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Dissolution Trapping of Carbon Dioxide in Reservoir Formation Brine – A Carbon Storage Mechanism

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1. Introduction

Carbon Capture and Storage (CCS) is a method to reduce anthropogenic greenhouse gas emissions thereby mitigating global warming. In CCS, carbon dioxide (CO₂) is captured from fossil fuel-fired power plants or other large point-source emitters, purified, compressed and injected deep underground into subsurface formations at depths of or greater than 800m. At such depths CO₂ is in a supercritical (sc) state increasing storage capacity (IPCC 2005).

In CCS, there are four main mechanisms which keep the buoyant CO₂ underground:

1. Structural/stratigraphic trapping – here an impermeable caprock prevents the CO₂ from flowing upwards,
2. Capillary trapping, where micrometer-sized disconnected CO₂ bubbles are formed and held in place by local capillary forces in the rock pore-network,
3. Dissolution trapping, where CO₂ dissolves in the formation brine and sinks in the reservoir as the CO₂-enriched brine has an increased density,
4. Mineral trapping, where the dissolved CO₂ reacts with the formation brine, forms carbonic acid which dissociates generating protons, HCO₃⁻ and CO₃²⁻ ions; these species subsequently react with the formation brine and/or host rock to form solid minerals which trap the CO₂ very safely.

The focus of this text is on dissolution trapping; how much CO₂ dissolves under which geothermal conditions and what happens to the CO₂-enriched brine, which is slightly denser than the original formation brine, in the formation.

Important open questions in this context are: How fast are these mass transfer processes in real geological porous media under realistic CCS conditions? Are there means of accelerating CO₂ dissolution? How do separate gas and/or oil phases (oil and/or gas reservoirs) in the reservoir affect CO₂ dissolution processes and reservoir fluid dynamics? How does the pressure drop due to CO₂ dissolution affect injectivity and storage capacity of CO₂?

2. Geological background of dissolution trapping

The International Panel on Climate Change (IPCC) (2005) has suggested several possible geological storage media, including deep saline aquifers, oil or gas reservoirs and unmineable

coal seams. In case of CO₂ storage in coal, a benefit is that additional methane is produced which is adsorbed on the coal surface and displaced by CO₂ (so-called enhanced coal-bed methane (ECBM) production). However, CO₂ injection leads to the highly detrimental effect of coal swelling which strongly deteriorates injectivity as observed from laboratory and pilot field studies (Reeves and Oudinot, 2005). This text focuses on aquifers and oil/gas reservoirs and will not discuss ECBM any further as low permeability and swelling characteristics limit the scale of exploitation of coalbeds as potential CO₂ storage sinks.

In terms of CO₂ storage, deep saline aquifers – too saline for drinking water or agricultural usage – are most promising, because they are geographically widespread and have large potential storage capacities. Published storage capacity estimates especially for aquifers vary widely based on the assumptions made. This is an active area of research with the objective to provide accurate basic information so that effective CCS schemes can be planned in order to store the large quantities of anthropogenic CO₂ emitted (circa 30 Gt CO₂/a, IPCC 2007).

To focus on dissolution trapping, the topic of this chapter, the main problem associated with it is addressed straight away: it is the slow speed of CO₂ dissolution and the two-phase (CO₂ and brine) reservoir flow dynamics – as long as the CO₂ is in a separate supercritical state it tends to flow upwards because of buoyancy forces, and it can potentially leak to the surface. Mass transfer of CO₂ from the supercritical phase into the aqueous phase is the time-determining step in dissolution trapping which therefore also determines leakage risk. In fact CO₂ is only stored safely once it is dissolved in the aqueous phase (or precipitated as a solid). Hence the study of CO₂ dissolution is an essential aspect of CCS risk assessment. Mass transfer and solubilities of CO₂ into brine are functions of pressure, temperature, salinity, local CO₂ concentration and subsequent chemical reactions (formation and dissociation of carbonic acid and following rock dissolution/precipitation). Moreover interfacial areas scCO₂-brine play a vital role in the mass transfer kinetics, and they are closely related to the two-phase flow dynamics in the reservoir. All these aspects will be discussed in this chapter. In addition several reservoir scale computer simulations will be presented which analyze fluid flow and CO₂ storage in CCS schemes.

In this context it is worth noting that CO₂ is a naturally abundant species in the subsurface. Rumble et al. (1982) suggested two possible chemical reactions between calcite and quartz which formed this naturally occurring CO₂ over geological times. A result of this is that CO₂ content in oil or gas reservoirs can be very high. In gas reservoirs CO₂ content can reach concentrations larger than 90 mol% and in oil reservoirs CO₂ content can be as high as 70-80 mol% (Badessich et al. 2005). As an example Ballentine et al. (2001) state that the CO₂ concentration in gas fields in Texas varies from 3% to 97% depending on the geographical location.

In summary dissolution trapping is a feasible mechanism to store large quantities of CO₂, and if a route could be found to quickly dissolve scCO₂ into brine CO₂ emissions could be dramatically, rapidly and economically reduced this way, maybe even solving the climate change problem caused by CO₂ gas emitted from large point-sources. However, although CO₂ contributes the largest chunk to greenhouse gas emissions, other gases such as CH₄, CO, N₂O, halogenated carbons, etc., also need to be eliminated to completely stop global warming. One route for disposing these gases may also be dissolution into formation brines.

3. Reservoir fluid dynamics

In actual ongoing CCS projects large quantities of CO₂ are injected deep underground. The largest injection time for a pure CCS project has been achieved in Norway in the Sleipner

project, where 1 Mt CO₂/a are injected into the Utsira sandstone formation in the Norwegian sector of the North Sea at 800m depth (Iglauer 2011). This project started in 1996, and reservoir CO₂ monitors confirm reservoir simulations which predict that the CO₂ rises upwards and accumulates beneath the caprock (Hesse et al. 2008).

CO₂ from this rising CO₂-plume dissolves in brine as it migrates upwards (Pruess and Garcia 2002, Bachu and Adams 2003). The CO₂-enriched brine has a slightly higher density than the original brine (Ennis-King and Paterson 2005, Moortgat et al. 2011). This leads to gravitational flow instabilities in the reservoir (Riaz et al. 2006, Pau et al. 2010), and it is believed that the CO₂-rich brine sinks in the reservoir over hundreds to millions of years (Bachu 2000, Ennis-King and Paterson 2005, Lindeberg and Wessel-Berg 1997) in the form of thick and thin fingers (cp. Figures 1 and 2), however this is an active area of research and it has been suggested that this mechanism is considerably faster (Moortgat et al. 2011).

Again, this storage mechanism is very safe, but if the dissolution process is a very slow process then that means that the leakage risk is high in the short term (= initial several hundreds of years) since the CO₂ may escape before it can dissolve.

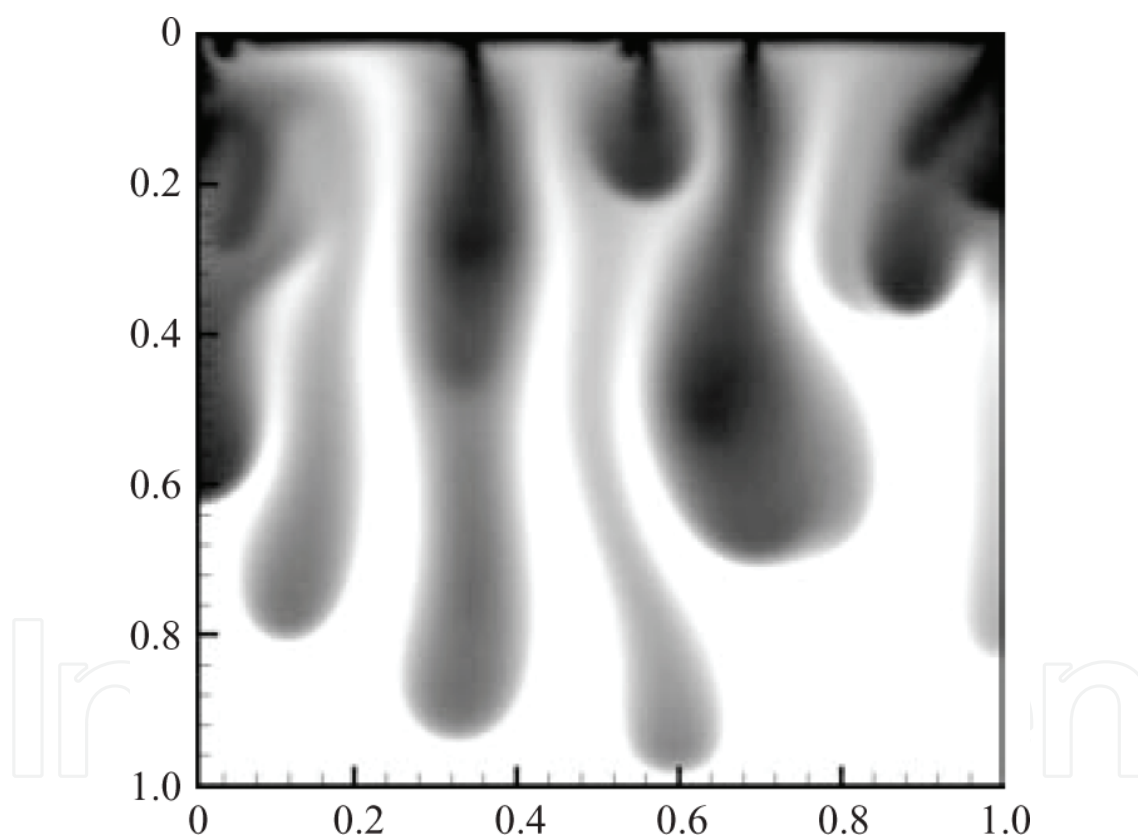


Fig. 1. CO₂-enriched brine sinks in a normalized simulated reservoir over hundreds to thousands of years (from Riaz et al. 2006 with permission from Cambridge University Press). The CO₂-concentration contours are shown in greyscales. The x- and y-axis are normalized lengths, the corresponding absolute values are in the kilometer range

4. Thermodynamics of CO₂ dissolution into formation brine

It has been reported that 0.9-3.6 mol% of CO₂ can be dissolved in brine, depending on pressure, temperature and brine composition (Rumpf et al. 1994, Koschel et al. 2006, Bando et al.

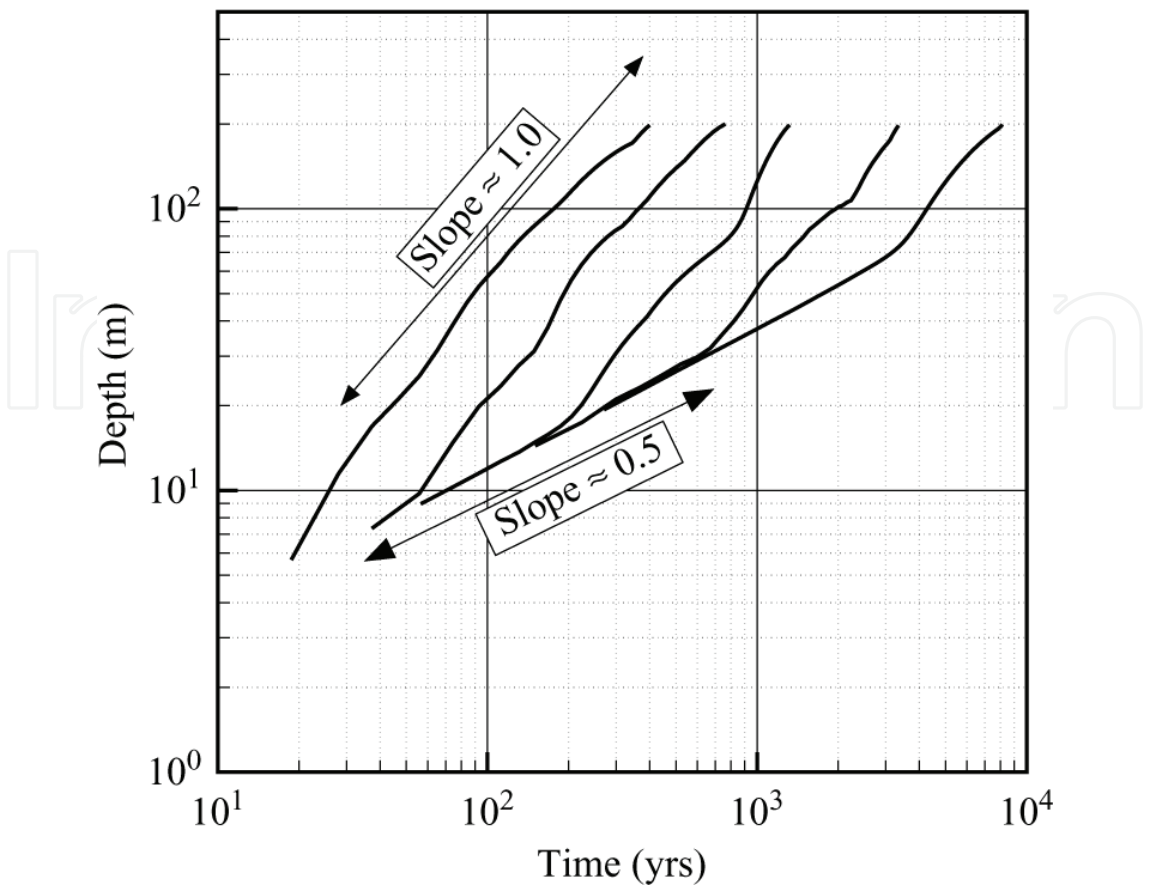
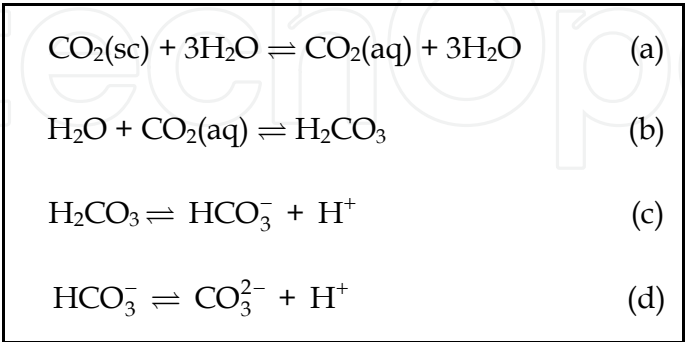


Fig. 2. The advance of the fastest finger front is shown for different permeabilities (represented by different lines) (from Riaz et al. 2006 with permission from Cambridge University Press)

2003, Kiepe et al. 2002). Before analyzing these relationships in more depth, it should be pointed out that CO₂ and brine are a reactive system, CO₂ reacts with water to form carbonic acid which subsequently dissociates (scheme 1) through a proton-relay mechanism that is catalyzed by several water molecules (Adamczyk et al. 2009) lowering the pH value of the brine.



Scheme 1. Formation and dissociation of carbonic acid. Reaction scheme (a) assumes that scCO₂ is dissolved in an analogous way to gaseous CO₂ (Adamczyk et al. 2009)

Adamczyk et al. (2009) studied these reactions at atmospheric pressure and found that the slowest step in scheme 1 is the forward reaction of (b), the hydration of CO₂(aq) resulting in

H₂CO₃. The dehydration of H₂CO₃ is also relatively slow, with a dehydration rate constant of $k_{de} = 18 \text{ s}^{-1}$ (Pocker and Bjorkquist 1977). The deprotonation rate in scheme (c) is $k_{off} = 10^7 \text{ s}^{-1}$ (Pocker and Bjorkquist 1977) and the associated pK_a value is 3.45. Note that this pK_a value published by Adamczyk et al. (2009) is considerably different from the normally assumed pK_a = 6.35 for the CO₂(aq)/H₂O system at atmospheric pressure conditions. The protonation rate of scheme (c) then results in $k_{on} = k_{off}/K_a$. The K value at 50.66 MPa for the reaction in scheme (d) is 5.13×10^{-11} (Hirai et al. 1997).

In addition, one consequence of an increase in CO₂ solubility with increasing pressure (or decreasing temperature or salinity) is that the aqueous phase increasingly acidifies because more CO₂ is present in the aqueous phase and reaction (b) is shifted to the right side according to Le-Chatelier's principle. It can therefore be expected that the pK_a value of scheme (c) drops further at increased CO₂ pressure.

In laboratory measurements pH values between 3.2-3.6 were observed within a temperature range between 300-343 K, a pressure range between 4-11 MPa and a salinity range 1-4 M NaCl solutions (Schaeff and McGrail 2004). In siliclastic and carbonate gas fields however pH values between 5-5.8 have been observed (Gilfillan et al. 2009); the discrepancy between lab and field data is most likely caused by complex geochemical buffering reactions, e.g. with carbonate host rock or carbonate based cements.

The increased proton concentration in the brine has significant implications for geochemical reactions (Stumm and Morgan 1996, Gauss 2010) generally leading to more rock dissolution and higher dissolution rates. When the pH value has increased again to sufficiently high levels CO₂ can be trapped as a solid phase – so-called mineral trapping (IPCC 2005, Gauss 2010). This could in principle also be engineered in the future although the physical and chemical phenomena associated with this process are highly complex and coupled.

Such reactions bring a range of problems and advantages with them:

- It is possible that too much rock is dissolved and high permeability channels are formed; this is especially a problem in carbonates (Egermann et al. 2005, Luquot and Gouze 2009). Injected CO₂ will preferentially flow through such channels, which are also termed "wormholes". This reduces reservoir sweep efficiency which again decreases capillary trapping as only low initial CO₂ saturations are achieved. Low initial CO₂ saturations however result in low residual CO₂ saturations (Pentland 2010 and 2011a,b; Al-Mansoori et al. 2010; Iglauer 2009). In addition such high permeability flow paths increase the risk of CO₂ leakage, especially if caprock material is affected.
- Should so much host rock be dissolved that the mechanical rock integrity is affected, then this can result in wellbore instability or even landmass subsidence.
- In case of precipitation of solid minerals due to geochemical reactions (when the pH value has increased again) rock permeability can be significantly reduced, e.g. by blockage of small pore throats (which determine the permeability value). This can result in serious injectivity problems, e.g. injection rates may have to be reduced dramatically which may render CCS schemes ineffective.
- Rock dissolution increases permeability and enhances injectivity rendering CCS schemes more economical.
- Precipitation of CO₂ in solid minerals (after chemical reactions) is the safest form of CO₂ storage in CCS as the CO₂ cannot escape to the surface anymore. This trapping mechanism is believed to take between thousands to billions of years (IPCC 2005, Xu et al. 2003).

In light of the new results published by Adamczyk et al. (2009) it is important to note that carbonic acid has a considerable acidity as it acts like a carboxylic acid on nanosecond timescales; this may have significant implications for geochemical reactions, rock surface alterations and associated possible rock wettability changes. Rock wettability strongly influences multi-phase fluid dynamics and capillary trapping.

On an important side issue these chemical reactions also happen in the oceans when CO₂ gas in the atmosphere dissolves in seawater thereby reducing its pH value. With the increasing CO₂ concentration in the atmosphere (from 190 ppm in 1750 to 380 ppm in 2005, IPCC 2005) more carbonic acid is formed in the oceans and the seawater pH value decreases with possible massive effects on sea life, starting with the sensitive but all important sea plankton. Therefore disposing anthropogenic CO₂ by dissolving it into the ocean seems to be a risky enterprise, as the pH value would drop further and locally reach substantially lower numbers.

4.1 Effect of pressure on CO₂ solubility in brine

CO₂ solubility (mole fraction of CO₂ per mass unit of brine) in formation brine is a strong function of pressure as shown in Figure 3. The data curve (open diamonds) in Figure 3 was computed with Duan and Sun (2003) and Duan et al. (2006)'s online CO₂ solubility calculator. The temperature was held constant at 323 K and brine salinity was 1 mol NaCl/kg brine. CO₂ solubility rapidly increases when pressure is raised from 0.1 MPa to 10 MPa, then the increase flattens out although a slight solubility increase follows. Three experimentally measured points at CCS pressure conditions are also added to the graph.

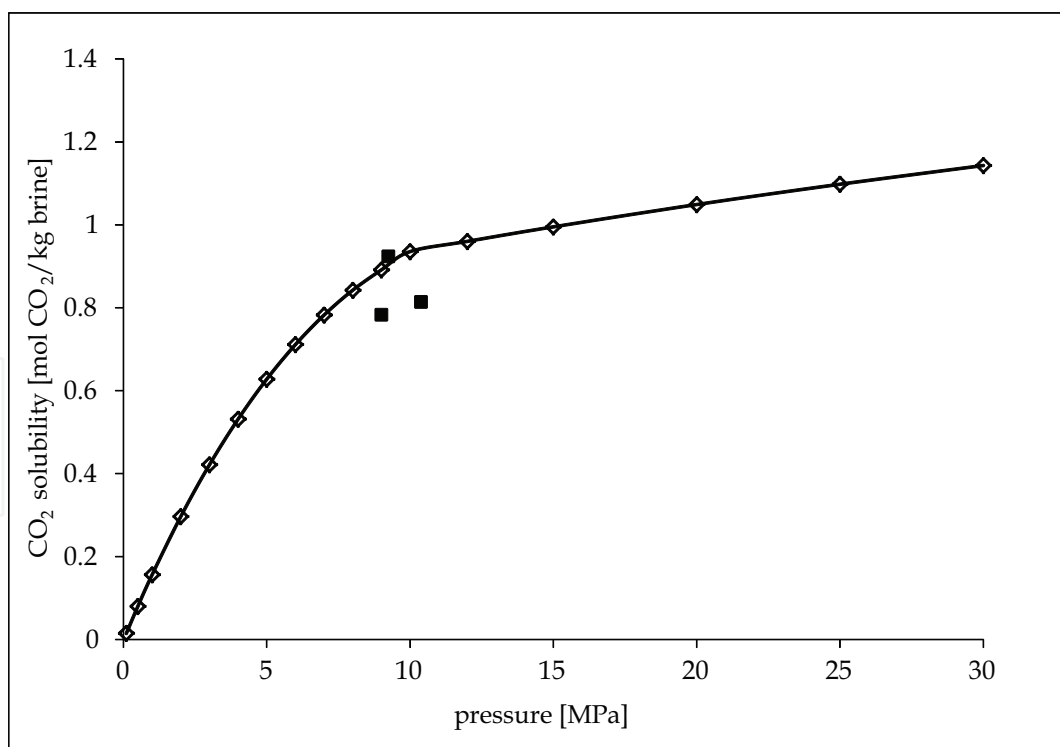


Fig. 3. CO₂ solubility increases with pressure increase. The data shown was computed with Duan and Sun (2003) and Duan et al. (2006)'s CO₂ solubility calculator. The black squares show experimental data points measured by Nighswander et al. (1989), Li et al. (2004) and Kiepe et al. (2002)

4.2 Effect of temperature on CO₂ solubility in brine

CO₂ solubility decreases with increasing temperature as shown in Figure 4. Experimental data relevant for CCS and simulated data are displayed. The computational data curve (open diamonds) was calculated with Duan et al. (2003+2006)’s solubility calculator setting the pressure to 10 MPa and salinity to 1 mole NaCl/kg.

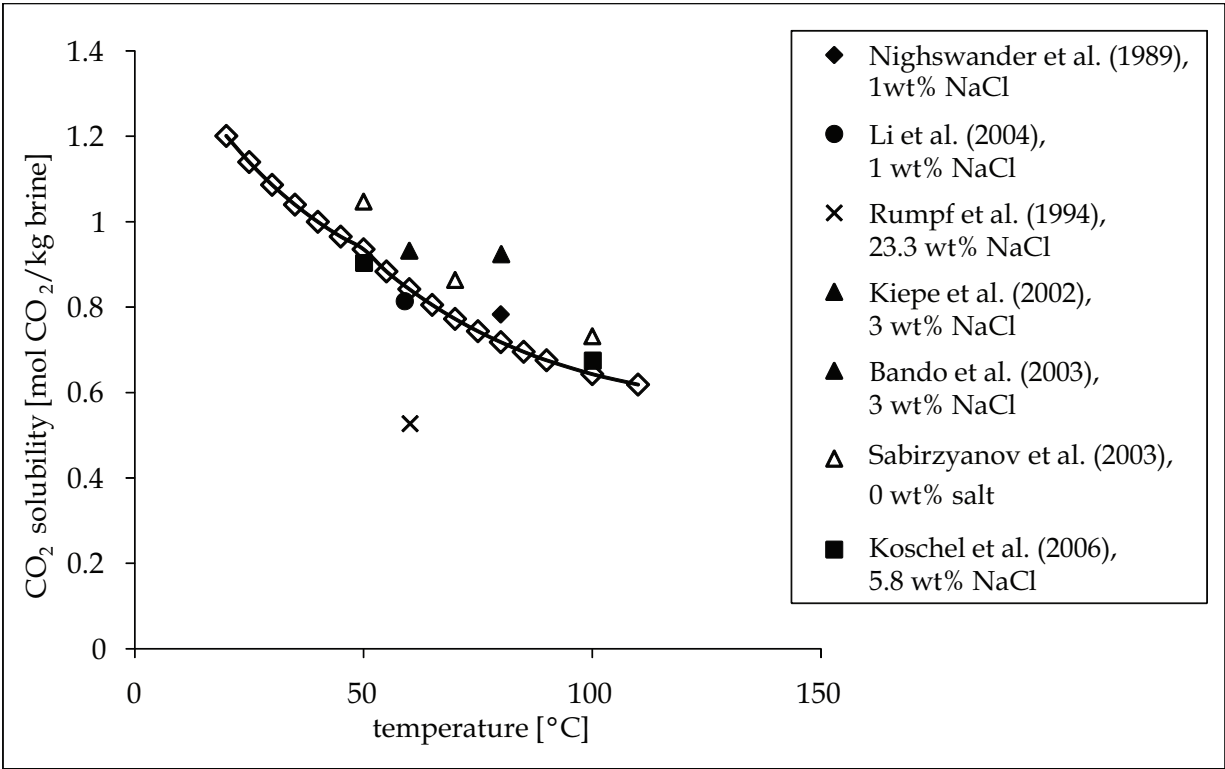


Fig. 4. CO₂ solubility versus temperature at high pressures. The substantially lower value measured by Rumpf et al. (1994) is caused by the high brine salinity (cp. section 4.3). Duan et al.’s (2003+2006) data is simulation data (open diamonds); the other points are experimentally measured values

4.3 Effect of brine salinity on CO₂ solubility in brines

CO₂ solubility decreases with increasing salinity as show in Figure 5. The open diamonds show simulated data calculated with Duan et al.’s (2003+2006) CO₂ solubility calculator setting the temperature to 323 K and the pressure to 10 MPa. It appears that Duan et al.’s. (2003+2006) model slightly over predicts CO₂ solubilities. The other points shown are experimentally determined values.

Moreover, the type of dissolved salt has an influence on CO₂ solubility. Yasunishi and Yoshda (1979) studied CO₂ solubilities at atmospheric pressure in a wide variety of salt solutions, these salts included NaCl, KCl, Na₂SO₄, MgCl₂, CaCl₂, K₂SO₄, MgSO₄, BaCl₂, AlCl₃, Al₂(SO₄)₃ among others. They found that for the same electrolyte concentration, KCl solutions can absorb more CO₂ than NaCl solutions, while CaCl₂ and MgCl₂ solutions absorb approximately the same amount of CO₂. Monovalent NaCl or KCl solutions with the same salt concentration absorb more CO₂ than their divalent CaCl₂ or MgCl₂ counterparts. For example Yasunishi and Yoshda (1979) measured at atmospheric pressure and 298 K that a 4.216 mol/L NaCl solution absorbs $L = 0.3144$ (L is the Ostwald coefficient, $L = V_g/V_l$ with

V_g = volume of CO_2 absorbed and V_l = volume of absorbing brine) while a 4.131 mol/L KCl solution absorbed $L = 0.4703$. For a 3.955 mol/L MgCl_2 solution they measured $L = 0.1648$. Chloride salt solutions absorbed more CO_2 than the corresponding sulphate solutions (that was tested for Na^+ , K^+ , Al^{3+} and Mg^{2+}).

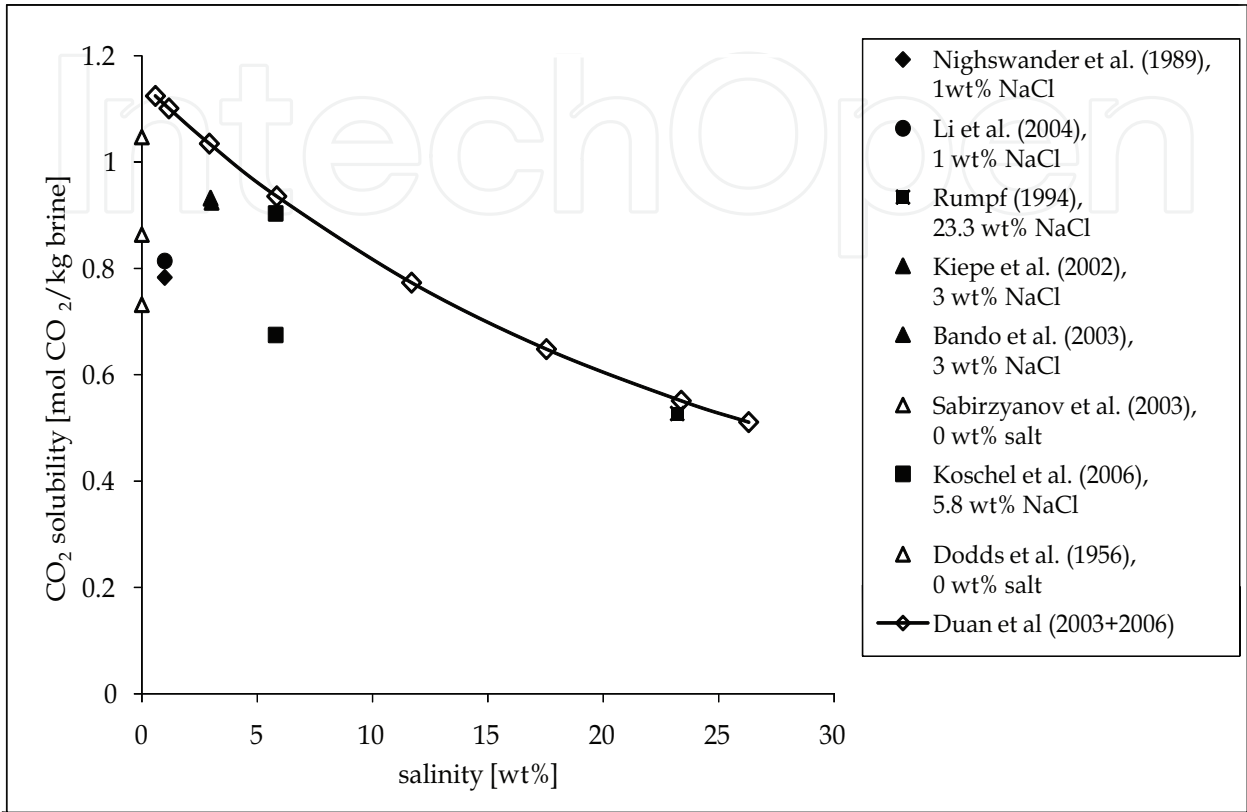


Fig. 5. CO_2 solubility as a function of brine salinity. The open diamonds represent data calculated with Duan et al. (2003+2006)’s CO_2 solubility calculator; the other points are experimentally measured values

Enick and Klara (1990) tested the influence of dissolved solids on CO_2 solubility in the temperature and pressure ranges 298-523 K and 3.40-72.41 MPa. Based on their results they developed an empirical equation for estimating salinity effects on CO_2 solubility (equation 1).

$$Y_{\text{CO}_2, \text{brine}} = Y_{\text{CO}_2, \text{pureH}_2\text{O}} \left(1 - 0.04893414 \cdot S + 0.001302838 \cdot S^2 - 0.00001871199 \cdot S^3 \right) \tag{1}$$

where

- $Y_{\text{CO}_2, \text{brine}}$ = CO_2 solubility in brine (mass fraction)
- $Y_{\text{CO}_2, \text{pureH}_2\text{O}}$ = CO_2 solubility in pure water (mass fraction)
- S = salinity of brine (weight percent)

4.4 Theoretical model for computing CO_2 solubilities

Duan and Sun (2003, 2006) developed an equation (equation 2) which can predict CO_2 solubilities in brine as a function of temperature (range 273-533 K), pressure (range 0-20 MPa) and salinity (different salts/ions can be considered: Na^+ , K^+ , Mg^{2+} , Ca^{2+} , Cl^- , SO_4^{2-}).

$$\begin{aligned} \ln m_{\text{CO}_2} = & \ln y_{\text{CO}_2} F_{\text{CO}_2} p - \mu_{\text{CO}_2}^{l(0)} / RT \\ & - 2\lambda_{\text{CO}_2-\text{Na}} (m_{\text{Na}} + m_{\text{K}} + 2m_{\text{Ca}} + 2m_{\text{Mg}}) \\ & - \zeta_{\text{CO}_2-\text{Na-Cl}} m_{\text{Cl}} (m_{\text{Na}} + m_{\text{K}} + m_{\text{Mg}} + m_{\text{Ca}}) + 0.07m_{\text{SO}_4} \end{aligned} \quad (2)$$

Here T is the temperature, p the pressure, R is the universal gas constant, m is the molality of components dissolved in water, y_{CO_2} is the mole fraction of CO_2 in the vapour phase, F_{CO_2} is the fugacity coefficient of CO_2 , $\mu_{\text{CO}_2}^{l(0)}$ is the standard chemical potential of CO_2 in the liquid phase, $\lambda_{\text{CO}_2-\text{Na}}$ is the interaction parameter between CO_2 and Na^+ and $\zeta_{\text{CO}_2-\text{Na-Cl}}$ is the interaction parameter between CO_2 and Na^+ , Cl^- .

The fugacity F_{CO_2} can be calculated via a fifth-order virial equation of state (equation 3). The coefficients c_i are stored in a look-up table (Duan et al. 2006) and they vary with the pressure and temperature regime.

$$\begin{aligned} F_{\text{CO}_2} = & c_1 + [c_2 + c_3T + c_4 / T + c_5 / (T - 150)]p + [c_6 + c_7T + c_8 / T]p^2 \\ & + [c_9 + c_{10}T + c_{11} / T] \ln p + [c_{12} + c_{13}T] / p + c_{14} / T + c_{15}T^2 \end{aligned} \quad (3)$$

For the pressure and temperature regime most relevant to CCS, i.e. for a pressure below 100 MPa and a temperature range 273-340 K, the coefficients are inserted and shown in equation (3b).

$$\begin{aligned} F_{\text{CO}_2} = & -0.71734882 + [0.00015985379 - 4.9286471 \cdot 10^{-7}T]p \\ & + [-2.7855285 \cdot 10^{-7} + 1.1877015 \cdot 10^{-9}T]p^2 \\ & + [-96.539512 + 0.44774938 \cdot T] / p \\ & + 101.81078 / T + 5.3783879 \cdot 10^{-6}T^2 \end{aligned} \quad (3b)$$

The mole fraction of CO_2 in the vapour phase y_{CO_2} can be computed with equation (4)

$$y_{\text{CO}_2} = (p - p_{\text{H}_2\text{O}}) / p \quad (4)$$

where $p_{\text{H}_2\text{O}}$ is the water vapour pressure which can be estimated with the empirical equation (5) (Duan and Sun 2003).

$$p_{\text{H}_2\text{O}} = \left(\frac{p_c T}{T_c} \right) \left[\frac{1 - 38.640844(-t)^{1.9} + 5.8948420 \cdot t + 59.876516 \cdot t^2}{+26.654627 \cdot t^3 + 10.637097 \cdot t^4} \right] \quad (5)$$

where $t = (T - T_c) / T_c$ and T_c and p_c are the critical temperature and critical pressure of water ($T_c = 647.29$ K, $p_c = 22.085$ MPa).

The parameters $\mu_{\text{CO}_2}^{l(0)}$, $\lambda_{\text{CO}_2-\text{Na}}$ and $\zeta_{\text{CO}_2-\text{Na-Cl}}$ are estimated with equation (6) and Table 1 (Duan and Sun 2003).

$$\begin{aligned} \text{Par}(T, p) = & a_1 + a_2 \cdot T + a_3 / T + a_4 \cdot T^2 + a_5 / (630 - T) + a_6 \cdot p + a_7 \cdot p \cdot \ln T \\ & + a_8 \cdot p / T + a_9 \cdot p / (630 - T) + a_{10} \cdot p^2 / (630 - T)^2 + a_{11} \cdot T \cdot \ln p \end{aligned} \quad (6)$$

T-p coefficient	$\mu_{CO_2}^{(0)}/RT$	λ_{CO_2-Na}	$\zeta_{CO_2-Na-Cl}$
a ₁	28.9447706	-0.411370585	3.36389723E-4
a ₂	-0.0354581768	6.07632013E-4	-1.98298980E-5
a ₃	-4770.67077	97.5347708	
a ₄	1.02782768E-5		
a ₅	33.8126098		
a ₆	9.04037140E-3		
a ₇	-1.14934031E-3		
a ₈	-0.307405726	-0.0237622469	2.12220830E-3
a ₉	-0.0907301486	0.0170656236	-5.24873303E-3
a ₁₀	9.32713393E-4		
a ₁₁		1.41335834E-5	

Table 1. CO₂ solubility interactions parameters (Duan et al. 2003, 2006)

4.5 Effect of injection depth on CO₂ solubilities

In a deep saline aquifer or oil reservoir high pressures and elevated temperatures are found. The pore pressure at depth is usually assumed to be equal to the hydrostatic pressure; a typical hydrostatic pressure gradient is 10.35 MPa/1000m (Dake 2007). In addition a geothermal gradient exists, the reservoir temperature increases with depth. Average typical geothermal gradients are 25-30 K/1000 m (Fridleifsson et al. 2008). Average temperatures and pressures at depth are listed in Table 2, they were calculated assuming typical pressure and temperature gradients and a surface temperature of 293 K. The surface temperature needs to be adjusted for each specific geographical location, e.g. average temperature is low in Norway (average temperature throughout the year is around 281 K) while average yearly temperature is high in Saudi Arabia (298 K).

As stated above CO₂ solubility decreases with increase in temperature, but increases with increase in pressure. In Table 2 CO₂ solubilities calculated with Duan’s web based CO₂-solubility calculator (Duan et al. 2003, 2006) are shown. The pressure effect over compensates the temperature effect so that CO₂ solubility increases with reservoir depth up to a depth of approximately 900m when it reaches a plateau.

With regard to storage of CO₂ in a supercritical phase optimal CCS conditions are conditions where the CO₂ density ρ_{CO_2} is maximal, because then a maximum mass of CO₂ can be stored in the same rock pore space. Thermodynamically ρ_{CO_2} increases with pressure but decreases with temperature. ρ_{CO_2} as a function of depth increases monotonically as the pressure effect also over compensates the temperature effect (Table 2).

4.6 Effect of presence of oil (CCS in oil reservoirs)

CO₂ can also be injected into depleted oil reservoirs although storage capacities are much smaller than in aquifers (IPCC 2005). It is estimated that 50 Gt of CO₂ can be stored in this way worldwide (Firoozabadi and Cheng 2010) which is roughly 1.5 times of what is emitted per year. So this is clearly not the solution to mitigate global warming, however CO₂ solubility in oil is very high, up to 60-80 mol% of CO₂ can be dissolved (De Ruiter et al. 1994, Kokal and Sayegh 1993, Emera and Sarma 2006, Firoozabadi and Cheng 2010).

CO₂ solubility generally increases with pressure and it is higher at lower temperatures. If the temperature is below the critical CO₂ temperature ($T_c = 304.13$ K), then CO₂ solubility increases until the CO₂ liquefaction pressure is reached (circa 5.88 MPa), then it levels off

with further pressure increase. CO₂ solubility also depends on oil composition and for light oils CO₂ can be completely miscible. For example De Ruiters et al. (1994) measured a strong increase of CO₂ solubility with pressure in two crude oils, at low pressures (0.69 MPa) the gas-oil ratio (GOR) was approximately 5.3 m³/m³, and GOR increased rapidly up to the CO₂ liquefaction pressure when it reached 71 m³/m³ and 102 m³/m³, respectively. With a further pressure increase GOR stayed approximately constant. The experimental temperature in De Ruiters et al. experiments was low (290 K). If the temperature is above T_c as expected for CCS conditions, then CO₂ solubility monotonically increases; but it is nominally lower as compared to lower temperatures (Kokal and Sayegh 1993).

Depth [m]	Temperature [K]	Pressure [MPa]	ρ_{CO_2} [kg/m ³]*	CO ₂ solubility [mol/kg]**
0	293	0.1	1.8	0.0307
100	296	1.135	21.8	0.3036
200	299	2.17	43.7	0.5037
300	302	3.205	68.3	0.6496
400	305	4.24	96.5	0.7542
500	308	5.275	130	0.8274
600	311	6.31	171.7	0.8769
700	314	7.345	221.8	0.9082
800	317	8.38	311.8	0.9260
900	320	9.415	391.9	0.9338
1000	323	10.45	412.8	0.9353
1100	326	11.485	449.6	0.9344
1200	329	12.52	486.3	0.9334
1300	332	13.555	522.7	0.9330
1400	335	14.59	561.3	0.9335
1500	338	15.625	576.1	0.9348

* estimated from Span and Wagner (1996).

** 1 mol/kg NaCl brine, calculated with Duan et al.'s (2003, 2006) calculator.

Table 2. Variation of temperature, pressure, CO₂ solubility and CO₂ density with depth

In case of heavy oils CO₂ dissolves into the oil phase while some light oil fractions are extracted into the CO₂ phase. Depending on the oil and thermophysical condition, vapour-liquid, liquid-liquid, liquid-supercritical fluid, liquid-liquid-vapour phase behaviours are observed. The densities of CO₂-saturated oil increase at lower temperature (294 K) while they decrease at higher temperature (e.g. 413 K) (Kokal and Sayegh 1993).

This makes CO₂ a very efficient solvent for crude oil extraction in tertiary oil recovery processes (Green and Willhite 1998, Blunt et al. 1993). The dissolved CO₂ reduces oil viscosity significantly which improves the mobility ratio oil-injected fluid (for improving production) and results in a much better reservoir sweep efficiency. The flow of oil in the reservoir is improved by the improved oil relative permeability, which leads to increased oil production. In addition, CO₂ which dissolves into the oil causes oil swelling (up to 50-60%, Firoozabadi and Cheng 2010) which also leads to enhanced oil production. One side effect of CO₂ addition to crude oil is that large asphaltene molecules precipitate (crude oil is a very complex fluid (cp. Table 3) with a multitude of components including such large asphaltene

components which are dissolved in the oleic phase under reservoir conditions, Dandekar 2006) which renders the rock surface more oil-wet which again changes multi-phase fluid dynamics in the reservoir. According to contact angle studies (Dickson et al. 2006, Espinoza and Santamaria 2010) this can result in CO₂-wet surfaces which would eliminate the possibility of capillary trapping of CO₂. Also the surface area CO₂-brine would most likely be affected by such wettability effects, which in turn would affect CO₂ dissolution kinetics (cp. section 5.1 and equations 12 and 17).

Component	mole %*	mole %**
Methane, CH ₄	45.93	36.47
Ethane, C ₂ H ₆	7.32	9.67
Propane, C ₃ H ₈	6.42	6.95
n-Butane, C ₄ H ₁₀	3.87	3.93
i-Butane, C ₄ H ₁₀	1.42	1.44
n-Pentane C ₅ H ₁₂	2.05	1.41
i-Pentane C ₅ H ₁₂	1.68	1.44
Hexanes	2.93	4.33
C ₇	2.30	C ₇₊ 33.29
C ₈	2.21	
C ₉	1.66	
C ₁₀	1.97	
C ₁₁	1.61	
C ₁₂	1.39	
C ₁₃	1.36	
C ₁₄	1.28	
C ₁₅	1.22	
C ₁₆	1.09	
C ₁₇	1.04	
C ₁₈	0.98	
C ₁₉	0.77	
C ₂₀₊	6.63	
Hydrogen sulphide, H ₂ S	0.60	0
Carbon dioxide, CO ₂	1.47	0.91
Nitrogen, N ₂	0.81	0.16

Table 3. Typical composition of black crude oil (*Dandekar 2006, **McCain 1990). Of course the exact compositions of crude oils are extremely complex and vary widely depending on the exact geographical location. The C fractions C₆₊ upwards contain many isomers and also hydrocarbons with additional functional groups (e.g. alcohol, ester, carbonyl, amine, etc. pp.). Crude oil also contains metal cations (e.g. Vanadium)

Moreover there is a very important reservoir engineering aspect associated with depleted oil reservoirs; reservoir pressure is low (because of oil production) and CO₂ can be injected at fairly high rates and comparatively large quantities of CO₂ can be stored. It is important not to exceed the fracture pressure of the caprock which would result in catastrophic leakage of

CO₂. Exceeding the capillary entry pressure of CO₂ into the caprock should also be avoided (then CO₂ will also flow through the caprock although very slowly because of the very low permeability of the caprock shale) resulting in potential CO₂ leakage to the surface.

Estimates suggest that many millions of tons of crude oil are produced yearly via enhanced oil recovery with CO₂ (CO₂-EOR) (Firoozabadi and Cheng (2010)). Crude oil production could be further increased if more CO₂ would be used but such CO₂-EOR schemes should have a CO₂ storage element.

In principle oil would be a very good storage medium for CO₂ (provided that the oil does not migrate upwards after CO₂-takeup, so ideally the process would be designed in such a way that oil density increases), but of course oil is an economically valuable commodity and will be produced, so oil production schemes need to be combined with CCS schemes and optimized, essentially as much oil as possible needs to be recovered while storing as much CO₂ as possible.

Reservoir simulations can calculate such CO₂-EOR recovery/injection schemes over several years (Qi et al. 2008, Firoozabadi 2011), one complication here is the three-phase flow and the associated complex fluid thermodynamics occurring in the reservoir. This includes mass transfer of CO₂ into the oil and aqueous phases.

In summary, most of the current CCS schemes which are online are actually EOR processes because of profitability. Example projects are the Weyburn-Midale project in Canada, which started in the year 2000. 1.8 Mt/a of CO₂ are injected into a depth of 1500m into a depleted oil reservoir (PTRC 2011, Pentland 2011). 225 m³ of CO₂ produce 0.12 m³ extra crude oil there. Another CO₂-EOR project is underway in the Salt Creek field in Wyoming, USA; here 2.09 Mt of CO₂ are injected yearly and more than 1.2×10^6 m³ of incremental crude oil have been recovered so far and it is planned to store 50 Mt of CO₂ in total (Andarko 2010, Pentland 2011).

4.7 Effect of presence of gas (gas reservoirs or oil reservoirs with a gas cap)

CO₂ can also be injected into depleted gas reservoirs in order to produce additional gas, this is called enhanced gas recovery (EGR). The injected CO₂ increases reservoir pressure which supports gas production. As in the case of oil reservoirs or indeed aquifers the caprock failure stress must not be exceeded. Natural gas is a mixture of various components (cp. Table 4); the exact composition varies with the location of the gas fields and it is determined by the original hydrocarbon generation (Dandekar 2006).

In the reservoir, the CO₂ flood front mixes with the natural gas by dispersion and diffusion. In parallel to the CO₂ – gas mixing process, CO₂ also equilibrates with the formation brine, similar to the mixing processes occurring in deep saline aquifers. The main advantage of CO₂-EGR is profitability as in CO₂-EOR, and an optimum between additional gas production and CO₂ sequestration needs to be found. There are several CO₂-EGR pilot units where these processes are tested, e.g. in the Lacq demonstration project in southwest France, 10⁵ t of CO₂ will be injected and stored in a depleted gas field at a depth of 4500m (Total 2011).

A thorough study of nine natural gas fields (including sandstone and carbonate reservoirs) concludes that the main trapping mechanism over millennial timescales is dissolution trapping. At most 18% of injected CO₂ is stored as a solid mineral phase (Gilfillan et al. 2009) and mineral trapping is predicted to happen only for siliclastic reservoirs.

In the case of oil reservoirs with a gas cap, the mixing thermodynamics are a combination of CO₂-gas mixing, CO₂ dissolution in oil and CO₂ dissolution in brine. These complex

processes are topic of current research (DaVega 2011). These mixing processes result in three-phase flow in the reservoir (oil, gas and water flow as separated phases); in addition it is possible that additional phases are formed (e.g. a second immiscible oil phase or a solid asphaltene phase) which can further complicate fluid dynamics at the pore-scale and in the whole reservoir. Depending on rock surface wettabilities CO₂ dissolution into brine can be slowed down, e.g. in case of a water-wet surface water covers the rock surface, and an oil layer may separate the brine from the CO₂ (Piri and Blunt 2005). This oil layer then essentially acts as a barrier through which the CO₂ has to pass in order to reach the brine and to be stored there safely by the dissolution trapping mechanism.

Component	mole %
Methane, CH ₄	70-98
Ethane, C ₂ H ₆	1-10
Propane, C ₃ H ₈	trace - 5
Butanes, C ₄ H ₁₀	trace - 2
Pentanes, C ₅ H ₁₂	trace - 1
Hexanes C ₆ H ₁₄	trace - 0.5
Heptanes C ₇ H ₁₆	trace - 0.5
Carbon dioxide, CO ₂	trace - 5
Nitrogen, N ₂	trace - 15
Hydrogen sulphide, H ₂ S	trace - 3
Helium, He	0 - 5

Table 4. Typical composition of natural gas (McCain 1990). Apart from methane and ethane traces of medium sized hydrocarbons can be found. In addition, natural gas can contain significant amounts of H₂S, CO₂ or N₂ – up to 90 mol% (Firoozabadi and Cheng 2010). Such non-hydrocarbon gases usually need to be separated out of the production stream in order to achieve sellable gas quality

5. Kinetics of CO₂ dissolution into formation brines

Dissolution kinetics of CO₂ into brine in a reservoir are driven by four main factors, namely molecular diffusion of CO₂ into brine, dispersion during flow, convection of CO₂-saturated (heavier) brine in the reservoir and flow of the scCO₂ phase in the reservoir. These mechanisms are described in more detail in the following paragraphs.

5.1 Molecular CO₂ diffusion into reservoir brines.

Molecular diffusion in natural groundwater systems is usually a time-dependent unsteady-state process. This is described by Fick’s second law (equation 7). The driving force behind molecular diffusion is the concentration gradient, essentially the entropy of the system is increased by molecular diffusion.

$$D \frac{\partial^2 c}{\partial z^2} = \frac{\partial c}{\partial t} \quad (7)$$

where

D = diffusion coefficient

c = concentration

t = time

z = depth

A limited number of measurements of the CO₂ diffusion coefficient in water D_{CO₂-H₂O} at high pressure have been conducted. Renner (1988) measured D_{CO₂-H₂O} for 0.25 N NaCl brine at 311 K for a pressure range 1.54-5.86 MPa and recorded D_{CO₂-H₂O} values in the range 3.07-7.35 × 10⁻⁹ m²/s. More measurements at atmospheric pressure were conducted and D_{CO₂-H₂O} values between 1.8 × 10⁻⁹ – 8 × 10⁻⁹ m²/s (Mazarei and Sandall 1980, Unver and Himmelblau 1964) were reported. Based on these datasets, Renner (1988) developed an empirical-statistical expression (equation 8).

$$D_{CO_2-H_2O} = 6391 \mu_{CO_2}^{-0.1584} \mu_{H_2O}^{6.911} \quad (8)$$

where

μ_{CO₂} = CO₂ viscosity

μ_{H₂O} = H₂O viscosity

Renner's analysis (1988) indicated that water viscosity and CO₂ viscosity were highly correlated with the diffusion coefficient, but molecular weight of CO₂, molar volume of CO₂, pressure or temperature were not statistically significant. However Renner states in his paper and Renner's data show that D_{CO₂-H₂O} increases with an increase in pressure. Therefore it can be expected that CO₂ diffusion processes under CCS conditions are faster than at atmospheric pressure conditions – which is positive news for dissolution trapping as it minimizes leakage risks by absorbing the mobile CO₂ faster in the aqueous phase.

Hirai et al. (1997) measured D_{CO₂-H₂O} via laser-induced fluorescence at 286 K and 29.4 and 39.3 MPa (D_{CO₂-H₂O} = 1.3 × 10⁻⁹ and 1.5 × 10⁻⁹ m²/s). Their results fit perfectly with the empirical equation (9) suggested by Wilke and Chang (1955). ι is an association parameter equal to 2.26 for water. The experimental data measured by Shimizu et al. (1995) (D_{CO₂-H₂O} is approximately 1.8 × 10⁻⁹ m²/s at 286 K and 9-13 MPa) is however 40% larger than predicted by equation (9). Hirai's data and the Wilke-Chang equation both indicate that D_{CO₂-H₂O} increases slightly with pressure.

$$D_{CO_2-H_2O} = 7.4 \cdot 10^{-8} \left(\iota M_{H_2O} \right)^{0.5} T / \left(\mu_{CO_2} V_{CO_2}^{0.6} \right) \quad (9)$$

More recently, Mutoru et al. (2010) developed a semi-empirical model for calculating diffusion coefficients for infinitely diluted CO₂ and water mixtures (equation 10) based on 187 experimental data points. The subscript 1 denotes CO₂ and the subscript 2 denotes water. However, in case of water diffusing into the CO₂ phase, subscript 1 denotes water and subscript 2 CO₂.

$$D_{CO_2-H_2O} = \frac{k_1 (M_{12} \epsilon_{12})^{k_2} T_{r,2}^{k_3}}{p_{r,2}^{k_4} (\mu_2 c_2)^{k_5}} \quad (10)$$

where

$$M_{12} = \left(\frac{1}{M_{H_2O}} + \frac{1}{M_{CO_2}} \right)^{-1} \quad (11)$$

M = molecular mass

ϵ_1 = dipole moment

ϵ_{12} = ratio of dipole moments, ϵ_1 / ϵ_2

$k_1 = 10^{-7.23389}$

$k_2 = 0.135607$

$k_3 = 1.84220$

$k_4 = 2.41943 \times 10^{-3}$

$k_5 = 0.858204$

$T_{r,2}$ = reduced temperature

$p_{r,2}$ = reduced pressure

μ = viscosity

c_2 = solvent molar density (the solvent is defined here as the dominant component)

The advantage of Mutoru et al.'s (2010) model is that it incorporates the temperature and pressure effects on the total dipole moment of water and the induced dipole moment of CO₂. In addition, it can predict $D_{CO_2-H_2O}$ over the complete range from infinitely diluted CO₂ to infinitely diluted H₂O. From equation (10) it is clear that temperature has a stronger influence on $D_{CO_2-H_2O}$ than pressure. This is due to the strong dependence of the viscosity and the solvent molar density on the temperature. However, pressure influences are also strong as pressure determines equilibrium compositions (Mutoru et al. 2010).

An interesting perspective on CO₂ dissolution into brine is the consideration of the CO₂ droplet diameter (Hirai et al. 1997). Especially in the context of residual trapping; here the rising CO₂ plume is split into a large number of small disconnected CO₂ clusters at the trailing edge of the plume due to natural water influx or chase brine injection (Iglauer et al. 2010).

The drop diameter is expected to have a highly significant effect on CO₂ dissolution speed. A strong enhancement of CO₂ dissolution is expected for such small CO₂ bubbles as their CO₂-brine surface area is significantly increased compared with that of a single-cluster CO₂ plume.

The dissolution rate of CO₂ can be described by equation (12) (Hirai et al. 1997).

$$d(\rho_{CO_2} V) / dt = -kA(C_0 - C_\infty) \quad (12)$$

where

V = volume of the scCO₂ droplet,

A = surface area of the scCO₂ droplet

k = mass transfer coefficient

C_0 = surface concentration of the droplet

C_∞ = concentration at infinity

The mass transfer coefficient k is expressed in equation (13) for high Schmidt (Sc) numbers

$$Sh = 1 + (Sc + 1 / Re)^{1/3} 0.752 Re^{0.472} \quad (13)$$

where

Sh = Sherwood number ($= k \cdot d_{\text{CO}_2} / D_{\text{CO}_2-\text{H}_2\text{O}}$)

Sc = Schmidt number ($= \nu / D_{\text{CO}_2-\text{H}_2\text{O}}$)

Re = Reynolds number ($= u \cdot d_{\text{CO}_2} / \nu$)

d_{CO_2} = CO₂ droplet diameter

ν = kinematic viscosity

u = flow velocity

According to equation (13) small CO₂ droplets dissolve faster than large ones (assuming that $D_{\text{CO}_2-\text{H}_2\text{O}}$ is a constant). Essentially this is a formal description of how residual trapping enhances dissolution trapping. More research in this area would certainly improve understanding of the relation between residual and dissolution trapping.

Suekane et al. (2006) studied such mass transfer processes of scCO₂ dissolution into pure water in packed glass beads (measurement conditions were 313 K, 8.3 MPa, 70 μm bead diameter) and developed relation (14)

$$Sh' = 0.029 Re^{0.92} \quad (14)$$

with the modified Sherwood number Sh' . Sh' can be calculated with equations (15) and (16).

$$Sh' = \frac{\hat{k} d^2}{D} \quad (15)$$

$$\hat{k} = \frac{u}{L} \ln \left(1 - \frac{C}{C_\infty} \right) \quad (16)$$

where \hat{k} is the total mass transfer rate ($= kA$) and L is the length of the glass bead pack. Equation (16) is the solution of the one-dimensional steady state mass balance equation (17) with the boundary conditions $C = 0$ at $x = 0$.

$$u \frac{dC}{dx} = -kA(C - C_\infty) = \hat{k}(C - C_\infty) \quad (17)$$

Another interesting suggested correlation for $D_{\text{CO}_2-\text{H}_2\text{O}}$ has been put forward by Bahar and Liu (2008); they measured $D_{\text{CO}_2-\text{H}_2\text{O}}$ at 17.8 MPa and 356 K in 2 wt% NaCl brines and developed an empirical correlation between $D_{\text{CO}_2-\text{H}_2\text{O}}$ and the pressure p , temperature T , molecular weight MW , volume V and viscosity μ of the liquid (equation 18).

$$D_{\text{CO}_2} = 1.3678 \cdot 10^{-7} \frac{T^{1.47} MW^{2.2}}{V^{0.5} p \mu} \quad (18)$$

Bahar and Liu (2008) found that $D_{\text{CO}_2-\text{H}_2\text{O}}$ is higher for unsteady-state systems, and that the duration of the unsteady-state system strongly depends on the pressure and temperature.

5.1.1 Effect of temperature on CO₂ diffusion in water

As stated in equation (10), temperature has a clear effect on $D_{\text{CO}_2-\text{H}_2\text{O}}$. Unver and Himmelblau (1964) developed an empirical equation (19) for the dependence of $D_{\text{CO}_2-\text{H}_2\text{O}}$ on

temperature for atmospheric pressure within a temperature range between 279-338 K. D increases monotonically with temperature T (equation 19).

$$D = (A + BT + CT^2) \cdot 10^{-9} \quad (19)$$

$A = 0.95893$, $B = 0.024161$, $C = 0.00039813$ are constants for CO_2 , A , B and C adopt different values for other gases. D is given in m^2/s . Again, diffusion-dominated CO_2 dissolution is more effective at higher temperatures.

5.1.2 Effect of pressure on CO_2 diffusion in water

According to measurements conducted by Wilke-Chang (1955), Shimizu et al. (1995) and Hirai et al. (1997) $D_{\text{CO}_2\text{-H}_2\text{O}}$ (approximately $1.5 \times 10^{-9} \text{ m}^2/\text{s}$) is quasi independent of pressure in the tested range of circa 9-40 MPa. Their measurements were all performed at 286 K. However, Renner's (1988) measurements show that $D_{\text{CO}_2\text{-H}_2\text{O}}$ increases with pressure, this is supported by Mutora et al.'s (2010) analysis.

5.2 CO_2 diffusion into oil

Renner (1988) measured diffusion coefficients of CO_2 in decane $D_{\text{CO}_2\text{-C}_{10}}$ at a temperature of 311 K and in a pressure range 1.54-5.86 MPa. The results for $D_{\text{CO}_2\text{-C}_{10}}$ ranged from $1.97\text{-}11.8 \times 10^{-9} \text{ m}^2/\text{s}$. An increase in pressure led to an increase in $D_{\text{CO}_2\text{-C}_{10}}$ and measured diffusion coefficients in a vertical sandstone core were significantly higher than in a horizontal core; this might have been due to convective forces in the vertical core. Renner developed the empirical-statistical equation (20) for $D_{\text{CO}_2\text{-HC}}$ estimates.

$$D_{\text{CO}_2\text{-HC}} = 10^{-9} \mu_{\text{CO}_2}^{-0.4562} M_{\text{CO}_2}^{-0.6898} V_{\text{CO}_2}^{-1.706} p^{-1.831} T^{4.524} \quad (20)$$

where V_{CO_2} is the molar volume of CO_2 .

Model equation (10) can also estimate the diffusion coefficients of CO_2 -hydrocarbon (HC) systems $D_{\text{CO}_2\text{-HC}}$; predictions can be made for small alkane molecules (e.g. methane, ethane, butane) and polar H_2S .

5.3 Water diffusion into scCO_2 phase

Although not essential for CO_2 storage, it is noted for completeness that water diffuses and dissolves into the scCO_2 phase. Water solubility in scCO_2 is low, it increases with pressure. For a pressure range from 8.31-20.54 MPa at 313 K water mole fractions between 0.00053-0.00596 were measured, for a pressure range between 2.51-10.20 MPa at 323 K the water mole fraction measured ranged between 0.00251-0.0120 (Sabirzyanov et al. 2002).

Measured diffusion coefficients of water in CO_2 $D_{\text{H}_2\text{O-CO}_2}$ are reported to be much higher than diffusion coefficients of CO_2 into water. Values between 1.5×10^{-8} to $1.8 \times 10^{-9} \text{ m}^2/\text{s}$ were published for a pressure range between 7-20 MPa and a temperature of 298 K (Espinoza and Santamaria 2010). However the investigated temperature was lower than the expected temperature at CCS storage depths. The cited numbers could therefore be slightly different for actual CCS conditions. In addition, the semi-empirical model equation (10) can also estimate $D_{\text{H}_2\text{O-CO}_2}$ diffusion coefficients.

5.4 Dispersion of dissolved CO₂ due to flow through a porous medium

In addition to diffusion, dispersion occurs when a solute flows through a porous medium. This can essentially be understood as an unsteady irreversible mixing process of two miscible fluids which have different solute concentrations (e.g. brine saturated with CO₂ = fluid 1, and brine undersaturated with CO₂ = fluid 2). Dispersion can therefore influence CO₂ mass transfer as it changes the CO₂ concentration gradient. Dispersion is caused by several effects (Bear 1972, Özgür 2006, 2010):

1. The flow velocity profile in a pore (the flow velocity has a maximum in the middle of a pore).
2. Different flow velocities in different pores (the pores in a geological rock have a pore size distribution and therefore different flow resistances to fluids according to their size; faster flow happens in the pores with a larger diameter).
3. The complex tortuosity of the pores in the rock; some pores are longer and fluid flow takes longer.
4. Interactions of the solute with the rock matrix/rock surface.
5. Chemical reactions, e.g. ion exchange, of the solute with species in the brine or on the rock surface.

Bear (1972) distinguishes between mechanical dispersion and hydrodynamic dispersion. He defined hydrodynamical dispersion as the sum of mechanical dispersion plus molecular diffusion. The dispersion described above - and all dispersion mentioned in this text - is the same as Bear's mechanical dispersion.

At reservoir scale dispersion can be described by equation (21) where D_{dis} is the dispersion coefficient, u is the average pore flow velocity and α the dispersivity (Bear, 1972; Özgür, 2006).

$$D_{dis} = u\alpha \quad (21)$$

The dispersivity α is a property of the reservoir and it depends on the heterogeneity of the porous medium and the length of flow. Schulze-Makuch (2005) reviewed 307 datasets and suggested $\alpha_L = c L^{0.5}$ (where α_L is the longitudinal dispersivity, L is flow distance and c varies between 0.01 m for sandstones and unconsolidated material and 0.8 m for carbonates). A detailed discussion of dispersion and dispersivities is given by Bear (1972).

5.5 Convection of CO₂-enriched brine in the reservoir

Convection - here defined as flow at the reservoir scale induced by gradients in density, concentration or heat - can potentially move large quantities of dissolved CO₂ through the formation. A resistance threshold has to be overcome for convection to commence, this threshold can be assessed with the Rayleigh number Ra (equation 22) (Riaz et al. 2006).

$$Ra = \frac{K\Delta\rho gH}{\phi D\mu} \quad (22)$$

where

K = permeability

$\Delta\rho$ = density difference (between brine with high CO₂ concentration and brine with low CO₂ concentration)

g = gravitational constant

H = reservoir depth

D = diffusion coefficient

μ = brine viscosity

ϕ = porosity

Above a critical Rayleigh number Ra_c convection will occur; Ra_c is a function of the boundary conditions of the system (Weatherill et al. 2004), e.g. for a homogenous reservoir where the horizontal boundaries are impermeable and perfect heat conductors $Ra_c = 4\pi^2$ (Lindeberg and Wessel-Berg 1997).

One mechanism which can trigger convective flow is dissolution of CO_2 into brine, which can increase brine density by 1% under CCS conditions (Ennis-King and Paterson 2005).

Based on an expression suggested by Garcia (2001) Özgür (2006) developed an equation (23) with which the effect of aqueous CO_2 concentration on brine density can be estimated.

$$\rho_{CO_2, brine} = \frac{\rho_{brine}}{1 - Y_{CO_2} \left(1 - \frac{V_{m, brine}}{M} \rho_{brine} \right)} \quad (23)$$

where

$\rho_{CO_2, brine}$ = density of CO_2 -enriched brine

ρ_{brine} = original brine density

Y_{CO_2} = dissolved CO_2 mass fraction

$V_{m, brine}$ = apparent molar volume of CO_2 in brine

M = molecular weight of CO_2

6. Reservoir scale dissolution trapping

In the context of reservoir flow where dissolved CO_2 molecules are transported, convection, dispersion, diffusion and maximum CO_2 solubilities can all play a significant role. This is an active field of research, and three literature examples are presented where CO_2 solute transport was modelled at reservoir scale.

Interesting conclusions extracted from these computations are that dissolution trapping is favourably done in a high permeability reservoir. Moreover CO_2 dissolution can significantly reduce reservoir pressure (the pressure is increased by CO_2 injection, but only a maximum reservoir pressure is tolerable, the fracture pressure), improving injectivity, i.e. CO_2 can be injected at a faster rate, and more CO_2 can be stored in total – provided that CO_2 dissolves at an adequate rate or CO_2 injection is slow enough.

6.1 Özgür model (2006)

Özgür (2006, 2010) modelled diffusion and convection in an aquifer with the diffusion-

convection equation (24). $D \frac{\partial^2 c}{\partial z^2}$ is the diffusion term and $\frac{u}{\phi} \frac{\partial c}{\partial z}$ is the convection term.

$$D \frac{\partial^2 c}{\partial z^2} - \frac{u}{\phi} \frac{\partial c}{\partial z} = \frac{\partial c}{\partial t} \quad (24)$$

This model considers one-dimensional vertical flow in an aquifer, temperature and CO_2 gas cap pressure are assumed to be constant; chemical reactions are ignored and the porous

medium is considered to be isotropic and homogenous. The aquifer thickness is set to 100m, porosity to 20%, temperature to 323 K, pressure to 7.6 MPa and $D_{\text{CO}_2\text{-H}_2\text{O}}$ to $3 \times 10^{-9} \text{ m}^2/\text{s}$. Equation 24 can be re-written in dimensionless form

$$\frac{1}{Pe} \frac{\partial^2 C_D^2}{\partial Z_D^2} - \frac{\partial C_D}{\partial Z_D} = \frac{\partial C_D}{\partial t_D} \quad (25)$$

with

$$C_D = \frac{C_{\text{CO}_2}}{C_{\text{CO}_2, \text{saturated}}}, \text{ normalized concentration of CO}_2 \text{ in brine at CCS conditions}$$

$$Z_D = \frac{z}{H}, \text{ normalized reservoir height, H is the total thickness of the aquifer}$$

$$t_D = \frac{ut}{\phi H}, \text{ dimensionless time}$$

$$Pe = \frac{uH}{\phi D}, \text{ Peclet number} = \text{ratio between transport by convection/transport by molecular diffusion (Bear 1972)}$$

Özgür (2006, 2010) solved the diffusion-convection equation (25) for one set of initial and boundary conditions:

Initial conditions:

$$C_D = 0 \text{ for } t_D = 0 \text{ and all } Z_D$$

Boundary conditions:

$$\text{at } Z_D = 0 : \quad C_D = 1 \text{ for } t_D > 0$$

$$\text{at } Z_D = 1 : \quad \frac{\partial C_D}{\partial Z_D} = 0$$

The solution for equation (25) is then (Lake 1989)

$$C_D = \frac{1}{2} \operatorname{erfc} \left(\frac{Z_D - t_D}{2\sqrt{\frac{t_D}{Pe}}} \right) + \frac{e^{Z_D Pe}}{2} \operatorname{erfc} \left(\frac{Z_D - t_D}{2\sqrt{\frac{t_D}{Pe}}} \right) \quad (26)$$

where

$$\operatorname{erf}(\beta) = \frac{2}{\sqrt{\pi}} \int_0^\beta e^{-t^2} dt \text{ is the error function}$$

and $\operatorname{erfc}(\beta) = 1 - \operatorname{erf}(\beta)$ is the complementary error function.

Özgür (2006, 2010) also conducted numerical modelling studies, and found that convection rate strongly increases with increasing permeability and dissolution trapping is strongly accelerated thereby. In diffusion dominated systems the dissolution rate is very slow, however, and only after 10^7 years the considered aquifer was completely saturated with CO_2 . Higher dispersivity generally supports dissolution trapping.

It is moreover interesting to note that the mixing zone length ΔZ_D (which is defined as the distance between the points $C_D = 0.1$ and $C_D = 0.9$ in the reservoir, Lake 1989) reaches a value of 0.9 in diffusion dominated systems only after 10^5 years in Özgür's model. Increased porosity slightly increases ΔZ_D , and ΔZ_D strongly increases with permeability once

a threshold (when convection sets in) is reached. In convection dominated systems, porosity decreases Δz_D , while dispersivity slightly increases Δz_D .

6.2 The Lindeberg/Wessel-Berg model (1997)

Lindeberg and Wessel-Berg (1997) modelled the onset of convection in aquifers into which CO_2 has been injected. The water column in such aquifers can be unstable because of the density gradient introduced by molecular CO_2 diffusion into the brine.

In their simulation they solved the Darcy equation, heat conduction equation, equation of continuity and energy equation (details are described very thoroughly by Bear (1972)). They also included the equation of diffusion (27) so that diffusive mass transfer was considered.

$$\frac{\partial \rho}{\partial t} + \frac{j}{\phi} \nabla \rho = \nabla \cdot \left(\frac{D}{\phi} \nabla \rho \right) \quad (27)$$

where j is the volume flux.

They calculated the Rayleigh numbers for an array of model reservoirs, spanning a temperature range from 303-363 K, a pressure range from 10-30 MPa, a permeability range from 100-2000 mD, while the porosity was a constant 30%. Variations in pressure and temperature resulted in brine density variations between 1013.5-1036 kg/m³ and molecular diffusion coefficients between $2.2\text{-}6.3 \times 10^{-9}$ m²/s. The brine density difference $\Delta \rho$ was 14.42 kg/m³ due to difference in dissolved CO_2 concentration, while $\Delta \rho$ was only 2.847-2.910 kg/m³ due to differences in temperature.

In Lindeberg and Wessel-Berg's model the water column is stable if only thermal gradients are considered, Ra lies then in the range 3.53-29.3 and no convection occurs ($Ra_c = 39.5$ in this case). However, if molecular CO_2 diffusion is considered (resulting in a significantly higher $\Delta \rho$), convection is predicted to occur. Lindeberg and Wessel-Berg define a stability criterion S which is the sum of the temperature and concentration effect on convective stability. S is analogous to Ra , and for infinite CO_2 dilution or an infinite molecular diffusion coefficient S becomes equal to Ra . The computed S values range from 1046-24204, and they are much higher than Ra_c . This means that convection will occur in aquifers under CCS conditions, which strongly enhances dissolution trapping and storage security. This convection is caused by the concentration gradient, not the temperature gradient. Lindeberg and Wessel-Berg (1997) suggest improvements for their model, especially a more sophisticated description of the concentration gradient should be implemented (they used a linearized concentration gradient). Moreover the Soret effect should be considered.

6.3 The Riaz model (2006)

Riaz et al. (2006) conducted a linear analysis and numerical simulations of the stability of the diffusive boundary layer (i.e. the brine layer adjacent to the scCO_2 phase into which the CO_2 diffuses) in a semi-infinite domain. Their calculations are based on Boussinesq-flow in a horizontal porous layer. The model neglects dispersion and geochemical reactions and assumes a homogenous and isotropic porous medium. Riaz et al. (2006) describe a critical time t_c (equation 28) which is a criterion for the onset of gravitational instability. For times larger than t_c convection will occur.

$$t_c = 146 \frac{\phi \mu^2 D}{(K \Delta \rho g)^2} \quad (28)$$

t_c influences the penetration depth of the diffusive boundary layer $\delta(t)$, which again influences $Ra_c = Ra(\delta(t))$. t_c can vary over several magnitudes, mainly because permeability can span several magnitudes in a geological formation. For a permeability increase from 1 mD to 3 Darcy Riaz et al. (2006) calculated a t_c decrease from 2000 years to below 10 days, and associated with that $\delta(t)$ changed from 55 m (for 1 mD permeability) to 0.07 m.

The model also demonstrates that Ra has a strong influence on the finger-like flow in the reservoir, including finger thickness and shape. In terms of the numerical model they found that grid size plays an important role and a fine grid is required to resolve disturbances at small times. Correct identification of such early disturbances is necessary to obtain reliable results.

One important conclusion they make is that dissolution trapping is strongly enhanced in high permeability reservoirs. They estimate that the onset of gravitational instabilities - essentially induced by molecular diffusion mass transfer processes - occurs after several hundreds of years for typical aquifers with average permeability. It should be noted that these estimates are quite rough because of the assumptions made.

6.4 Summary of reservoir models

There are other reservoir models described in the literature, e.g. Ennis-King and Paterson (2005) conclude that anisotropy of the reservoir has a strong effect on dissolution trapping, but this is beyond the scope of this book chapter and the reader is encouraged to check current research; this is an active area of research.

The simulation results are very important for CCS assessments and project planning, but it must be emphasized that more experimental research should be conducted, in the laboratory and especially at field scale to evaluate the quality of the model predictions. In addition, it is important to stress the approximative character of these models, real field situations are much more complex, e.g. it is not clear whether Fick's law can describe diffusion in the field or whether very heterogeneous pore structures (for instance in carbonate reservoirs) enhance convection or slow it down.

7. Multiphase flow in the reservoir – flow of the scCO₂ phase

The flow of the scCO₂ phase affects the dissolution process as it determines interfacial areas and overall position of the CO₂ in the reservoir. Reservoir models predict that the injected CO₂ phase rises upwards and is stopped by the caprock (Qi et al. 2009, Juanes 2006, Hesse 2008). This behaviour has been confirmed experimentally in the Sleipner formation by seismic imaging (Iglauer 2011).

Small residual CO₂ clusters at the trailing edge of the rising CO₂ plume - trapped by capillary forces (Iglauer et al. 2010, Juanes et al. 2006) - strongly increase CO₂-brine interfacial areas. Hence CO₂ dissolution speed is predicted to be accelerated, especially if combined with convective flow of saturated/undersaturated brine. However experimental reservoir monitoring data is needed to confirm these predictions. Optimal conditions would be to bring undersaturated brine continuously into contact with residual micrometer-sized CO₂ bubbles while removing saturated or highly CO₂-enriched brine simultaneously. Engineering this dissolution phenomenon can be a promising topic for future research.

Moreover, and most likely even more significant in the short term - thereby strongly affecting the economics of CCS schemes are the fluid dynamics associated with CO₂ injection. CO₂ injectivity and CO₂-wellbore effects can strongly impact CCS schemes. For

example, flow in the reservoir is strongly influenced by changes in rock morphology and wettability, which can result in changes of relative permeabilities and capillary pressures of CO₂ and brine. Relative permeability and capillary pressures however strongly influence multi-phase fluid flow in the reservoir. As an example, there is evidence that wettability (Espinoza and Santamaria 2010, Chiquet et al. 2007) and rock pore morphology – especially carbonates (Luquot and Gouze 2009) are changed by scCO₂. More research work is required in this area to completely understand these changes and improve CCS risk assessment.

8. Conclusions

In summary it is clear that dissolution trapping is a potential solution for storing large quantities of anthropogenic CO₂ thereby reducing carbon emissions. More research is required, especially field testing with integrated monitoring to check how the CO₂ behaves under realistic injection and reservoir conditions in the medium-to-long term. The major advantages of dissolution trapping are that very substantial amounts of CO₂ can be stored very safely. The risk is that CO₂ dissolves too slowly so that a significant part of CO₂ is still in a mobile separate supercritical phase (separated from the brine phase) which is buoyant and could escape to the surface. There are however two other CCS mechanisms, structural and residual trapping which prevent or at least reduce the CO₂ leakage risk. It must also be guaranteed that no drinkable-water aquifers are contaminated with CO₂ or any harmful species mobilized by CO₂ injection (e.g. dissolution of heavy metal ions by the acidic brine generated), which may then be transported into drinking water reservoirs.

9. Acknowledgements

I would like to thank Prof. Martin Blunt, Prof. Tetsuya Suekane, Dr. Amer Syed and Prof. Abbas Firoozabadi for reviewing this book chapter and helpful comments. Many thanks go to Prof. Zhenhao Duan for supplying the CO₂ solubility calculator software.

10. Nomenclature

CCS	carbon capture and storage (of carbon dioxide)
CO ₂	carbon dioxide
HCO ₃ ⁻	hydrogen carbonate anion
CO ₃ ²⁻	carbonate anion
H ⁺	proton
HC	hydrocarbon
a	year
Gt	Gigatons = 10 ⁹ tons
φ	porosity [-]
K	permeability [m ²]
M	molar mass [g/mol]
μ	viscosity [Pa.s]
ρ _{CO2}	CO ₂ density [kg/ m ³]
GOR	gas-oil ratio [m ³ /m ³]
H	domain depth; reservoir height [m]
D	diffusion coefficient [m ² /s]

$D_{\text{CO}_2\text{-H}_2\text{O}}$	diffusion coefficient of CO_2 into brine [m^2/s]
$D_{\text{H}_2\text{O-CO}_2}$	diffusion coefficient of H_2O into scCO_2 [m^2/s]
$D_{\text{CO}_2\text{-HC}}$	diffusion coefficient of CO_2 into hydrocarbon [m^2/s]
$D_{\text{CO}_2\text{-C}_{10}}$	diffusion coefficient of CO_2 into n-decane [m^2/s]
p	pressure [Pa]
T	temperature [K]
μ	viscosity [Pa.s]
μ_{CO_2}	viscosity of CO_2 [Pa.s]
$\mu_{\text{H}_2\text{O}}$	viscosity of water [Pa.s]
ρ	density [kg/m^3]
ρ_{brine}	density of brine [kg/m^3]
$\rho_{\text{CO}_2,\text{brine}}$	density of CO_2 -enriched brine [kg/m^3]
$\Delta\rho$	density difference [kg/m^3]
g	gravitational constant [m/s^2]
y	molality [mol/kg]
R	universal gas constant [$\text{J}/\text{mol.K}$]
$\mu_{\text{CO}_2}^{1(0)}$	standard chemical potential of CO_2 [J/mol]
F	fugacity coefficient [-]
$\zeta_{\text{CO}_2\text{-Na-Cl}}$	interaction parameter between CO_2 and Na^+ and Cl^-
$\lambda_{\text{CO}_2\text{-Na}}$	interaction parameter between CO_2 and Na^+
H	reservoir height [m]
Z_D	normalized reservoir height [-]
$V_{\text{m,brine}}$	apparent molar volume of CO_2 in brine [m^3/mol]
V_{CO_2}	molar volume of CO_2 [m^3/mol]
Y_{CO_2}	dissolved CO_2 mass fraction [-]
t_D	dimensionless time [-]
t	time [s]
u	interstitial or pore flow velocity [m/s]
Ra	Rayleigh number [-]
Re	Reynolds number [-]
Sh	Sherwood number [-]
Sc	Schmidt number [-]
z	depth [m]
α	dispersivity [m]
α_L	longitudinal dispersivity [m]
D_{dis}	dispersion coefficient [m^2/s]
C_0	surface concentration of droplet [mol/L]
C_∞	concentration at infinity [mol/L]
A	surface area of CO_2 droplet [m^2]
k	mass transfer coefficient [m/s]
j	volumetric flux [m/s]

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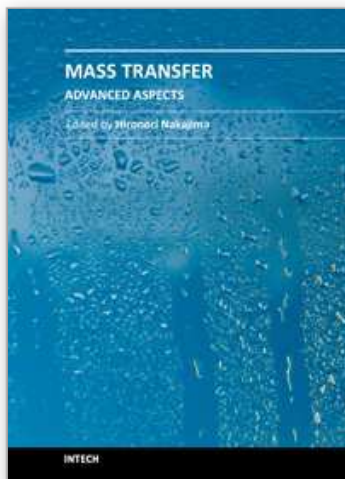
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Mass Transfer - Advanced Aspects

Edited by Dr. Hironori Nakajima

ISBN 978-953-307-636-2

Hard cover, 824 pages

Publisher InTech

Published online 07, July, 2011

Published in print edition July, 2011

Our knowledge of mass transfer processes has been extended and applied to various fields of science and engineering including industrial and manufacturing processes in recent years. Since mass transfer is a primordial phenomenon, it plays a key role in the scientific researches and fields of mechanical, energy, environmental, materials, bio, and chemical engineering. In this book, energetic authors provide present advances in scientific findings and technologies, and develop new theoretical models concerning mass transfer. This book brings valuable references for researchers and engineers working in the variety of mass transfer sciences and related fields. Since the constitutive topics cover the advances in broad research areas, the topics will be mutually stimulus and informative to the researchers and engineers in different areas.

How to reference

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Stefan Iglauder (2011). Dissolution Trapping of Carbon Dioxide in Reservoir Formation Brine – A Carbon Storage Mechanism, Mass Transfer - Advanced Aspects, Dr. Hironori Nakajima (Ed.), ISBN: 978-953-307-636-2, InTech, Available from: <http://www.intechopen.com/books/mass-transfer-advanced-aspects/dissolution-trapping-of-carbon-dioxide-in-reservoir-formation-brine-a-carbon-storage-mechanism>

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