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Modelling and Optimizing Operating Conditions of Heat Exchanger with Finned Elliptical Tubes

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1. Introduction

High performance heat exchangers are widely used in many industries such as petrochemical industry, automotive industry, energetic industry and many others. The name of this kind of heat exchangers is in consistency with their structure. Large heat transfer area leads to high thermal efficiency of the device. Its working principle is to cool rapidly large amount of gaseous or liquid medium. Because of its compact size, it is possible to use it for easy installation in various systems, like the heating, drying, air conditioning and the other systems.

High thermal performance of this type of heat exchangers (Nagatha, 2004; Ibrahim & Gomaa, 2009) is achieved due to large packed rectangular fins on tube surface, as well as the shape of the tubes. The elliptical shape ensures better distribution of velocity along the tube circumference during the flow, as well as lower pressure drop in comparison to the circular shape (Matos et al, 2004; Jang & Yang, 1998; Li et al, 2006; Nishiyama et al, 1988; Khan et al, 2004; Saboya & Saboya, 1981; Saboya & Sparrow, 1974; Bordallo & Saboya, 1999). Consequently, the coefficient of heat transfer in the elliptical shape is larger when compared to the circular shape (Idem et al, 1990; Rocha et al, 1997; Taler, 2009). It brings about that the heat exchange process is more efficient in elliptical than in circular tubes.

To use all the aforementioned advantages, it is necessary to ensure uniform distribution of velocity in all the tubes. Failing that, it leads to large differences in mean temperature in the tubes. Consequently, excessive thermal stress occurs, that may cause the heat exchanger to break down. Small volume of collectors of the heat exchangers implicates possibility of improper flow condition inside the tubes, causing unsuitable inner distribution of thermal and mechanical loads. Damages to a tube, Fig. 1, as well as a tube bundle prone to buckling, Fig. 2, prove, that such a failure is possible (Łopata & Ocłon, 2010; Ocłon & Łopata, 2011).

Compressible stresses are undesirable (Łopata & Ocłon, 2010; Łopata & Ocłon, 2011) in this type of construction, because the allowable compressible load (allowable buckling stress) for a slender tube is significantly lower than the safe tensile load.

The simplified method to assess failure of a high performance heat exchanger with finned elliptical tubes is presented in the paper. The diagram, shown on Fig. 3, presents the method. The geometrical representation of construction – a CAD model is prepared at first. Next, on the basis of prepared geometry, numerical representation of the model is created. Afterwards, the flow computation - CFD (Ferziger & Peric 1999; Anderson, 1995; Gresho &

Sani, 2000; Chung, 2010) is carried out using the commercial tool ANSYS CFX (ANSYS, 2008). It is desirable to obtain the uniform mass flow rates in all the tubes of the heat exchanger, because it ensures the appropriate heat transfer conditions for the device. When the mean velocities in tubes differ significantly (Łopata & Ocłon, 2010), new modified collectors are applied to improve flow distribution inside the device (Ocłon & Łopata, 2011). If the liquid distribution among all the tubes ensures safe heat transfer conditions, the CFD computation is stopped. The results obtained with the CFD analysis as the bulk temperatures T_{bulk} and the heat transfer coefficients h are transferred next to a code, that calculates temperature distribution inside the elliptical tube, using the Control Volume Method (Chung, 2010; Taler & Duda, 2006) as the mixed type boundary conditions (Incopera, 2001). The obtained temperature and pressure values are used as thermal loads for structural analysis, using the ANSYS Structural (ANSYS ADPL, 2009) software.



Fig. 1. View of damaged tube



Fig. 2. Buckled tubes of a high performance heat exchanger



Fig. 3. Calculation methodology

This commercial tool computes displacement and stresses for the numerical model of the heat exchanger using the Finite Element Method (Reddy, 1993; Logan, 2011; Smith & Griffiths, 2004; Zienkiewicz & Taylor, 2000; Hutton, 2003). Structural analysis allows finding the zones and locations where allowable stresses are exceeded. When it happens, the computational flow loops back to the CFD analysis and the other geometrical parameters of the collector are set. The loop continues until allowable stresses in the construction are exceeded. If the maximum stress is lower than the critical stress, the construction is safe to use, and the computations stop.

The methodology presented in the paper may help to assess structural safety. The computations described in this chapter were successfully applied in analysis of two failures of a high performance heat exchanger that happened in an industrial company in Poland.

2. Hydraulic and thermal performance of heat exchanger

The proper operating conditions of the high performance heat exchanger improve its thermal efficiency and lead to significantly smaller pressure exerted on tube material.

The operating conditions depend strongly on flow distribution inside the tubes of the high performance heat exchanger. The main task of hydraulic computation is to confirm, that liquid flowing inside the tubes is distributed properly in all the tubes. It is ideally, when the mass flow rate \dot{m} in each tube equals to the mass flow at the inlet \dot{m}_{in} divided by the number of tubes *n* in a single pass of the heat exchanger. Nevertheless, it is impossible in fact to obtain the uniform flow distribution in all the tubes. The improper flow distribution has an influence on heat transfer process in each pass of the heat exchanger, causing high temperature differences in the adjacent tubes that can cause the stress limit of the tube material to be exceeded.

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329



The heat exchanger, for which the computations are carried out, is presented schematically on Fig. 4.

Fig. 4. Scheme of a two pass high performance heat exchanger

The operating conditions of the heat exchanger, as presented on Fig. 4, are listed in Tab. 1:

Name	Symbol	Value, Unit
Mass flow rate of water	\dot{m}_{wat}	19 kg /s, 14 kg /s, 9.5 kg/s
Velocity of combustion gas	W _{cg-in}	4 m/s
Number of tubes	п	42
Inlet and outlet temperature of water	$T_{wat-in}/T_{wat-out}$ (max value)	110/130 °C
Inlet and outlet temperature of combustion gas	$T_{cg-inlet}/T_{cg-outlet}$ (min value)	450/350 °C
Operating pressure for water	p _{op}	6 bar

Table 1. Operating conditions for the analysed high performance heat exchanger

Working principle of the analysed heat exchanger is to cool rapidly combustion gas. Heat is transferred from hot gas to water which flows into the device through the inlet nozzle pipe at the mean temperature $T_{wat-in} = 110$ °C. Water fills in the upper part of the collector. Next, it is distributed in the 1st pass to n = 42 tubes of the heat exchanger, where it flows down to the lower collector. Then, it turns back and flows through the other part of the tubes (2nd pass of the heat exchanger) into the upper collector and finally leaves the device through the outlet nozzle pipe with the mean temperature $T_{wat-out}$, that equals to no more than 130 °C. The tubes are arranged into two rows. Combustion gas, with the mean inlet temperature

 $T_{cg-inlet} = 450$ °C flows perpendicularly to the tubes, exchanging heat with the first row of the tubes and subsequently with the second one. Finally, it leaves the device cooled to no less than $T_{cg-outlet} = 350$ °C. Later, the gas is cooled to 100 °C in other heat exchangers (not shown here) and released to the atmosphere. The presented heat exchanger is a part of the installation where the largest thermal stresses may occur, so its failure is most probable to happen.

2.1 Numerical model of flow around finned elliptical tube

Analysis of flow dynamic and heat transfer is divided into two parts. First, heat transfer from combustion gas to the tube is modelled, next the flow of water inside the heat exchanger is analysed, knowing the thermal boundary conditions - heat transfer coefficient and bulk temperature distribution on the outer surface of the tube. Discretization of flow domain (lower case) and boundary condition types (uppercase) for analysis of the flow around the finned elliptical tube are presented on Fig. 5.



Fig. 5. Numerical model of flow around finned elliptical tube

The combustion gas is modelled as an ideal gas with density $\rho_{cg} = 3 \text{ kg/m}^3$ and dynamic viscosity $\eta_{cg} = 27 \cdot 10^{-6} \text{ Pa/s}$. Mass, momentum and energy equation (Bird, 2001) are used in computation. In order to model the turbulence, the SST (shear-stress transport) transport equations (Egorov & Menter, 2008) are used in the CFD analysis. The mixed boundary condition simulates the heat transfer from water to a tube. The bulk temperature of water T_{wat-b} is assumed to be equal $T_{wat-out}$ and the heat transfer coefficient h_{wat} is calculated, knowing the mean velocity in a tube, which is calculated according to the following formula:

$$\overline{w_{wat}} = \frac{\dot{m}_{wat}}{n \cdot A_{tube} \cdot \rho_{wat}},\tag{1}$$

where A_{tube} is equal to the area of flow through an elliptical duct and ρ_{wat} is the density of water calculated for T_{wat-b} .

The Nusselt Number of water Nu_{wat} is calculated according to (Gnielinski, 1976):

$$Nu_{wat} = \frac{\frac{1}{8} f \cdot (Re_{wat} - 1000) \cdot Pr_{wat}}{1 + 12.7 \cdot (\frac{1}{2} \cdot f)^{\frac{1}{2}} \cdot (Pr_{wat}^{\frac{2}{3}} - 1)}.$$
 (2)

It was stated in the paper (Łopata & Ocłon, 2010), that velocity of water in this type of a heat exchanger can be significantly different for tubes in the 1st and in the 2nd pass of the high performance heat exchanger. The zones and locations in the tube were indicated where mean velocity is over 20 times lower than in the other tubes. The thermal load is significantly larger in those zones than in the ones with proper velocity distribution. Moreover, low velocities of liquid favour deposition of boiler scale (Najibi et al., 2007; Quan et al., 1997; Nebot et al., 2007; Freeborn & Lewis, 1962), especially when water is not prepared properly that causes the flow resistance growth. Consequently, water velocity decreases even more. With regard to the above, it is not the average velocity of water, but the $w_{wat} = 1/2 \cdot \overline{w_{wat}}$ that is used to calculate the Reynold's number of water - Re_{wat}. It simulates the real operating conditions of a device, when scale starts to deposit inside the tube and thermal load increases significantly.

The Prandtl number of water Pr_{wat} is obtained for the specific heat capacity c_{p-wat} , dynamic viscosity η_{wat} and thermal conductivity k_{wat} of water that on the turn are measured at bulk temperature equal to T_{wat-b} :

$$\Pr_{wat} = \frac{c_{p-wat} \cdot \eta_{wat}}{k_{wat}} \tag{3}$$

The Darcy friction factor *f* is calculated according to formula (White, 2003):

$$f = (0.79ln \cdot \text{Re}_{wat} - 1.64)^{-2}.$$
 (4)

Finally h_{wat} is computed as follows:

$$h_{wat} = \frac{\mathrm{Nu}_{wat} \cdot k_{wat}}{d_h}.$$
(5)

The equivalent diameter d_h of an elliptical tube is computed according to the following dependence:

$$d_h = \frac{4 \cdot A_{tube}}{P_{tube}},\tag{6}$$

where P_{tube} is equal to the perimeter of an elliptical duct. The values of h_{wat} , obtained from eq. (5) and T_{wat-b} for all the computational cases are listed in Table 2.

<i>ṁ_{wat}</i> [kg/s]	$h_{wat} \left[\frac{W}{m^2 \cdot K}\right]$	T_{wat-b} [°C]
19	3630	130
14	2720	130
9.5	1815	130

Table 2. Values of T_{wat-b} and h_{wat} used in analysis of flow around finned elliptical tube

The cross – section of the tube, which fractured during the device operation, is presented on Fig. 6. It is possible to observe deposition of scale on the inner surface of the tube. For computational cases, two thicknesses of scale $t_s = 1 \text{ mm}$ and 1.5 mm are analysed. Thermal conductivity of scale depends on quantity of calcium carbonate (Najibi et al. 1997). The thermal conductivity of boiler scale is assumed to be equal to 0.1 W/(m K) for the present

computation that models the location of the fracture. Scale deposition is modelled only for the tubes in the middle of computational domain, because the existence of scale in neighbouring tubes wasn't noticed during the site inspection.

It is possible to observe, on the outer wall of the elliptical tube, the layer of fouling that also influences heat transfer inside the device, worsening considerably thermal contact between the fin and the tube.



Fig. 6. Picture of damaged tube (cross - section)

This layer is modelled in ANSYS CFX as thermal resistance on the interface between the fin band and the tube. Thickness of material is assumed to be 0.1 mm and its thermal conductivity equals to 1 W/(m K).

Heat flux density \dot{q} is computed on the interface between the tube (Fig. 7) and the fin band so that the interface area is divided into n_{int} = 40 intervals along the *y* direction.



Fig. 7. Division of interface area into intervals

For the n_{el} number of elements that belong to each of the *y* intervals the area averaged heat flux is computed as follows:

$$\dot{q}_{int} = \frac{\sum_{j=1}^{n_{el}} q_j \cdot A_j}{\sum_{j=1}^{n_{el}} A_j}.$$
(7)

The area averaged wall temperature is calculated according to the formula:



$$T_{wall} = \frac{\sum_{j=1}^{n_{el}} T_j \cdot A_j}{\sum_{j=1}^{n_{el}} A_j}.$$
 (8)

The mass flow averaged temperature is computed in the middle of each *y* interval:

$$T_{cg-bulk} = \frac{\sum_{j=1}^{n_{el}} T_j \cdot m_j}{\sum_{j=1}^{n_{el}} m_j}.$$
(9)

Knowing values of the area averaged heat flux density for each interval \dot{q}_{int} , the area averaged wall temperature for intervals T_{wall} and the bulk temperature computed in the middle of interval $T_{cg-bulk}$, the heat transfer coefficient from combustion gas to the tube is determined according to the formula:

$$h_{cg}(y) = \frac{\dot{q}_{int}}{(T_{cg-bulk} - T_{wall})}$$
(10)

The sample results presented below were obtained for $\dot{m}_{wat} = 19$ kg/s. The bulk temperature and heat transfer coefficient distribution for all the \dot{m}_{wat} values listed in Table 2 are used as the thermal boundary condition for analysis of flow distribution inside the high performance heat exchanger.

On Figure 8 temperature and velocity distribution is presented when combustion gas flows around the finned elliptical tube.



Fig. 8. Temperature and velocity distribution for $t_s = 0$ mm

It is possible to observe, that the stagnation zones occur under the 1st row as well as above and under the 2nd row of tubes. Therefore the fins don't operate properly, and the largest amount of heat is transferred only to the 1st row of tubes. It may lead to high temperature differences on the circumference of the tubes and consequently may cause significant bending stresses to the tube.

The temperature distribution in the fin and in gas for $t_s = 1$ mm is presented on Fig 9. The boiler scale deposition significantly increases temperatures of the tube and the fin. Consequently thermal loading of the tubes in this case is considerably greater than when $t_s = 0$ mm.



Fig. 9. Temperature distribution for $t_s = 1 \text{ mm}$

Increase in scale thickness from 1 mm to 1.5 mm (see Fig. 10) causes the temperature of the tube, the fin and combustion gas to grow up. It is evident that boiler scale reduces heat flux from combustion gas to water, causing temperature of the fins and the tube to increase. This effect is more pronounceable with increase in boiler scale thickness.



Fig. 10. Temperature distribution for $t_s = 1.5$ mm



On Fig. 11 the h_{cg} and $T_{cg-bulk}$ distribution along y direction for t_s is presented

Fig. 11. Heat transfer coefficient and bulk temperature distribution for combustion gas - $t_s = 0$ mm, $t_s = 1$ mm and $t_s = 1.5$ mm

Heat transfer coefficient from combustion gas to the tube is larger when boiler scale thickness increases. Rapid growth of $T_{cg-bulk}$ combined with isolative behaviour of scale, causes high temperature increase of the tube. Therefore, operating condition of the device worsens significantly, because the normal stresses acting on the tube grow up with increase in temperature, too.

2.2 Flow distribution inside collectors and tubes of high performance heat exchanger

Analysis of flow distribution inside the high performance heat exchanger is carried out using the mass, momentum and energy balance equations (Bird, 2001; Incopera & De Witt, 2001). The turbulence model SST is applied to flow behaviour simulation. The numerical model of the flow is presented on Fig. 12.



Fig. 12. Numerical model used for analysis of flow distribution inside heat exchanger

The boundary layer is modelled using the hexahedral elements and the flow core is discretized using the tetrahedral elements. The boundary conditions set for \dot{m}_{wat} = 19 kg/s is presented on Fig. 13.



Fig. 13. Boundary condition for analysis of water flow - \dot{m}_{wat} = 19 kg/s

As mentioned in section 1, it is difficult to ensure the uniform flow distribution to all the tubes, due to limited space in the collectors. The zones of the tubes where the velocity is significantly lower than the $\overline{w_{wat}}$ are marked on red, see Fig. 14.



Fig. 14. Velocity distribution in high performance heat exchanger - \dot{m}_{wat} = 19 kg/s

These zones are formed because pressure changes occur (see Fig. 15) in the direction parallel to the perforated bottom of the collector.



Fig. 15. Pressure and temperature distribution inside high performance heat exchanger - \dot{m}_{wat} = 19 kg/s

The pressure is lower in the tubes with velocity bellow $\overline{w_{wat}}$ than in the others. Therefore, the driving force of the flow falls down and liquid doesn't come into this segment of the tubes. E.g., in the first pass of the heat exchanger, liquid flows into the tubes of the heat exchanger, those are located directly under the inlet nozzle pipe. Velocity of liquid in these tubes is larger than $\overline{w_{wat}}$. Amount of liquid flowing into the other tubes in this pass is much lower so are velocities. For the second pass of the heat exchanger, the component of velocity vector, parallel to the perforated bottom of the collector, is significantly larger than the other components. Therefore, due to the large inertia forces, water encounters difficulties when changing the flow direction by 90 degrees especially when entering the first tubes of the 2nd pass of the heat exchanger. The mean velocity in tube *w* (calculated as area averaged velocity) is the smallest there. It is even 20 times lower than the mean value (see Fig. 16).

Low flow velocities inside the tubes cause temperature of water to increase significantly. The zones and the locations where saturation temperature is exceeded are marked in red on Fig 15., the vapour bubbles starts to form there. When saturated steam flows inside the tube, the value h_{wat} increases rapidly. As a result, heat transfer conditions in the device change considerably and differ from the design predictions.

The ratio of average velocity in tube w, computed from numerical simulation, to mean velocity of water $\overline{w_{wat}}$, presented on Fig. 16, shows that difficulties with improper flow distribution inside the heat exchanger exist in all the computational cases. For $\dot{m}_{wat} = 14$

338

kg/s in the 2nd pass of the heat exchanger, *w* can be even 50 times smaller than $\overline{w_{wat}}$, that can lead to fracture of the first tube near the axis of symmetry (2nd pass), because temperature increases and stress level exceeds the allowable limits. Therefore the construction of the collector should be modified to improve the ratios $w/\overline{w_{wat}}$, for the 1st and the 2nd pass of the high performance heat exchanger.



Fig. 16. Ratios of mean velocity in tube *w* to estimated velocity $\overline{w_{wat}}$ for $\dot{m}_{wat} = 9.5$ kg/s, 14 kg/s and 19 kg/s

The modified version of the collector is presented in the next subchapter.

2.3 Analysis of flow distribution inside modified collector of high performance heat exchanger

As mentioned before, difficulties in ensuring uniform flow velocity inside all the tubes of the high performance heat exchangers have forced us to find a new shape of the collector that improves flow distribution inside the tubes of the 1st and the 2nd pass. The shape of the collector, presented on Fig. 17, is being proposed. The greatest advantages of using it are its simplicity and low costs of manufacturing.



Fig. 17. Traditional and modified shape of collector

The wedge is being welded to the wall of the collector pipe. It increases flow space to enable the more uniform flow distribution inside all the tubes. It is the influence on flow distribution of three parameters (Fig. 18): the h_w , which is the height of wedge, the tube height h_t and x that is the distance between the side of collector and the axis of nozzle pipe, that are investigated.



Fig. 18. Geometrical parameters of modified collector

The presented results from analysis are carried out for $\dot{m}_{wat} = 19$ kg/s. The ratio of maximal velocity in the tube for the 1st and the 2nd pass of the heat exchanger to mean velocity of water in the tube $w_{max}/\overline{w_{wat}}$, as well as the ratio of minimal velocity for the 1st and the 2nd pass of the heat exchanger to mean velocity of water in the tube $w_{min}/\overline{w_{wat}}$ can be observed on Fig. 19 for $h_w = 40$ mm.



Fig. 19. Highest ratios of maximal and minimal velocity in tube w_{max} and w_{min} to $\overline{w_{wat}}$ for $\dot{m}_{wat} = 19 \text{ kg/s}$ and $h_w = 40 \text{ mm}$.

The presented results show that it is difficult to obtain geometrical parameters of the collector for that the ratios $w_{max}/\overline{w_{wat}}$ and $w_{min}/\overline{w_{wat}}$ are optimal for the 1st and the 2nd pass of the heat exchanger. For instance, the parameters x = 500 mm, $h_t = 38$ mm and $h_w = 40$ mm ensure proper flow distribution to the 2nd pass of the heat exchanger, but in the 1st pass the ratio $w_{max}/\overline{w_{wat}}$ has the highest value. For the presented case it is not as dangerous as low values of $w_{min}/\overline{w_{wat}}$ for the construction; but when the function of the heat exchanger is to cool water, the large value of ratio $w_{max}/\overline{w_{wat}}$ is not desired. The worst results are for x = 200 mm, $h_t = 60$ mm and $h_w = 40$ mm, when the ratio $w_{min}/\overline{w_{wat}}$ equals 0.05 for the 1st pass of the heat exchanger. In the case of heating the medium that flows inside the heat exchanger it can bring about significant compressible stresses acting on the tube.

The results obtained for $h_w = 50$ mm doesn't differ significantly from the previously presented. However, it is possible to observe on Fig. 20 that for the 1st pass of the heat exchanger the ratio $w_{min}/\overline{w_{wat}}$ is lower for x = 500 mm than for the same values of x and h_t if $h_w = 40$ mm. This means, that the increase in h_w from 40 mm to 50 mm doesn't improve the flow condition inside the device considerably.



Fig. 20. Highest ratios of maximal and minimal velocity in tube w_{max} and w_{min} to estimated velocity $\overline{w_{wat}}$ for $\dot{m}_{wat} = 19 \text{ kg/s}$ and $h_w = 50 \text{ mm}$.

The results for $h_w = 60$ mm are shown on the Fig. 21. Extreme positioning of the nozzle pipe: x = 200 mm and x = 500 mm causes improper flow distribution for the 1st pass of the heat exchanger.



Fig. 21. Highest ratios of maximal and minimal velocity in tube w_{max} and w_{min} to estimated velocity $\overline{w_{wat}}$ for $\dot{m}_{wat} = 19 \text{ kg/s}$ and $h_w = 60 \text{ mm}$.

The last set of results, Fig. 22, is presented for $h_w = 70$ mm. In contrary to the results presented before, increasing h_w with increasing h_t up to 55 mm, improves $w_{min}/\overline{w_{wat}}$ in the 1st pass of the heat exchanger for x = 200 mm. The largest improvement is in ratio $w_{min}/\overline{w_{wat}}$ especially for the 1st pass, that is 0.4 for x = 200 mm.



Fig. 22. Highest ratios of maximal and minimal velocity in tube w_{max} and w_{min} to estimated velocity $\overline{w_{wat}}$ for $\dot{m}_{wat} = 19 \text{ kg/s}$ and $h_w = 70 \text{ mm}$.

All of the presented results show that it is difficult to find optimal solution for geometry presented on Fig.18. However, it is possible to reach compromise between large values of $w_{max}/\overline{w_{wat}}$ and low values of $w_{min}/\overline{w_{wat}}$, e.g. for the set of parameters x = 300 mm, $h_w = 40$ mm and $h_t = 55$ mm. In the 1st pass of the heat exchanger, the $w_{max}/\overline{w_{wat}}$ is slightly larger than for the traditional collector. However, the biggest advantage of using the modified shape of the collector, with the mentioned parameters is that the $w_{min}/\overline{w_{wat}}$ is similar for both the 1st

and the 2nd pass of the heat exchanger and doesn't exceed 0.6. Therefore the differences in the wall temperature will be not as large as for the traditional shape of the collector.

3. Temperature distribution inside elliptical tube

Numerical code in c++ has been written for the thermal analysis, in order to obtain temperature distribution inside the elliptical tube. The boundary condition for this analysis is taken from the CFD computation (water flow distribution inside the tubes). On the outer surface of the elliptical tube, the h_{cg} and $T_{cg-bulk}$ values are applied. On the inner surface of the elliptical tube, the h_{wat} values obtained for the CFX calculated mass flow rate and mass flow averaged temperature in each tube are taken into computation. If it is necessary, the code can incorporate the boiler scale layer into calculation, as well as fouling resistance on the outer surface of the tube.

The results obtained from thermal analysis for all the tubes are then used as the thermal loading in structural analysis.

Temperature distribution inside the elliptical tube is obtained in the Finite Volume Method via the Finite Element Method (Taler & Duda, 2006; Chang, 2010). The two-dimension steady state heat transfer is assumed; therefore the isoparametric linear quadrilateral elements (Chang, 2010, Lewis, 2004) are applied. The isoparametric element creating the control volume is presented on Fig. 23.



Fig. 23. Finite Volume Method via Finite Elements

The real coordinates (x,y) are mapped into the natural coordinates (ξ,η) . Therefore, the quadrilateral element can be treated as the rectangular element. The lines connecting the middles of the rectangle sides and the nodes of the quad element create the sub-control volumes, as presented on Fig. 23. Four sub-control volumes create one large control volume. Each sub-control volume is also treated as the quadrilateral isoperimetric element; therefore, the shape functions are applied to approximate temperature distribution and coordinates inside the quad element. For the isoperimetric elements, the temperature distribution inside the element domain is estimated, using the following shape function (Taler & Duda, 2006):

$$\Phi_{1}^{e} = \frac{1}{4} \cdot (1 - \xi) \cdot (1 - \eta)
\Phi_{2}^{e} = \frac{1}{4} \cdot (1 + \xi) \cdot (1 - \eta)
\Phi_{3}^{e} = \frac{1}{4} \cdot (1 + \xi) \cdot (1 + \eta)
\Phi_{4}^{e} = \frac{1}{4} \cdot (1 - \xi) \cdot (1 + \eta)$$
(11)

Temperature inside the element is obtained by multiplying the nodal temperature T_i and the function Φ_i^e that corresponds to the node shape:

$$T(\xi,\eta) = \sum_{i=1}^{4} T_i \cdot \Phi_i^e \tag{12}$$

The x and y coordinates inside the element domain are interpolated in a similar manner as temperature:

$$x = \sum_{\substack{i=1\\4}}^{4} x_i \cdot \Phi_i^e$$

$$y = \sum_{\substack{i=1\\i=1}}^{4} y_i \cdot \Phi_i^e$$
(13)

As mentioned before, the control volume consists of four sub-control volumes A, B, C, D. The heat flux balance for the control volume can be written in the form:

$$k_{A} \cdot \sum_{CS2,3} \left(\frac{dT}{dx} \cdot n_{x} + \frac{dT}{dy} \cdot n_{y} \right) \cdot \overline{\Delta\Gamma} + k_{B} \cdot \sum_{CS3,4} \left(\frac{dT}{dx} \cdot n_{x} + \frac{dT}{dy} \cdot n_{y} \right) \cdot \overline{\Delta\Gamma} + y = \sum_{i=1}^{4} y_{i} \cdot \Phi_{i}^{e}$$

$$k_{c} \cdot \sum_{CS4,1} \left(\frac{dT}{dx} \cdot n_{x} + \frac{dT}{dy} \cdot n_{y} \right) \cdot \overline{\Delta\Gamma} + k_{D} \cdot \sum_{CS1,2} \left(\frac{dT}{dx} \cdot n_{x} + \frac{dT}{dy} \cdot n_{y} \right) \cdot \overline{\Delta\Gamma} = 0$$

$$(14)$$

Where n_x and n_y are the components of the normal vector \vec{n} and $\vec{\Delta \Gamma}$ is the boundary vector. Heat flows between control volumes through control surfaces Cs (two for each sub-control volume). In numerical code, the energy equation is applied in the following manner for sub-control volumes A, B, C and D:

• For sub-control volume A:

$$Q_A = k_A \cdot \left(\frac{dT}{dx} \cdot \cos\theta_2 + \frac{dT}{dy} \cdot \sin\theta_2\right) \cdot \overrightarrow{\Delta\Gamma_2} + k_A \cdot \left(\frac{dT}{dx} \cdot \cos\theta_3 + \frac{dT}{dy} \cdot \sin\theta_3\right) \cdot \overrightarrow{\Delta\Gamma_3}$$
(15)

• For sub-control volume B:

$$Q_B = k_B \cdot \left(\frac{dT}{dx} \cdot \cos\theta_3 + \frac{dT}{dy} \cdot \sin\theta_3\right) \cdot \overrightarrow{\Delta\Gamma}_3 + k_B \cdot \left(\frac{dT}{dx} \cdot \cos\theta_4 + \frac{dT}{dy} \cdot \sin\theta_4\right) \cdot \overrightarrow{\Delta\Gamma}_4 \tag{16}$$

• For sub-control volume C:

$$Q_C = k_C \cdot \left(\frac{dT}{dx} \cdot \cos \theta_4 + \frac{dT}{dy} \cdot \sin \theta_4\right) \cdot \overline{\Delta \Gamma_4} + k_C \cdot \left(\frac{dT}{dx} \cdot \cos \theta_1 + \frac{dT}{dy} \cdot \sin \theta_1\right) \cdot \overline{\Delta \Gamma_1}$$
(17)

• For sub-control volume D:

$$Q_D = k_D \cdot \left(\frac{dT}{dx} \cdot \cos\theta_1 + \frac{dT}{dy} \cdot \sin\theta_1\right) \cdot \overrightarrow{\Delta\Gamma}_4 + k_D \cdot \left(\frac{dT}{dx} \cdot \cos\theta_2 + \frac{dT}{dy} \cdot \sin\theta_2\right) \cdot \overrightarrow{\Delta\Gamma}_2$$
(18)

Sum of fluxes to control volume equals to 0:

$$Q_A + Q_B + Q_C + Q_D = 0 (19)$$

Derivatives of temperature are obtained from the following equation:

$$\frac{dT}{dx} = \frac{1}{|J|} \cdot \left(\frac{dT}{d\xi} \cdot \frac{dy}{d\eta} - \frac{dT}{d\eta} \cdot \frac{dy}{d\xi}\right)$$

$$\frac{dT}{dy} = \frac{1}{|J|} \cdot \left(\frac{dT}{d\eta} \cdot \frac{dx}{d\xi} - \frac{dT}{d\xi} \cdot \frac{dx}{d\eta}\right)$$
(20)

And the Jacobian of isoparimetric transformation is obtained according to the formula:

$$|J| = \frac{dx}{d\xi} \cdot \frac{dy}{d\eta} - \frac{dy}{d\xi} \cdot \frac{dx}{d\eta}$$
(21)

 $T(\xi, \eta)$ is differentiated with respect to ξ and η , and the values of $\frac{dT}{d\xi}$ and $\frac{dT}{d\eta}$ are computed in the middle of each control surface (see Tab. 3):

Control Surface	ξ	η
Cs ₁	0	-0.5
Cs ₂	0.5	0
Cs ₃	0	0.5
Cs ₄	-0.5	0

Table 3. Values of isoparametric parameters ξ , η for control surfaces

The system of heat balance equations (see eq. 19) is being solved for each control volume. Assuming the mixed type boundary conditions on the inner and the outer wall of tube, the temperature distribution inside the elliptical tube is obtained. The sample results for the tube from the 1st row are presented on the Fig. 24 and Fig 25. Applied boundary conditions on the inner wall of the tube are: $h_{wat} = 5000 \text{ W/} (\text{m}^2 \text{ K})$, $T_{wat-bulk} = 120 \text{ °C}$. On the outer wall, h_{cg} and $T_{cg-bulk}$ - obtained for $t_s = 0 \text{ mm}$ and adequately $t_s = 1 \text{ mm}$ is assumed – compare with Fig. 11.



Fig. 24. Temperature distribution for $h_{wat} = 5000 \text{ W}/(\text{m}^2 \text{ K})$, $T_{wat-bulk} = 120 \text{ }^{\circ}\text{C}$, $t_s = 0 \text{ mm}$



Fig. 25. Temperature distribution for h_{wat} =5000 W/(m² K), $T_{wat-bulk}$ =120 °C, t_s = 1 mm It is possible to observe, on Fig. 25, that when boiler scale exists inside the tube, e.g. Fig. 26, temperature increases significantly.



Fig. 26. Boiler scale observed in tubes of heat exchanger

The large increase of temperature brings about considerable thermal stresses inside all the tubes that may lead to fracture of the tube; see Fig. 1 & Fig. 2.

4. Analysis of stress state inside construction of high performance heat exchanger

As mentioned before, structural analysis is carried out on basis of the Finite Element Method using the commercial code ANSYS. The manner in which the construction is supported is illustrated on Fig. 27. As it can be observed, the heat exchanger consists of two collectors with perforated bottoms to which all the tubes are welded. A shell connects the collectors and flat bars welded to the shell stiffen its structure. The frame increases the stiffness of the perforated bottom.

The other components of construction (a conffuser and a diffuser), not shown on Fig. 27 are welded to the shell. These parts constrain structure movement in the *y* direction. The heat exchanger is also clamped at the distance z = 525 mm from the upper perforated bottom. The isoparametric linear shell element, SHELL 181 (ANSYS, 2009), used for the structural analysis is presented on Fig. 28. The CAD model of the heat exchanger is divided into a set



of finite elements. Each part of the construction consists of elements with different thickness, t_{shell} , (see Tab. 4). Thickness of heat exchanger components are listed on Table 4.

Fig. 27. Construction and support conditions for heat exchanger

Component	t _{shell} [mm]
tube	2
flat bar	15
perforated bottom	12
nozzle pipe	2.5
frame	20
collector	6
shell	8

Table 4. Thickness of heat exchanger components

The midsurface is geometrical representation of the element. The maximal stress value is computed at the outer fibbers of the element.



Fig. 28. Isoparametric shell element for structural calculations

The set of equation being solved in structural analysis (ANSYS, 2009) is presented below:

$$[K_e] \cdot \{u\} - \{F_e^{th}\} = [M_e] \cdot \{\ddot{u}\} + \{F_e^{pr}\},\tag{22}$$

where the element stiffness matrix $[K_e]$ is obtained from the following formula:

$$[K_e] = \int_{\mathcal{U}} [B]^T \cdot [D] \cdot [B] dV, \tag{23}$$

where [B] is the strain displacement matrix based on the element shape functions, $\{u\}$ is the nodal displacement vector and [D] is the elasticity matrix. The thermal load vector is calculated on the basis of the equation:

$$\left|F_{e}^{th}\right| = \int_{V} \left[B\right]^{T} \cdot \left[D\right] \cdot \{\varepsilon_{th}\} dV, \qquad (24)$$

The thermal strain vector $\{\varepsilon_{th}\}$ is calculated on the basis of the following dependence:

$$\{\varepsilon_e^{th}\} = \alpha \cdot (T - T_{ref}),\tag{25}$$

where α is the thermal expansion coefficient. Nodal temperatures *T* are obtained from the thermal calculations, section 3. T_{ref} is the reference temperature assumed as 22 °C. The element mass matrix $[M_e]$ is calculated on the basis of the following equation:

$$[M_e] = \rho \cdot \int_{V} [N]^T \cdot [N] \cdot [D] dV, \qquad (26)$$

where [*N*] is the element shape function matrix and ρ is material density that the element consists of. For the computation steel P355NH is used. Yield strength Y_s of this material equals 355 MPa and the density $\rho = 7850 \text{ kg/m}^3$. The element pressure vector $\{F_e^{pr}\}$ is calculated according to the formula:

$$\left\{F_e^{pr}\right\} = \int_A \left[N_n\right]^T \cdot \{P\} dA,\tag{27}$$

where $[N_n]$ is the matrix of shape functions for normal motions at the surface and $\{P\}$ is the nodal pressure vector assumed equal to p_{op} – Tab. 2.

The material properties used in the structural analysis are presented in Tab. 5.

Name	Symbol	Value
Young's modulus	E[MPa]	$2.1 \cdot 10^{5}$
Tangent modulus	T[MPa]	470
Thermal expansion coeff.	α [1/K]	$1 \cdot 10^{-5}$
Poisson's ratio	ν	0.3

Table 5. Material properties of P355NH steel

4.1 Analysis of stress state for construction of high performance heat exchanger with traditional collectors

Results of numerical computation for nodal displacements, compressible stresses and equivalent stress acc. to von Misses are presented in this subchapter. Computations are carried out using the model of bilinear isotropic hardening (ANSYS, 2009; Chakrabarty, 2006), because computations in elastic range show that the allowable stresses inside the construction are exceeded.

Thermal loadings for all the tubes are obtained using a program described in section 3 (computational case with $t_s = 0$ mm). Because combustion gas flow is on the outer surface of the elliptical tubes, where heat flux density is the highest – avg. 80000 W/m². Heat flux density on the outer wall of perforated bottoms, collectors and nozzle pipes is low - avg. 5000 W/m². Therefore, temperature of walls for these parts is assumed to be equal to near wall temperature of liquid obtained from the flow distribution analysis – subsection 2.2. Thermal loading of construction is presented on Fig. 29. The contour map of nodal displacements {*u*} is presented on Fig. 30.



Fig. 29. Thermal loading for construction of high performance heat exchanger – traditional collector



Fig. 30. Nodal displacement – high performance heat exchanger with traditional collectors

Displacements of the elliptical tube, that is the nearest to the axis of symmetry of the high performance heat exchanger are the largest at all. Comparing the results presented on Fig. 30 to Fig. 2, it is evident that reasons for the heat exchanger to break down are improper heat transfer conditions inside the construction. In the tube that broke down the temperature was significantly higher than in the others.

The next two figures – Fig. 31 & Fig. 32 present the contour of equivalent stress distribution σ_e acc. to von Mises (Bower, 2009), and the map of compressible stresses σ_3 inside the construction. The von Mises stress is calculated according to the formula:

$$\sigma_e = \frac{1}{2} \cdot \sqrt{(\sigma_1 - \sigma_2)^2 + (\sigma_2 - \sigma_3)^2 + (\sigma_3 - \sigma_1)^2},$$
(28)

where σ_1 , σ_2 and σ_3 are principal stresses. It is possible to observe that in the shell and in the tube that was broken the stress level is the largest. According to the adopted plasticity model, maximum compressible stress in the broken tube equals to 205 MPa.



Fig. 31. Equivalent stress acc. to von Mises - heat exchanger with traditional collector



Fig. 32. Compressible stresses - heat exchanger with traditional collector

Because the tube is slender and thin, it is prone to buckling, that is caused by compressible loads. Buckling computations (Brush & Almroth, 1975) are carried out in order to understand the manner how the tube was destroyed. The assumed model of buckling is presented on Fig 33. The left end is clamped as opposed to the right one that is supported in vertical direction.



Fig. 33. Buckling model of elliptical tube

The effective length of the tube equals:

$$L_e = K \cdot L, \tag{29}$$

where K equals to the tube effective length factor and L is the unsupported length of the tube.

The slenderness ratio *s* for the tube presented on Fig. 33 equals:

$$s = \frac{L_e}{\sqrt{\frac{I_{min}}{A}}},\tag{30}$$

where I_{min} equals to minimal area moment of inertia for the tube cross section and A equals to the cross – sectional area of the tube. The critical slenderness ratio s_{cr} is calculated according to the formula:

$$s_{cr} = \pi \cdot \sqrt{\frac{E}{Y_s}}.$$
(31)

The following values, Tab. 6, are used in calculations:

Name	Symbol	Value
Tube effective length factor	K	0.5
Length of unsupported tube	L[m]	2.1
Minimal area moment of inertia	$I_{min}[m^4]$	$4.85 \cdot 10^{-9}$
Cross – sectional area of tube	$A[m^2]$	$1.7 \cdot 10^{-4}$
Poisson's ratio	ν	0.3
Critical slenderness ratio	S _{cr}	73.1
Slenderness ratio	S	168

Table 6. Parameters taken into buckling calculations

Since *s* is larger than s_{cr} , see Tab.6, buckling occurs in the elastic range. The following formula allows calculating the critical compressible stress:

$$\sigma_{crit} = \frac{\pi^2 \cdot E}{s^2} \,.$$

The allowable stresses obtained on the basis of procedure (28 – 32) are less than σ_{crit} = 66 MPa. Comparing σ_{crit} to the maximal value from the contour map of compressible stresses in the construction, it is evident, that the critical stress is exceeded in the tube nearest to the axis of symmetry. This explains the manner of tube deformation, after the damage, compare with Fig 2.

4.2 Analysis of stress state for construction of high performance heat exchanger with modified collectors

For the construction of the high performance heat exchanger, computations are carried out in the elastic range, because the σ_e for this computations is lower than Y_s . On Fig. 34 & Fig. 35 thermal loading and the nodal displacements of construction are presented.

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(32)



Fig. 34. Thermal loading for construction of high performance heat exchanger – modified collectors



Fig. 35. Nodal displacement - high performance heat exchanger with modified collectors

It is possible to observe that, for the proposed construction of the collector, displacements are significantly lower than for the typical one. The maximum displacement is over 20 times lower – compare Fig. 35 and Fig. 30.

For the modified construction, distribution of equivalent stress acc. to von Mises is presented on Fig. 36. The σ_e value obtained from computations does not exceed Y_s , therefore the construction is safe to use. Moreover the compressible stresses σ_3 in the tubes, Fig. 37, do not exceed 22 MPa and are significantly lower than σ_{crit} .

The modified construction of the collectors improves flow distribution inside the tubes through the heat exchanger. Consequently, thermal load acting on the tubes is lower than in the case of the traditional construction of the collectors. Comparing the maps of nodal displacements $\{u\}$, effective stresses σ_e according to von Mises and compressible stresses σ_3 for the traditional and the modified construction of the collectors, it is evident that the proposed construction carries thermal loads better than the one actually available in manufacturing.



Fig. 36. Equivalent stress acc. to von Mises - heat exchanger with modified collector



Fig. 37. Compressible stresses - heat exchanger with modified collector

5. Conclusions

Improper liquid distribution in the tubes of the presented heat exchanger causes the mass flow rates of medium flowing inside all the tubes being far from uniform. This implicates that heat transfer conditions inside the tubes differ significantly. Consequently, the temperatures as well as the thermal stresses inside all the tubes are different, often exceeding the allowable stress limits.

Reasons of high performance heat exchanger failure, on the basis of the real model applied in the industry, were investigated. The zones and the locations where the allowable stresses are exceeded were presented in the paper. Presented analysis may be applied not only in analysing the stress state in the high performance heat exchanger, but also in accessing safety of boilers, pressure vessels, shell and tube heat exchangers and many other devices applied in the thermal power engineering.

Furthermore, the new and easy in manufacturing model of collectors is presented. On the basis of numerical computation it was proven that application of the new collectors ensures proper flow distribution inside all the tubes that consequently reduces significantly thermal stresses inside the construction.

With continuous increase in computational power it will be possible to investigate in more detail the flow behaviour and heat transfer conditions inside the high performance heat exchanger, including e.g. the phase change effect. If liquid evaporates, the heat transfer coefficient increases significantly and due to the local pressure drop, the cavitation may occur inside the tube that leads to erosion of the tube wall. Consequently, material of the tube is damaged and less resistible on acting stresses. Therefore, the real values of σ_e and σ_3 may be larger than the presented in the paper. This effect should be encountered into calculations in the future.

Wall temperature increases significantly in the tubes where the mass flow rates are low. If water is not prepared properly, boiler scale formation is more intensive than in the other tubes. Influence of boiler scale thickness and its thermal conductivity on the stress state in the high performance heat exchanger should be investigated in the future.

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Fluid Dynamics, Computational Modeling and Applications Edited by Dr. L. Hector Juarez

ISBN 978-953-51-0052-2 Hard cover, 660 pages **Publisher** InTech **Published online** 24, February, 2012 **Published in print edition** February, 2012

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Stanisław Łopata and Paweł Ocłoń (2012). Modelling and Optimizing Operating Conditions of Heat Exchanger with Finned Elliptical Tubes, Fluid Dynamics, Computational Modeling and Applications, Dr. L. Hector Juarez (Ed.), ISBN: 978-953-51-0052-2, InTech, Available from: http://www.intechopen.com/books/fluid-dynamics-computational-modeling-and-applications/modelling-and-optimizing-operating-conditions-of-heat-exchanger-with-finned-elliptical-tubes

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