summary of mode-locked Yb:YAG ceramic laser

A diode-pumped femtosecond ytterbium laser with a host material of YAG ceramic was demonstrated. We successfully achieved passive mode locking at wavelengths of 1033.5 and 1048.3 nm. At 1033.5 nm, passive mode locking by a semiconductor saturable absorber mirror generated 286 fs pulses with an average power of 25 mW using a 0.1% output coupler. This is the shortest pulse width in the 1030 nm region. At 1048.3 nm, the shortest

pulse of 233 fs with an average output power of 20 mW was generated using a 0.1% output coupler. To the best of our knowledge, this is the first mode-locked Yb:YAG ceramic laser, and the shortest pulse for diode-pumped ceramic lasers without Kerr lens mode locking was obtained.

5. Summary

A high-power efficient ceramic Yb:YAG laser was demonstrated at a room temperature of 20°C with an Yb concentration of 9.8 at.%, a gain medium of 1 mm, a pumping power of 13.8 W, an output coupler of $T = 10\%$, and a cavity length of 20 mm. A 6.8 W cw output power was obtained with a slope efficiency of 72%.

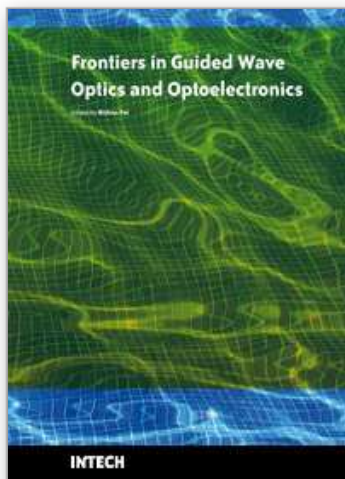
A diode-end-pumped Yb:YAG ceramic laser with a broad tunability was demonstrated. The continuous tunability in the spectral range of 104.52 nm from 994.35 to 1098.87 nm for $T = 1\%$ at a maximum output power of 1.4 W and in the spectral range of 118.31 nm from 992.52 to 1110.83 nm for $T = 0.1\%$ at a maximum output power of 0.16 W was also obtained at room temperature. To the best of our knowledge, this is the broadest tunability of Yb:YAG lasers. A diode-pumped passively mode-locked Yb:YAG ceramic laser was demonstrated. 417 and 286 fs pulses with average powers of 250 and 25 mW were obtained at 1030 nm using 1 and 0.1% output couplers, respectively. 233 fs pulses with an average power of 20 mW were also obtained at a center wavelength of 1048.3 nm using a 0.1% output coupler. To the best of our knowledge, this is the first demonstration of a diode-pumped mode-locked Yb:YAG ceramic laser.

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As the editor, I feel extremely happy to present to the readers such a rich collection of chapters authored/co-authored by a large number of experts from around the world covering the broad field of guided wave optics and optoelectronics. Most of the chapters are state-of-the-art on respective topics or areas that are emerging. Several authors narrated technological challenges in a lucid manner, which was possible because of individual expertise of the authors in their own subject specialties. I have no doubt that this book will be useful to graduate students, teachers, researchers, and practicing engineers and technologists and that they would love to have it on their book shelves for ready reference at any time.

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Phone: +86-21-62489820
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