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## Luminescent Devices Based on Silicon-Rich

## **Dielectric Materials**

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Additional information is available at the end of the chapter

http://dx.doi.org/10.5772/64894

#### Abstract

Luminescent silicon-rich dielectric materials have been under intensive research due to their potential applications in optoelectronic devices. Silicon-rich nitride (SRN) and silicon-rich oxide (SRO) films have been mostly studied because of their high luminescence and compatibility with the silicon-based technology. In this chapter, the luminescent characteristics of SRN and SRO films deposited by low-pressure chemical vapor deposition are reviewed and discussed. SRN and SRO films, which exhibit the strongest photoluminescence (PL), were chosen to analyze their electrical and electroluminescent (EL) properties, including SRN/SRO bilayers. Light emitting capacitors (LECs) were fabricated with the SRN, SRO, and SRN/SRO films as the dielectric layer. SRN-LECs emit broad EL spectra where the maximum emission peak blueshifts when the polarity is changed. On the other hand, SRO-LECs with low silicon content (~39 at.%) exhibit a resistive switching (RS) behavior from a high conduction state to a low conduction state, which produce a long spectrum blueshift (~227 nm) between the EL and PL emission. When the silicon content increases, red emission is observed at both EL and PL spectra. The RS behavior is also observed in all SRN/SRO-LECs enhancing an intense ultraviolet EL. The carrier transport in all LECs is analyzed to understand their EL mechanism.

**Keywords:** silicon-rich dielectrics, photoluminescence, electroluminescence, conduction mechanisms



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### 1. Introduction

The use of photonic signals instead of electrons to transmit information through an electronic circuit is an actual challenge. Unfortunately, it is well known that bulk silicon (Si) is an indirect bandgap semiconductor, making it an inefficient light emitter. Therefore, great efforts have been taken to obtain highly luminescent Si-based materials in order to get Sibased photonic devices, especially a light emitting device [1–3]. Such circumstances have led to explore new options for converting silicon into a luminescent material. Si nanoparticles (Si-nps) embedded in a dielectric material as silicon-rich oxide (SRO) or silicon-rich nitride (SRN) show a prominent photoluminescence (PL) emission in red and blue-green region, respectively [4–10]. Thus, SRN or SRO films have been considered as promising candidates for emissive materials due to their potential applications in Si-based optoelectronic devices, and their fully compatibility with the complementary metal-oxide-semiconductor (CMOS) processes [11–16].

Two main strategies have been explored: those that focus on the intrinsic emission from the matrix, either through emissions from defects or by the presence of Si-nps. The second one focus on extrinsic emission, which is produced by doping the material (usually introducing rare earth ions) [3, 17].

The most common strategy to obtain intrinsic emission is through silicon nanostructures, which significantly increase the emission due to the quantum confinement effect (QCE) [5, 18]. Furthermore, the dependence on the size of the Si-nps on the forbidden gap width allows that the emission can be adjusted in the visible and the red-near infrared region of the electromagnetic spectrum [19]. Several studies have reported both red and near infrared electroluminescence (EL) SRO which is mainly attributed to the recombination of excitons in Si-nps [20, 21]. On the other hand, the emission of green or blue light has been attributed to defects associated with oxygen [22, 23]. Some studies have reported that red (620 nm) EL emission could be attributed to non-bridging oxygen hole center (NBOHC) defects whose origin has been corroborated by the fact that the peak position does not change if the film is excited with different energies [24].

Another alternative to obtain intrinsic light emission is through an ordered structure of Sinps by a superlattice, which is formed by the alternating of SRO and SiO<sub>2</sub> nano-films. Red or near infrared emission has been observed in these structures and has been related to both excitonic recombination taking place in confined states within Si-nps or relaxation of hot electrons [25, 26].

Intrinsic emission has been also observed in SRN films [27–31]. For example, an orange emission at 600 nm was observed at room temperature and has been related to the electron-hole pairs' recombination within Si-nps [27]. Also, green emission has been observed in nitrogenrich silicon nitride, which was attributed to radiative recombination in localized states related to Si-O [28]. Some other authors have shown significant improvement of the green emission intensity using oxidized silicon-rich nitride [29]. Also, when a silicon nitride film is implanted with Si ions, violet and green-yellow emission bands are observed, giving rise to an intense

white EL emission [30]. The violet band was related with the presence of defects states related to silicon dangling bonds (=Si<sup>0</sup> or centers K<sup>0</sup>) located near the middle of the forbidden gap of silicon nitride and defect states related to the unit Si–Si= located near the edge of the valence band, while the green-yellow band was attributed to the transition from the =Si<sup>0</sup> state to nitrogen dangling bonds (=N–) in the tails of valence bands.

Red-near infrared EL (800 nm) has been reported in superlattices combining SRN and SiO<sub>2</sub> films and explained by the bipolar recombination of electron-hole pairs in Si-nps present within the SRN films [31]. A yellow EL emission has been reported when an SRO film instead of SiO<sub>2</sub> layer is used in the multilayer structure [32]. All of these promising results have proved the first implemented all-silicon-based photonic device [33]. Nevertheless, despite all these promising results in luminescent silicon-based materials, the improvement of the efficiency of the light emitting devices is still necessary.

This chapter shows a review about our experience on the PL and EL properties of SRN and SRO films deposited by low-pressure chemical vapor deposition (LPCVD). The effect of the combination of the SRN and SRO luminescent properties is also analyzed as an SRN/SRO structure. A study about the composition, structural, optical, and electro-optical properties of these films will be discussed. The study also includes the analysis of the charge transport mechanism through the SRO, SRN, and SRN/SRO films to understand their electroluminescence behavior and its correlation with the different luminescent centers (LCs) within the active material.

### 2. Experimental procedure

In this chapter, SRN, SRO, and SRN/SRO films were deposited in a homemade LPCVD hotwall reactor. In these silicon-rich dielectrics materials, the Si content was controlled by a ratio of partial pressure of reactant gases;  $R_{\rm N}$  and Ro for SRN and SRO, respectively.

The SRN films were deposited on N-type ((100)-oriented) Si wafers with a resistivity of 1–5  $\Omega$ -cm at 750°C using ammonia (NH<sub>3</sub>) and 5% nitrogen (N<sub>2</sub>)-diluted silane (SiH<sub>4</sub>) as the reactant gases by the ratio  $R_N = 20^*P(NH_3)/P(SiH_4)$ . SRN films with  $R_N$  values of 5, 20, and 80 were deposited with the parameters shown in **Table 1**. SRO films were deposited on N-type Si wafers with a resistivity of 5–10  $\Omega$ -cm ((100)-oriented) at a temperature of 720°C using pure nitrous oxide (N<sub>2</sub>O) and 3.3% nitrogen (N<sub>2</sub>)-diluted silane (SiH<sub>4</sub>) as the reactant gases by the ratio  $R_0 = 30^*P(N_2O)/P(SiH_4)$ . SRO films with  $R_0$  values of 30 and 20 were deposited with the parameters shown in **Table 1**. Finally, SRN/SRO bilayers were deposited on P-type silicon substrates ((100)-oriented) with resistivity of 5–10  $\Omega$ -cm.  $R_0 = 20$  and 30 and  $R_N = 80$  values were used for the SRO and SRN films, respectively. SRO and SRN films were deposited at 730 and 760°C, respectively. The deposition conditions are also summarized in **Table 1**. About 3.3% of nitrogen diluted silane was used for the bilayer structures.

After deposition, SRN, SRO and SRN/SRO samples were thermally annealed at 1100°C under nitrogen atmosphere conditions for 180 min.

	Sample name	$R_{_{\mathbf{N}}}$	Ro	Pressure of gases (Torr)			Time (min)	Thickness (nm)
				N <sub>2</sub> O	NH <sub>3</sub>	SiH <sub>4</sub>	_	
SRO	M20		20	0.53		0.80	12	$55.72 \pm 5.0$
	M30		30	0.80		0.80	15	$64.00 \pm 3.4$
SRN	N5	5			0.22	0.85	10	$102.83 \pm 3.62$
	N20	20			0.85	0.85	15	112.67 ± 6.19
	N80	80			2.00	0.50	-13	66.93 ± 2.12
SRN/SRO	B20	80			1.08	0.41	4	$16.32 \pm 1.54$
			20	0.53		0.80	15	$55.72 \pm 5.0$
	B30	80			1.08	0.41	4	$16.32 \pm 1.54$
			30	0.80		0.80	15	$64.00 \pm 3.4$

**Table 1.** Process parameters of the SRO, SRN and SRN/SRO films deposited by LPCVD and thickness of the sample after thermal annealing.

For electrical and electroluminescence studies, light emitting capacitive (LEC) structures were fabricated. For SRN-LECs, a transparent 300-nm thick fluorine-doped tin oxide SnO<sub>2</sub>:F (FTO) film was deposited onto the surface of the SRN by ultrasonic spray pyrolysis. Square-shaped patterns with 1 mm<sup>2</sup> area were defined by a photolithography process step to act as gate contact. For SRO-LECs, ~400-nm thick semitransparent n+ polycrystalline silicon (poly) gate was deposited onto the SRO film surface by LPCVD. After a photolithography process step, square-shaped gates of 4 mm<sup>2</sup> area were defined. For SRN/SRO-LECs, ~300-nm thick indium tin oxide (ITO) film was deposited by RF sputtering onto the surface of the films as gate contact. Square-shaped patterns with area of 1 mm<sup>2</sup> were defined by a photolithography process step to act as anode gate contact. Approximately 700-nm thick aluminum (Al) film was evaporated onto the backside of the silicon substrates as cathode contact in all of the LECs. A thermal annealing process at 460°C in N<sub>2</sub> atmosphere for 20 min was used to form the ohmic contact.

The thickness of thermally annealed SRO and SRN films was measured with a Gaertner L117 ellipsometer with a 70° incident laser with wavelength of 632.8 nm and is also shown in **Table 1**. Chemical bonding characteristics was analyzed by means of Fourier transform infra-red spectroscopy (FTIR) with a Brucker V22 equipment in the 4000–350 cm<sup>-1</sup> range with a resolution of 5 cm<sup>-1</sup>. The PL spectra were measured with a Fluoromax 3 of Horiba Jobin Yvon. The samples were excited using a 300 nm radiation, and the PL emission signal was collected from 400 to 900 nm with a resolution of 1 nm. The depth analysis profile of thermally annealed SRN, SRO, and SRN/SRO films was analyzed by means of X-ray photoelectron spectroscopy (XPS) Escalab 250Xi of Thermo Scientific equipment, with an Al K $\alpha$  monochromated source. Current-voltage (I-V) measurements of SRN, SRO, and SRN/SRO-LECs were

performed using a Keithley 4200-SCS parameter analyzer at the same time that the EL was collected with an optical fiber, which was located right on the surface of the device and connected to an Ocean Optics QE-65000 spectrometer.

### 3. Composition

The composition of SRN and SRO films play an important role in order to understand their luminescence, electrical, and electro-optical properties. In this sense, some techniques such as FTIR and XPS spectroscopies have been used.

#### 3.1. Silicon-rich nitride (SRN) film

The Si-N bonds of SRN films were determined by FTIR measurements. **Figure 1(a)** shows the IR spectra measured from SRN films with  $R_N = 5$  (N5), 20 (N20), and 80 (N80) before (B-TA) and after (A-TA) thermal annealing.

IR peaks at 460 and 840 cm<sup>-1</sup> ascribed to Si-N wagging and stretching modes, respectively, were observed for all samples [35–37]. An IR band appears at 1080 cm<sup>-1</sup> after thermal annealing, being more evident in the N80 sample. The presence of this peak has been observed before and attributed to a reordering in the films toward  $a-Si_3N_4$  bonding configuration [38, 39]. Nevertheless, it could be related to the Si–O stretching mode due the oxygen incorporation in the samples.



**Figure 1.** (a) FTIR spectra of SRN films with  $R_N = 5$  (N5), 20 (N20) and 80 (N80) before and after thermal annealing, and (b) atomic concentration (at.%) of Si, O, and N as a function of  $R_N$ . Inset shows the XPS depth profile of SRN film with  $R_N = 80$  (N80). From Cabañas-Tay et al. [34].

In order to comprehend the stoichiometry and the presence of some oxygen into the SRN films, analysis of their composition was performed by means of XPS. **Figure 1(b)** shows information about the chemical composition of the thermally annealed SRN samples. The inset of **Figure 1(b)** exhibits the depth profile composition of the thermally annealed  $R_N$  = 80 (N80) sample. As can be observed, some oxygen is present at the outmost part of the layer through the film depth, being present mainly at the film surface. This behavior is similar for different  $R_{N'}$  but with different concentrations. **Figure 1(b)** shows the atomic concentration of Si, O, and N as a function of  $R_N$ . Mean silicon content values along the layer of  $45.55 \pm 0.38$ ,  $43.58 \pm 0.48$ , and  $42.88 \pm 1.39$  at.% were obtained for the SRN films with  $R_N = 5$ , 20, and 80, respectively. Moreover, the presence of oxygen increases from 5.6 to 11.8 at.% as the  $R_N$  value increases from 5 to 80, respectively.

#### 3.2. Silicon-rich oxide (SRO) film

The Si–O bonds of SRO films were also determined by FTIR measurements. The IR spectra measured from SRO films with Ro = 20 (M20) and Ro = 30 (M30) after thermal annealing are shown in **Figure 2(a)**. Typical vibration bands at 460, 810, and 1070 cm<sup>-1</sup> related to Si–O rocking, Si–O bending, and Si–O stretching modes, respectively, were present in both SRO films [40, 41]. The shoulder from ~1100 to ~1300 cm<sup>-1</sup> observed in both SRO films has been attributed to Si–O stretching out of phase [42]. It has been reported that this shoulder is less pronounced for the suboxides compared to the stoichiometric oxide [43, 44]. The absorption at 610 cm<sup>-1</sup> due to unsaturated Si–Si bonds (phonon-phonon interactions) was observed only in M20 (SRO film with higher proportion of silicon precursor) showing the presence of structural imperfections at the SiO<sub>2</sub>/silicon nanoparticles (Si-nps) interface [45].

**Figure 2(b)** exhibits the depth profile composition of the thermally annealed SRO films. Mean silicon content values of  $41.85 \pm 1.1$  and  $39.98 \pm 0.8$  at.% were obtained for the SRO films with



**Figure 2.** (a) FTIR spectra of SRO films with Ro = 20 (M20) and 30 (M30) after thermal annealing, and (b) XPS depth profile of SRO films with Ro = 20 (M20) and 30 (M30).

*Ro* = 20 (M20) and 30 (M30), respectively. The presence of nitrogen was also observed in both SRO films being slightly more evident for M20. An analysis of the Si2p XPS signal of SRO film with *Ro* = 30 thermally annealed at 1100°C (not showed here) [46] shows that the microstructure of this film is almost a stoichiometric SiO<sub>2</sub> (Si<sup>4+</sup>) and that the silicon excess is present as Si–O compounds, which could include the neutral oxygen vacancy (NOV) and non-bridging oxygen hole center (NBOHC) defects. As expected by the Si2p XPS signal, crystalline silicon nanoparticles were not observed in this films; however, energy-filtered transmission electron microscopy (EFTEM) analysis has shown (although not clearly) the presence of silicon nanoparticles with a mean size of 1.5 nm and density of ~1.1 × 10<sup>12</sup> cm<sup>-2</sup> as reported in [47, 48]. In fact, it has been reported that a largely disordered system is expected when the size of Si-nps is between 1 and 2 nm. Thus, the small Si-nps are no longer crystalline but amorphous like, as also indicated by theoretical calculations [49, 50]. Then, the dispersed Si atoms in the M30 films are preferentially arranged as very small Si-nps (E'<sub>8</sub>) and some defects like oxygen-deficiency center (ODC) and NBOHC as reported for thin SiOx films (<100 nm) [51].

The presence of the Si<sup>0</sup> peak in the Si2p XPS signal correlates well with the presence of silicon nanocrystal observed by high-resolution transmission electron microscopy (HRTEM, not shown here) in the M20 film. HRTEM reveals Si nanocrystals with an average size (and density) of  $2.91 \pm 0.40$  nm ( $8.66 \times 10^{11}$  cm<sup>-2</sup>) [46].

#### 3.3. Silicon-rich nitride/silicon-rich oxide (SRN/SRO) bilayer

The Si–N and Si–O bonds of SRN/SRO bilayers with Ro = 20 (B20) and Ro = 30 (B30) were also determined. **Figure 3(a)** shows the FTIR spectra of SRO and SRN/SRO films with Ro = 20 (M20 and B20, respectively) after thermal annealing. The IR peak at 820 cm<sup>-1</sup> observed only in the B20 bilayer is ascribed to Si–N stretching present in SRN films [45]. The intensity of a shoulder



**Figure 3.** (a) FTIR spectra of SRO and SRN/SRO films with *R*o = 20 (M20 and B20, respectively) after thermal annealing, and (b) XPS depth profile of SRO and SRN/SRO films with *R*o = 20 (M20 and B30, respectively) after thermal annealing.

(shadow area) from 900 to 1000 cm<sup>-1</sup> increased when the B20 bilayer was formed, and it was attributed to the formation of Si–N–O bonds in silicon oxynitride [45].

As shown in **Figure 3(a)**, the shoulder from ~1100 to ~1300 cm<sup>-1</sup> was observed in both samples (M20 and B20) and it was attributed to Si–O stretching out of phase [42]. The IR peak at 610 cm<sup>-1</sup> observed in the SRO monolayer (M20) disappeared for the bilayer (B20) and it could be related to the nitrogen incorporation within the SRO that creates Si–N–O–Si bridges. The presence of these bridges decreased the quantity of strained bonds and Si dangling bonds at the SiO<sub>2</sub>/Si-np interface [45]. The SRN/SRO film with Ro = 30 (B30) showed the same IR peaks than the B20 bilayer, except for the 610 cm<sup>-1</sup> IR peak.

Figure 3(b) showed the depth profile composition of the thermally annealed SRO monolayer and SRN/SRO bilayer, both with Ro = 20 (M20 and B20, respectively). The average Si content within the SRO monolayer was about 41.85±1.1 at.%. When the SRN/SRO bilayer was formed, the oxide-nitride interface became an imprecise oxynitride (SiON) layer almost for 30 s etching time (shadow area in **Figure 3(b)**). In that region, a gradual increasing of both nitrogen and silicon was observed toward the interface. The average Si content slightly increased up to  $\sim$ 43.21 ± 0.7 at.% in the SRO layer, whereas the oxygen content in the SRO film was reduced from  $56.78 \pm 1.3$  to  $53.08 \pm 1.0$  at.%. The Si, O, and N diffusion could be enhanced by the high annealing temperature. This is supported by two facts: first, the presence of a nitrogen concentration within the oxide layer which goes from 2 to 10%, and second, the oxygen presence in the nitride layer (~25%). The nitrogen content inside of the SRO in the SRN/ SRO bilayer structure could modify its optical and structural properties. It has been reported that the nitrogen hinder the diffusion of Si atoms and prevents the phase separation in the amorphous SiOx:N films [52]. Thus, the mobility of the Si atoms is smaller and the growth of the Si-nps during the thermal diffusion process is reduced, giving as a result smaller Si-nps. A similar behavior was observed for the B30 bilayer, but with different concentration.

### 4. Photoluminescence

#### 4.1. Silicon-rich nitride (SRN) film

The PL spectra of SRN films before and after thermal annealing are shown in **Figure 4**. The PL intensity was normalized to the thickness of each SRN film. The as-deposited SRN film with  $R_N = 80$  (N80) emit a broad PL band with the main peak at 490 nm, as shown in **Figure 4(a)**. As the  $R_N$  value decreases, the PL emission band becomes narrower and shifts to a higher wavelength reaching 590 nm for  $R_N = 5$  (N5, higher silicon excess) but with a lower intensity.

The PL band of the SRN films blueshifts after thermal annealing, particularly for  $R_N = 80$  where the main band shifts from 490 to 420 nm, as shown in **Figure 4(b)**. It is well known that the silicon excess in silicon-rich dielectrics agglomerates forming Si-nps after high temperature annealing, resulting in a redshift of the PL band due to the quantum confinement effects [53–57]. However, contrary to this assumption, in this work, the PL blueshifts after the annealing, as shown in **Figure 4(b)**. This effect has been observed before and it has been ascribed to compositional-dependent changes in the concentration of defect states within the samples [58].



**Figure 4.** (a) Normalized PL spectra of as-deposited, and (b) maximum PL peak position before and after thermally annealed, SRN films with  $R_N = 5$  (N5), 20 (N20), and 80 (N80). Lines are plotted as an eye-guide. From Cabañas-Tay et al. [34].

Therefore, defects are created rather than Si-nps, after thermal annealing. Moreover, due to the low diffusivity of silicon atoms in  $Si_3N_4$ , a high Si content (>52 at.%) is needed to form Si-nps [59, 60]. Because of the low Si content present in the SRN films from this work (<46 at.%), as shown above in the XPS results, these films can be explained as a sub-stoichiometric nitride with structural defects, as reported in the study of Cabañas-Tay et al. [34].

PL bands between 380 and 600 nm (2.0–3.2 eV) have been observed before in SRN films, and they have been ascribed to the radiative recombination of carriers in band tail states, which are related to defect energy levels within the gap of amorphous silicon nitride [61–64]. Therefore, the PL bands emitted by the SRN films from this work can be explained by the excitation of different defects as discussed in the study of Cabañas-Tay et al. [34]. As observed in **Figure 4(a)**, the as-deposited SRN films emit at 590, 580, and 490 nm, whereas the thermally annealed films emit at 580, 540, and 420 nm for  $R_N$  of 5, 20, and 80, respectively. PL emission bands observed at 580–590 nm (~2.1 eV) have been related to electronic transitions from the conduction band minimum (CBM) to K<sup>0</sup> centers located near of the mid-gap, the emission band at 540 nm (~2.3 eV) is related to electronic transition from the K<sup>0</sup> centers to =N<sup>-</sup> centers located near of the valence band maximum (VBM), and the emission band at 420 nm (~2.9 eV) has been related to electronic transitions from the K<sup>0</sup> centers to VBM. Hence, the PL bands observed in this work could have a similar origin, as discussed in the study of Cabañas-Tay et al. [34].

Some studies have shown that electronic transition related with K<sup>0</sup> centers to =Si–O–Si states are observed when oxygen is incorporated in the SRN film [65, 66]. The presence of oxygen in SRN films creates a gap state of Si–O above the VBM, giving rise to the 485 nm (~2.55 eV) emission. The XPS analysis demonstrated that SRN films contain oxygen, being the higher concentration for  $R_N$  = 80 (11.8 at.%). Therefore, the PL emission band at 490 nm (~2.53 eV)

observed in the as-deposited sample with  $R_N = 80$  could be related to the electronic transition from the K<sup>0</sup> centers to =Si–O–Si states.

In summary, the analysis of the PL emission observed in the SRN films before and after thermal annealing indicates that it could be mainly originated from the radiative recombination via luminescent Si dangling bonds, N dangling bonds, and Si–O bonds existing in the silicon nitride matrix.

### 4.2. Silicon-rich oxide (SRO) film

**Figure 5** shows the normalized PL spectra of SRO films before and after thermal annealing. The as-deposited SRO film with Ro = 30 (M30) emits a broad PL band with the main peak at 460 nm. After thermal annealing, the PL emission band becomes narrower and redshifts reaching 690 nm, as shown in **Figure 5(a)**. As the Ro value decreases, the PL emission band becomes wider and shifts to a higher wavelength reaching 660 nm for Ro = 20 (M20, higher silicon excess), as shown in **Figure 5(b)**. After thermal annealing, the PL of M20 redshifts, reaching 740 nm.

The PL emitted by M30 exhibit a significant redshift of the main peak after thermal annealing as shown in **Figure 5(a)**. Nevertheless, when the silicon content is increased (M20), the PL band appears mainly at the red side of the spectrum. The redshift of the main PL peak (after thermal annealing) has been widely observed and ascribed to the agglomeration of silicon excess and a subsequent silicon nanoparticle formation as a result of thermal annealing process [67–71]. Nevertheless, some point defects are also present within the SRO films. It is widely accepted that violet-blue (400–460 nm), green (520), and even red (630) emission bands, obtained from the deconvolution of the PL spectrum [72], can be related with oxygen



**Figure 5.** Normalized PL spectra of SRO monolayer (a) with *R*o = 30 (M30) and (b) *R*o = 20 (M20) before and after thermal annealing. From Palacios-Huerta et al. [46].

defect centers (ODC),  $E'_{\delta}$  (Si $\uparrow$ Si=Si) and NBOHC defects, respectively [67–71]. The  $E'_{\delta}$  center is one of the at least four different E' centers [43], which comprises an unpaired spin delocalized over five silicon atoms and suggest the presence of very small Si-nps in the films.

#### 4.3. Silicon-rich nitride/silicon-rich oxide (SRN/SRO) bilayers

In previous studies, it has been reported that the combination of  $\text{Si}_3\text{N}_4/\text{SRO}$  structure improves luminescent emission properties [73, 74]. Previous studies have also shown that a  $\text{Si}_3\text{N}_4$ -SRO bilayer structure improves the operation of light-emitting devices, such as a reduced leakage current, a reduced electric field on the oxide layer, and results an improvement in efficiency and a longer device life [73–76]. In this chapter, the effect of a SRN film on a SRO film (SRN/SRO bilayers) on their optical properties is analyzed. SRN films, deposited by LPCVD with  $R_N = 80$  (N80), show luminescence in the blue region, while the SRO films emit in the red region, as observed in **Figures 4** and **5**, respectively. Therefore, a combination of SRN with  $R_N = 80$  and SRO films could allow reaching a broad emission spectrum.

**Figure 6** shows the PL spectra of SRN/SRO bilayers with Ro = 20 (B20) and 30 (B30) after thermal annealing. The emission spectra of the SRN (N80) and SRO monolayer films with Ro = 20 (M20) and 30 (M30) are shown as references, respectively. The emission spectra of the samples were normalized to the SRO thickness. Each PL spectrum of SRN and SRO



**Figure 6.** Normalized PL spectra after thermal annealing of (a) SRN/SRO film with  $R_0 = 20$  (B20), SRN film with  $R_N = 80$  (N80), and SRO with  $R_0 = 20$  (M20), and (b) SRN/SRO film with  $R_0 = 30$  (B30), SRN film with  $R_N = 80$  (N80), and SRO with  $R_0 = 30$  (M30).

monolayers was fitted three Gaussian, while four Gaussians were used to fit the PL spectrum of the SRO/SRN bilayer.

For the SRN film (N80), PL bands at about ~420, 505, and 680 nm were identified, which are related to electronic transitions from the K<sup>0</sup> centers to the VBM, K<sup>0</sup> centers to =Si–O–Si, and states of defects nitrogen (N<sup>0</sup><sub>2</sub>), respectively [64]. For the SRO film with Ro = 20 (M20), PL bands at ~420, 735, and 820 nm were identified and have been widely related to weak oxygen bond (WOB) defects, interactions that take place at the Si-np/SiOx interface, and quantum confinement in silicon nanocrystals, respectively [72]. The PL spectrum of the SRN/SRO bilayer with Ro = 20 (B20) exhibits the same PL bands of the SRO films at ~420, ~735, and 812 nm, but with an additional band at ~505 nm (green band), which is related to transitions from K<sup>0</sup> centers to =Si–O–Si featured of the SRN film, as shown in **Figure 6(a)**.

As observed in **Figure 6(b)**, the SRO monolayer with Ro = 30 (M30) exhibits PL bands at ~420, 720, and 805 nm, which have been also linked to WOB defects, interactions that take place at the Si-nps/SiOx interface, and silicon nanocrystals (Si-ncs) through quantum confinement, respectively [72]. The PL spectrum of the SRN/SRO with Ro = 30 (B30) is composed by the same luminescent bands than that observed in the bilayer with Ro = 20 (B20), but with an increased intensity of the bands at ~420 and ~505 nm related with transitions from K<sup>0</sup> centers to VBM (violet band), transitions from K<sup>0</sup> centers to =Si–O–Si (green band), respectively.

In summary, it is observed that the emission intensity in the blue and green bands (~420 and 505 nm) is improved when a SRN/SRO bilayer is formed compared to SRO monolayers, being higher when a Ro = 30 (B30) is used. The XPS results show a diffusion of oxygen from the SRO layer to the SRN layer and nitrogen from the SRN layer to the SRO layer during SRN deposition (onto SRO layer). Therefore, some defects could be passivated into the SRO layer and others generated into the SRN layer like WOB or =Si–O–Si defects.

### 5. Electroluminescence

SRN, SRO, and SRN/SRO films exhibit intense and visible photoluminescence. In this section, luminescent characteristics of the samples are present but under electrical excitation.

### 5.1. Silicon-rich nitride (SRN)-LECs

**Figure 7** shows the EL spectra of the light emitting capacitors using the as-deposited SRN film with  $R_{\rm N}$  = 20 as an active dielectric layer at different injected currents. At forward bias (FB) (positive voltage to FTO contact respect to Si substrate), the SRN-LEC shows a broad spectrum with the maximum emission centered at around 580 nm, and it remains at the same wavelengths for different voltages as shown in **Figure 7(a)**. This luminescence has been reported as a characteristic behavior of defect-related EL [67]. In fact, the luminescent properties of SRN films in this chapter are related to the presence of defects. Therefore, the EL in SRN-LECs may originate from the same electronic transitions of the CBM to K<sup>0</sup> centers,



**Figure 7.** EL spectra of the as-deposited SRN LEC with  $R_N = 20$  for different injected currents at (a) forward bias and (b) reverse bias. In the inset, pictures for the device with 100 mA of injected current for each bias. FTO was used as gate electrode. From Cabañas-Tay et al. [34].

as observed in PL [77, 78]. This behavior is similar for all SRN films in both as-deposited and after thermal annealing, as reported in the study of Cabañas-Tay et al. [34].

On the other hand, at reverse bias (RB), the EL spectrum changes and now three main emission bands are observed at around 600, 680, and 780 nm, as observed in **Figure 7(b)**. Namely, the maximum emission band blueshifts when the polarity is changed from RB to FB. The EL bands at 600, 680, and 780 nm also remains at the same wavelength when the voltage is increased indicating the EL emission is also produced by defects. These EL bands have been ascribed to electronic transitions from the K<sup>0</sup> centers to valence band tail states [58, 77]. The EL emission for all the SRN-LECs in both polarities is through shine dots as shown in the inset of **Figure 7(a)** and **(b)**. The device's area is covered with more shine dots when the current increases. The EL emission composed by shining spots is attributed to the formation of a finite number of preferential conductive paths within the SRN films, which connect the top and bottom electrodes, as discussed in the study of Cabañas-Tay et al. [34].

Negligible spectral shift is observed between the EL at forward bias and PL spectra of SRN films with  $R_{\rm N}$ =5 and 20 before and after thermal annealing, indicating that both PL and EL emissions are originated from the same radiative centers [34]. Nevertheless, a long spectral shift is observed for the SRN films with  $R_{\rm N}$ =80 (before and after thermal annealing). This behavior was explained as a significant smaller electrical pumping of electrons allowing the holes to relax to the lowest defect states before recombining. For the EL at RB, a similar behavior is presented but the electrical pumping of electrons is slightly improved allowing only few of the holes to relax to the lowest defect states before recombining [34].

#### 5.2. Silicon-rich oxide (SRO)-LECs

The presence of defects including the Si-nps, either crystalline or amorphous, and their density and size in silicon-rich dielectric materials affect clearly the current transport, and therefore the EL, as in the SRO case. **Figure 8** shows J-E curves of SRO-LECs with *R*o = 30 and 20; samples were thermally annealed at 1100°C. SRO-LECs are forwardly biased (accumulation mode) considering the substrate as reference.

SRO-LECs with Ro = 30 (M30) films show a high current state (HCS) at low electric fields, then after the applied voltage increased, the current was switched to an LCS, as shown in **Figure 8**. This behavior was observed by our group before in SRO films with Ro = 30, and it was related to the creation and annihilation of preferential conductive paths generated possibly by adjacent stable Si-nps and unstable silicon nanoclusters (Si-ncls) through structural changes and by the possible creation of defects (breaking off Si–Si bonds) [70, 79, 80]. Indeed, a clear correlation between current jumps/drops and EL dots appearing/ disappearing on the LEC surface was observed [70, 79, 80]. This RS behavior is independent on the thermal annealing temperature [46]. Recent studies regarding the same electrical RS behavior in SRO films was observed and related with a conductive filament formed by Si-nps [81–84], which undergoes structural changes through a crystallization and amorphization process of the Si-nps, as discussed in [82]. Such observations are in agreement with the presented asseverations about the behavior of the SRO with Ro = 30 (M30)-based LECs.

Once the current fluctuations disappear, through the electrical annealing, the current behavior stabilized (see I–V curve marked as M30-after in dark blue line, **Figure 8**) and EL on the whole area (WA EL) was observed at higher electric fields. On the other hand, the electrical behavior of SRO-LEC with Ro = 20 (M20) did not show current fluctuations. The latter could



**Figure 8.** J-E curve of SRO-LECs with Ro = 20 (M20) and 30 (M30) thermally annealed at 1100°C. From Palacios-Huerta et al. [46].



**Figure 9.** EL spectra and images of SRO-LECs with (a) *R*o = 20 (M20) and (b) 30 (M30) annealed at 1100°C. Semitransparent n+ polycrystalline silicon was used as gate electrode. From Palacios-Huerta et al. [46].

be related to the presence of well-separated and crystalline silicon nanoparticles, mainly on the density of Si-nps, as shown by transmission electron microscopy (TEM) analysis [46]. Therefore, a uniform network of conductive paths becomes possible as the Si-nps density increases, allowing a uniform charge flow through the whole capacitor area.

**Figure 9** shows the EL spectra from the SRO-LECs. Blue EL in the whole area was observed with Ro = 30 (M30) only after the current drop. The main EL peak remains at 468 nm when the electric field increases, as observed in **Figure 9(a)**. When the silicon content is increased (Ro = 20), the LECs emit a broad EL spectrum in the red region, as observed in **Figure 9(b)**. The EL peak is placed at 710 nm. This value is blueshifted with respect to the value found in PL at 740 nm. Differences in the PL and EL peak wavelength have been related to the transmittance of the top electrode, which influences the real light emission from the active layer [85]. Indeed, the process of carrier injection and the several transport mechanism taking place through the active luminescent materials produces a complicated understanding of the EL process, as compared to the PL one [86]. Nevertheless, both EL and PL spectra appear in the same red region, which could indicate that the emission is originated by the same radiative centers. Images of LECs are shown as insets in **Figure 9**. As can be observed, the LEC devices emit an intense whole area EL.

#### 5.3. Silicon-rich nitride/silicon-rich oxide (SRN/SRO)-LECs

**Figure 10** shows the J(E) characteristic and the EL spectra of SRN/SRO-LECs at forward bias considering the substrate like reference. SRN/SRO-LECs show a high current state (HCS) at low electric fields, and then after the applied voltage increases, the current is switched to a LCS, as shown in **Figure 10(a)**. The resistive switching in the B20 bilayer occurs at a lower electric field (~2 MV/cm) compared to the B30 bilayer. When SRN/SRO-LECs, both with *R*o = 20



**Figure 10.** (a) J-E curve of SRN/SRO-LECs with *R*o = 20 (B20) and 30 (B30) and EL spectra of SRN/SRO LECs at forward bias for different currents applied with (b) *R*o = 20 (B20) in the HCS, (c) *R*o = 20 (B20) in the LCS, and (d) *R*o = 30 (B30) in the LCS. Inset of each graph are shown photographs of the LEC. ITO was used as gate contact.

(B20) and Ro = 30 (B30), are in the HCS (<20V), a broad spectrum with the maximum emission centered at about 730 nm and a shoulder at ~900 nm is observed. The 730 nm emission increases as the voltage also increases, as observed in **Figure 10(b)**. Nevertheless, first the shoulder at 900 nm increases its intensity when the voltage is between -17 and -18 V and then remains at the same value for larger voltages. As we can see in **Figure 10(b)**, the EL is observed as randomly scattered luminescent spots (mainly red) over the entire LEC surface. The spectrum exhibits that the EL emission follows the PL, showing that the emission is caused by the same luminescent centers.

After current switching, in the LCS (>20V), the EL spectra changes, as observed in **Figure 10(c)**. Also, a narrow (width of  $7 \pm 0.6$  nm) and highly intense UV EL peaks appear at ~250, 270, 285, 305, 325, and 415 nm. A narrow blue EL peak with the highest intensity is observed at ~450 nm. All EL peaks remain at the same wavelength as the current increases. The red EL band observed in the HCS is still present at the LCS, but with a slight redshift and with a low intensity. The emission intensity of the blue EL is about 20 times higher than the red emission. The EL emission is observed as luminescent blue-violet dots randomly distributed on the surface of the contact, as shown in the insets of **Figure 10(c)**.

**Figure 10(d)** shows the EL emission of the SRN/SRO-LEC with Ro = 30 (B30) in the range of LCS (>35V). The same narrow and highly intense UV EL peaks at ~250, 270, 285, 305, 325, 415, and 450 nm with an average width of  $7 \pm 0.6$  nm are also observed in these devices. Nevertheless, these LECs do not show EL emission in the red-near infrared region.

Four narrow UV EL peaks have been reported at 293.78, 316.10, 403.07, and 444.82 nm, with an average width at half peak of 4 nm, in ITO/Y<sub>2</sub>O<sub>3</sub>/Ag EL devices. This emission was attributed to the characteristic radiation of indium ions [87]. However, as the best of our knowledge, narrow UV EL peaks in silicon rich dielectric materials have not been reported before. The most intense EL peaks emitted by the present SRN/SRO-LECS (~305, 325, 415, and 450 nm) are very similar to those obtained in reference [87], but displaced ~10 nm toward longer wavelengths. Thus, they could have a similar origin; however, a deeper analysis of these narrow emission bands needs to be done.

### 6. Conduction mechanisms

The origin of the EL emission from the SRO, SRN, and SRN/SRO-LECs can be determined by identifying the charge transport mechanism that takes place within the different silicon-rich dielectric materials as discussed below.

### 6.1. Silicon-rich nitride (SRN)-LECs

**Figure 11** shows the J(E) dependency for the SRN-LECs with  $R_N = 5$ , 20, and 80. The conduction mechanism of the carriers in the SRN films was studied by analyzing the different reported mechanism including P-F conduction [88, 89], TAT [90], and space charge limited current (SCLC) [91]. As observed in **Figure 11**, the TAT mechanism is dominating the carrier transport at both forward and reverse bias, and it can be ascribed to the defect states generated during the SRN deposition. A model based on the trap-assisted tunneling carrier transport is shown with the EL radiative recombination process in the reference [34].

### 6.2. Silicon-rich oxide (SRO)-LECs

**Figure 12** shows the experimental J-E data from SRO-LECs with Ro = 20 and 30 fitted to the carrier transport mechanisms including Poole-Frenkel (P-F) conduction [88, 89] and trapassisted tunneling (TAT) [90]. As can be observed, the TAT conduction mechanism predominates in the SRO-LEC with Ro = 30 (M30).

The trap energy was estimated to be  $\phi_t \cong 1.75 \text{ eV}$ , if a uniformly distributed trap concentration is assumed. It is worth to mention that these trap energy levels correlates well with the one obtained for Er-doped SRO/SiO<sub>2</sub> superlattice (2.1 eV) and Er-doped SiO<sub>2</sub> (1.9 eV), which are assumed to be deep traps inherent to the SiO<sub>2</sub> [92] possibly neutral oxygen vacancy [93]. Mehonic et al. have reported that the barrier height of the trap is modified by the concentration of oxygen vacancies in the connecting tissue; lower barrier heights correspond to a higher concentration of oxygen vacancies, which defines the device resistance [84]. Moreover, it has been reported that a poor quality of nanocrystalline composite material (formed by Si



**Figure 11.** J-E characteristic and conduction mechanism fitting of SRN-LECs with  $R_N = 5$  (N5), 20 (N20), and 80 (N80) after thermal annealing at: (a) reverse bias (RB) and (b) forward bias (FB). Symbols are the experimental current data, black line represents the TAT conduction, green dash line SCLC, and violet dash dot line the Poole-Frenkel conduction. From Cabañas-Tay et al. [34].



**Figure 12.** J-E characteristic and conduction mechanism fitting of SRO-LECs with *R*o = 30 (M30) and *R*o = 20 (M20). Red line represents the TAT conduction and blue line the Poole-Frenkel conduction. Symbols are the experimental data. From Palacios-Huerta et al. [46].

nanocrystals and SiOx) yields parasitic current paths, while low density of nanocrystals turns difficult the direct charge injection into the nanocrystals [94]. A similar behavior is obtained for SRO-LECs with *R*o = 30 (M30), where the annihilation of conductive paths, created by adjacent stable Si-nps and unstable silicon nanoclusters (Si-ncls), produce structural changes including the possible creation of defects. Then, it can be assumed that M30 films exhibit a poor quality of the silicon oxide resulting from the phase separation and the low silicon content. This behavior is well correlated to the blue EL emission observed in the M30-based LECs.

On other hand, the P-F conduction fits well the charge transport in SRO-LECs with Ro = 20 (M20), as observed in **Figure 12**. A relative permittivity ( $\varepsilon_r$ ) value of 9.16 was obtained from the P-F fit for M20, which is closer to the relative permittivity of silicon ( $\varepsilon_{si} = 11.9$ ), similar to other reports [92]. Relatively high permittivity values are a good indication of the large amount of silicon present as Si-nps within the M20 films. Moreover, the Si-nps size calculated by using the P-F estimate relative permittivity are very close to that obtained through the TEM analysis, as reported in [46].

#### 6.3. Silicon-rich nitride/silicon-rich oxide (SRN/SRO)-LECs

When the SRN/SRO bilayer structure is used as active layer in LECs, several charge transport mechanisms are present, as observed in **Figure 13**. For the SRN/SRO-LEC with *R*o = 20 (B20), and in the HCS where red EL emission is observed, the hopping tunneling is the predominant mechanism (pink line), as shown in **Figure 13**(a).

The average distance between traps (*a*) was estimated to be about 0.25 nm. When the current switches to the low conduction state, the tunneling mechanism changes, then the P-F conduction



**Figure 13.** J-E characteristic and conduction mechanism fitting of SRN/SRO-LECs with (a) *R*o = 20 (B20) and (b) *R*o = 30 (B30) thermally annealed at 1100°C. Pink line represents the hopping conduction, orange line the TAT conduction, and blue lines the Poole-Frenkel conduction. Symbols are the experimental data.

dominates (dark blue line). The dielectric constant ( $\varepsilon_r$ ) obtained by the P-F fit was about 7.96. However, when the deep UV EL emission starts, the conduction transport is dominated by two different mechanisms: TAT (orange line) and P-F (blue line). In this region, the traps depth obtained by the TAT fit was about 2.06 eV, and the dielectric constant ( $\varepsilon_r$ ) obtained by the P-F fit was about 7.81.

The SRN/SRO-LECs with Ro = 30 (B30) behaves similar to the B20 bilayer, as observed in **Figure 13(b)**. In the high conduction state, the dominant mechanism is hopping tunneling (pink line), where the average distance between traps is about 0.15 nm, indicating a higher density of traps compared to the B20 sample (0.25 nm). When the current drops to the low conduction state, the dominating mechanism changes to the P-F conduction (blue line). The dielectric constant ( $\varepsilon_r$ ) was estimated at 5.98, indicating a lower content of silicon related to B20 sample (8.61), as expected. However, when the EL emission starts, the conduction begins to be dominated by TAT (orange line). The traps depth was estimated at about 2.15 eV, very similar to that obtained in B20 sample (2.06) and very close to the depth at which are the centers K<sup>0</sup> (~2.1 eV) [64].

In summary, the conduction mechanism in the region of low electric fields at HCS is hopping for SRN/SRO bilayers. It was observed that when the nitrogen content in the SRO layer increases, the average distance between defects also increases either due to the defects passivation or due to the reduction of the Si-nps size. In the region of high electric fields at low conduction, both TAT and P-F mechanisms take place simultaneously. The trap depth, obtained by the TAT fit, was ~2.1 eV from the minimum of the conduction band for both bilayers (B20 and B30), which relates to the location of the centers K<sup>0</sup> in SRN film.

### 7. Conclusions

The compositional, structural, optical, and electro-optical properties of SRO, SRN films as well as combination of SRN/SRO bilayers deposited by LPCVD were studied. The EL of the SRN-LECs showed a broad emission spectrum where the maximum peak blueshifts when the polarity changed from reverse to forward bias. The EL spectrum was nearly similar to that of PL when LECs were forwardly biased and the silicon excess was increased. Analyzing the current-voltage characteristics, it was found that TAT was the main carrier transport mechanism in SRN films in both biases, where typical EL was observed. In SRO films, it was demonstrated that the silicon content affects the luminescence centers density obtaining the EL emission at lower electric field as the silicon excess increased. SRO with Ro = 30films exhibited the presence of structural defects like NOV, NBOHC, and E'<sub>8</sub> (Si↑Si=Si) centers, which produced blue to red PL band after thermal annealing. SRO with Ro=20 films showed mainly red PL produced by the presence of well-defined Si-ncs. It was found that SRO-LECs with Ro = 30 shows an intense blue EL, after a resistance switching behavior was reached. The RS was originated by the annihilation of preferential conductive paths created by Si-nps and silicon nanoclusters through structural changes that result in a long spectral blueshift (~222 nm) between the PL and EL. On the other hand, the RS was not observed in SRO-LECs with *Ro* = 20 due to the well-defined Si-nps with 2.7 nm size. These well-defined Si-nps produce a better charge injection through multiple conductive paths enhancing an intense red EL, which is related mainly to defects at the Si-nps/SiOx interface. The transport mechanism was affected by the composition and structure of the films, being TAT and PF the dominant transport mechanism in SRO with *Ro* = 30 and 20, respectively. When the SRN was deposited onto the SRO to form the bilayer, compositional and electro-optical changes were observed. The SRN/SRO layers showed a broad PL emission respect to a simple SRO or SRN layer. In addition, multiple defects or centers were responsible of PL emission. The same luminescent centers in SRN/SRO were present in SRO and SRN monolayers. SRN/SRO-LECs showed an intense UV EL at high electric fields related to indium ions radiation. The conduction mechanism in SRN/SRO bilayers was found to be dominated by Pool-Frenkel and TAT mechanisms at LCS, and hopping at HCS.

### Acknowledgements

This work has been partially supported by the project CONACyT-180992. The authors acknowledge technicians Pablo Alarcon, Armando Hernández, and Victor Aca from INAOE and Luis Gerardo Silva from CIMAV.

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