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Heat Transfer and Phase Change in Deep CO₂ Injector for CO₂ Geological Storage

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1. Introduction

CO₂ capture and storage (CCS) is expected to reduce CO₂ emissions into the atmosphere. Various underground reservoirs and layers exist where CO₂ may be stored such as aquifers, depleted oil and gas reservoirs as well as unmined coal seams.

Coal seams are feasible for CCS because coal can adsorb CO₂ gas with roughly twice volume compared with CH₄ gas originally stored (Yee et al., 1993). However, the coal matrix is swelling with adsorbing CO₂ and its permeability is reduced. Supercritical CO₂ has a higher injection rate of CO₂ into coal seams than liquid CO₂ because its viscosity is 40% lower than the liquid CO₂ (see Harpalani and Chen, 1993).

The Japanese consortium carried out the test project on Enhanced Coal Bed Methane Recovery by CO₂ injection (CO₂-ECBMR) at Yubari City, Hokkaido, Japan during 2004 to 2007 [Yamaguchi et al. (2007), Fujioka et al.(2010)]. The target coal seam at Yubari was located about 890 to 900 m below the surface (Yasunami et al., 2010). However, liquid CO₂ was injected from the bottom holes because of heat loss along the deep injection tubing. The absolute pressure and temperature at the bottom hole was approximately 15.5MPa and 28°C. The regular tubing was replaced with thermally insulated tubing that included an argon gas layer but the temperature at the bottom was still lower than the critical temperature of CO₂.

This chapter provides a numerical model of heat transfer and calculation procedure for the prediction of CO₂ temperature and pressure that includes a phase change (supercritical or liquid) by considering the heat loss from the injector to surrounding casing pipes and rock formation. Furthermore, this study provides numerical simulation results of the temperature distribution of the coal seam after the injection of CO₂.

2. Prediction model for CO₂ injection temperature

2.1 CO₂ flow rate injected into a reservoir

As shown in Fig. 1, a schematic radial flow model in a reservoir, such as coal seam or aquifer, is targeted for CO₂ injection with vertical injection well (injector). The reservoir with radius R and thickness h_R , is saturated with water and open with constant pressure at its outer boundary. Assume omitting well pressure loss, the initial CO₂ mass flow rate, $M(0)$, at time $t = 0$, that is injected into the reservoir from its bottom hole, is equal to radial water flow rate in the reservoir [Michael et al. (2008) and Sasaki & Akibayashi (1999)],

$$M(0) = \rho_{BH} \frac{P_{BH}(0) - P_R}{\frac{\mu_w}{2\pi K_w h_R} \cdot \ln\left(\frac{R}{r_w}\right)} ; \quad P_{BH}(0) = P_{WH}(0) + g \int_0^H \rho(x,0) dx \quad (1)$$

where $\rho(x,t)$ and $\rho_{BH} = \rho(H,t)$ are CO_2 density in the injector and bottom hole respectively, g is acceleration of gravity, r_w is outer radius of the bottom hole, K_w is reservoir permeability, P_{WH} , P_{BH} and P_R are pressures at well head, bottom hole and injector outer boundary, μ_w is water viscosity in the reservoir, and H is length of vertical injector. The reservoir initial pressure is also equal to P_R .

After starting CO_2 injection, the CO_2 mass flow rate $M(t)$ and bottom hole pressure $P_{BH}(t)$ are changing with elapsed time t , since bottom hole pressure depends on CO_2 density distribution through the injector and water is replaced with CO_2 . Therefore, flow rate after becoming steady-state Q is given with P_{BH} and CO_2 viscosity μ_f at $t = \infty$.

$$M(\infty) = \rho_{BH}(\infty) \frac{P_{BH}(\infty) - P_R}{\frac{\mu_f}{2\pi K_w h_R} \cdot \ln\left(\frac{R}{r_w}\right)} ; \quad P_{BH}(\infty) = P_{WH}(\infty) + g \int_0^H \rho(x,\infty) dx \quad (2)$$

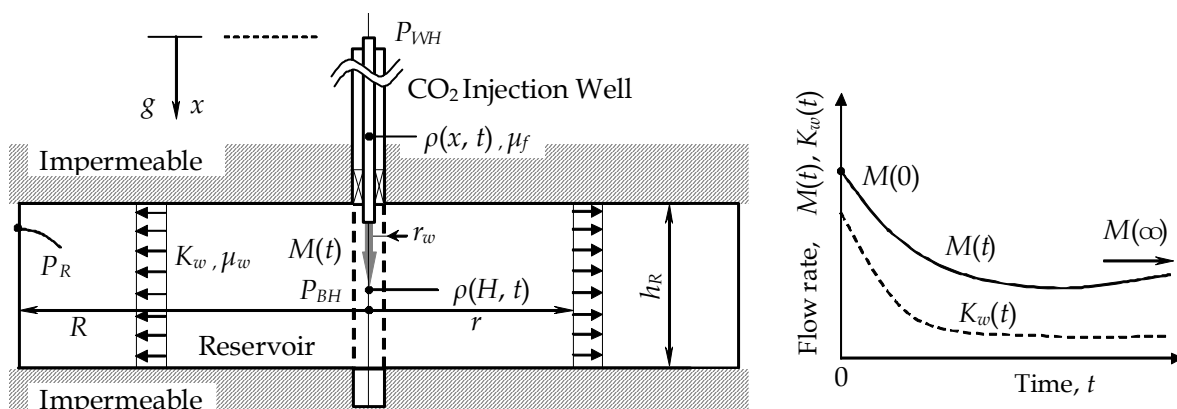


Fig. 1. Schematic radial flow model for injected CO_2 into a reservoir filled with water

Generally, CO_2 viscosity (30°C, 15MPa) is much smaller than water (roughly 1/30), thus the flow rate increases with t . Furthermore, viscosity of supercritical CO_2 is smaller than liquid CO_2 . On the other hand, the flow rate Q strongly depends on reservoir permeability times height ($=K_w h_R$). Especially coal seams have relatively low permeability of order 10^{-15} m^2 . It has been reported by some projects that permeability of coal seams decreased with rough ratio of 1/10 to 1/100 after CO_2 injection due to swelling of coal matrix by CO_2 adsorption [Clarkson et al. (2008) and Sasaki et al. (2009)].

2.2 Unsteady heat conduction equation

Figure 2 shows schematic diagram of radial heat loss from a vertical injection well (injector) that is consisting tubing pipe, casing pipes and well annulus. CO_2 is flowed down through the tubing pipe, and injected from bottom of the well with perforated holes. The annulus between two coaxial pipes is not used for CO_2 injection, and possibly needed to prevent heat loss from the tubing.

In present analytical approaches, inside area of the casing pipe is assumed as quasi-steady and outer region of the casing pipe ($r \geq r_{cao}$) is analyzed by unsteady equation of heat conduction. For the outer cement and rock region at a level, Fourier’s second law in cylindrical coordinates (r, x) is expressed as;

$$\frac{\partial \theta}{\partial t} = a_r \left(\frac{\partial^2 \theta}{\partial r^2} + \frac{1}{r} \frac{\partial \theta}{\partial r} + \frac{\partial^2 \theta}{\partial x^2} \right) \cong a_r \left(\frac{\partial^2 \theta}{\partial r^2} + \frac{1}{r} \frac{\partial \theta}{\partial r} \right) \tag{3}$$

where θ (°C) is rock temperature, t (s) is elapsed time, r (m) is radius, a_r (m²/s) is the heat diffusivity of rock. Heat conduction in vertical direction, x , can be omitted by comparing with that of radial direction. Analytical solution has been presented by Starfield & Bleloch (1983) for unsteady-state rock temperature distribution around underground airways. Especially, they presented a method to simulate internal surface temperature using with Biot number and elapsed time factor function of Fourier number (see section 2.7).

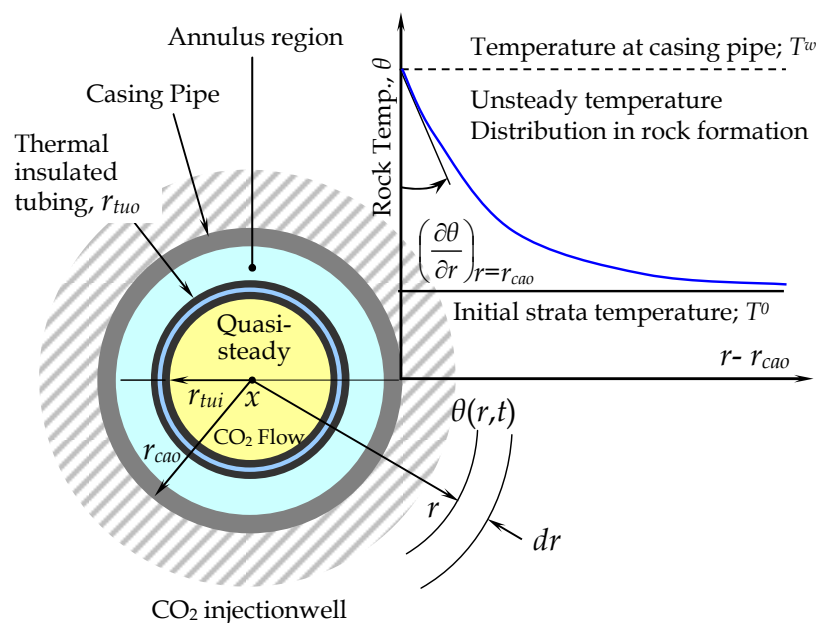


Fig. 2. Schematic diagram of radial heat flow from a vertical injection well (cross section)

2.3 Four thermal phenomena considered along CO₂ injection well

Figure 3 shows a schematic of heat transfer phenomena at an injection well. Four thermal phenomena were considered for the construction of the numerical model that is used for predicting CO₂ temperature and pressure at the bottom hole.

1. Natural convection in the annulus, filled with N₂ or water, increases heat transfer from tubing to casing, cement and rock formation. The heat transfer coefficient or Nusselt number at a specific depth is determined by using a formula reported by Choukairy et al. (2004).
2. The thermal performance of insulated tubing containing an argon shield layer was evaluated by considering the vertical convection flow of argon, thermal radiation between inner surfaces of the argon layer and thermal conduction at the tubing joints. Thermal characteristics of the insulated tubing are able to be corrected against the

original heat conductivity of argon gas using a number n determined by a field test and also by well logging data (see section 2.5).

3. The CO₂ phase was determined by its specific enthalpy which can be calculated from the pressure, temperature and heat loss along the well.
4. An unsteady analytical solution of the outer-surface temperature of casing pipe, expressed with Eq.(1), can be applied against the elapsed time from the start of CO₂ injection.

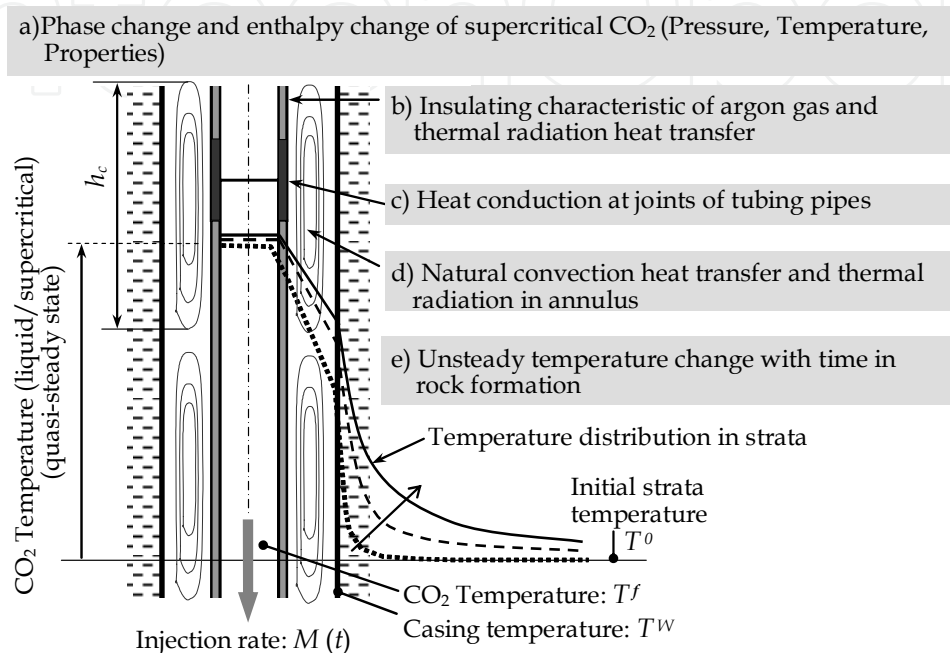


Fig. 3. Heat transfer phenomena from fluid flow in injector to surrounding rock formation

2.4 Overall thermal conductivity of the quasi-steady state region of the injection well

Figure 4 shows an example of the well structure (Yubari CO₂-ECBMR pilot-test site). CO₂ heat loss occurs during flow down to the bottom and propagates through various cylindrical combinations of steels and fluids with various thermal properties in the well configuration. To evaluate heat loss the overall heat conductivity that consists of conductivities of well materials and convective heat transfer rates of fluid flows that are contained in the well are important. Equations (4) and (5) represent single tubing and thermally insulated tubing, respectively (Nag, 2006).

$$\lambda = \frac{1}{\frac{\ln \frac{r_{cao}}{r_{cai}}}{\lambda_{Steel}} + \frac{\ln \frac{r_{cai}}{r_{ruo}}}{N_u \cdot \lambda_f} + \frac{\ln \frac{r_{tuo}}{r_{tui}}}{\lambda_{Steel}} + \frac{1}{r_{tui} \alpha_i}} \quad (4)$$

$$\lambda = \frac{1}{\frac{\ln \frac{r_{cao}}{r_{cai}}}{\lambda_{Steel}} + \frac{\ln \frac{r_{thco}}{r_{thci}}}{N_u \cdot \lambda_f} + \frac{\ln \frac{r_{thco}}{r_{thci}}}{\lambda_{Steel}} + \frac{\ln \frac{r_{thci}}{r_{tho}}}{n \cdot \lambda_{Ar}} + \frac{\ln \frac{r_{tho}}{r_{thi}}}{\lambda_{Steel}} + \frac{1}{r_{thi} \alpha_{thi}}} \quad (5)$$

Where α_{thi} is the heat transfer coefficient at the inner wall of the tubing pipe, λ_f is the heat conductivity of the fluid (water) in the annulus, λ_{Steel} is the heat conductivity of the casing and tubing pipes, N_u ($=\alpha_f r_{thco}/\lambda_f$) is the Nusselt number for the annulus and n is a correction number to adjust the heat conductivity of the argon gas layer in the insulated tubing.

2.5 Evaluation of performance of thermal insulated tubing

Thermal insulated tubing pipe is sometime used for geo-thermal wells through cold formation in order to prevent heat loss from produced hot spring water/steam. In case of the Yubari injected CO₂-ECBMR test, connected thermal insulated tubing pipes 20 m in length were used partially in 2005-2006 and totally in 2007. The insulated tubing includes argon gas shield layer is enclosed between inner and outer pipes to prevent heat loss from inside ideally with low thermal conductivity of argon gas; 0.116 W/m°C. However, joints between pipes are not shielded, and natural gas convection flow in the shield is expected to make increase the heat loss trasfered from the flow to outer tubing.



Fig. 4. Test to evaluate of equivalent thermal conductivity in the thermal insulated tubing using by pulsed heating carried at Yubari CO₂-ECBMR test field (Oct. 10, 2006) (see Yasunami et al., 2010)

To evaluate the thermal performance of the insulated tubing, tests using a insulated tubing pipe were carried out by pulsed heating from inside and measurements of outer and inner surface temperatures of the pipe placed horizontally as shown Fig. 4. Furthermore, the equivalent thermal conductivity was analyzed with Choukairy et al.'s equation (see section 2.5) and the history matching study for the well logging data. The thermal conductivity correction factor for conductivity of argon gas, n , is evaluated as shown in Fig. 5.

The equivalent heat conductivity including inside convective heat transfer was evaluated as three times larger as that of original argon gas without longitudinal heat loss through to connected tubing pipes. The correction factor, n , was introduced to adjust the equivalent heat conductivity of the tubing based on the original heat conductivity of argon gas. It was determined to be $n = 3$ but heat loss through the joints that are between the insulated tubing was not included in the test. The thermal equivalent conductivity of the insulated tubing was determined to be $n = 4$ or $\lambda = 0.21\text{W/m}^\circ\text{C}$ based on the well logging temperature at the Yubari CO₂-ECBMR test site and the measurement data were obtained from the heater response test carried out in the test field.

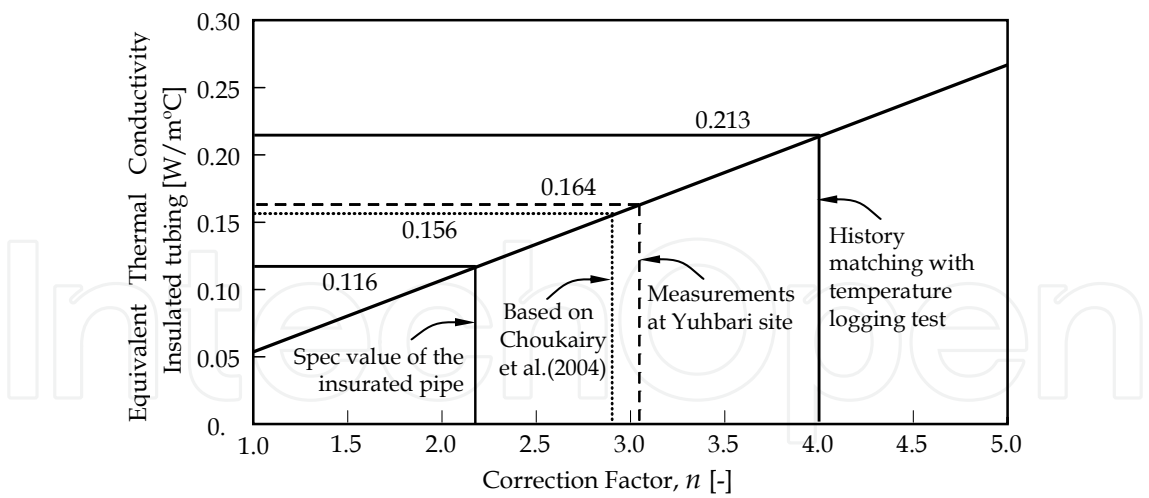


Fig. 5. Thermal conductivity correction factor for shielding with argon gas (*; spec value provided by a steel pipe maker)

2.6 Convective heat transfer in the annulus

Natural convection of annulus fluids makes influences on the heat transfer rate from the tubing pipe to the surrounding casing pipe and the formation. Choukairy et al. (2004) presented the following formula for the Nusselt number, N_u , for natural convection flow in an annulus with various radius ratios:

$$N_u = \frac{\alpha_f L}{\lambda_f} = \frac{\kappa}{mA} (P_r R_a)^{1/4} \cdot T_m^{5/4} \tag{6}$$

where α_f denotes the natural convection heat transfer coefficient on the inner surface of the casing, $L (=r_{cai} - r_{tuo})$ is the width of the annulus, κ is the radius ratio, m is a constant defined by Choukairy et al., A is the aspect ratio, P_r is the Prandtl number, R_a is the Rayleigh number and T_m is a dimensionless temperature defined by following equations:

$$A = \frac{h}{L} \tag{7}$$

$$\kappa = \frac{r_{cai}}{r_{tuo}} \tag{8}$$

$$R_a = \frac{g\beta_T(T^f - T^w) \cdot h_c^3}{a_f \cdot \nu_f} \tag{9}$$

$$T_m = \frac{1}{1 - \kappa^{4/5}} \tag{10}$$

where h_c is the circulation height of natural convection flow, g is the acceleration of gravity, β_T is the coefficient of thermal expansion of the fluid, ν_f is the dynamic viscosity and a_f is heat diffusivity of fluid in the annulus. The Nusselt number, N_u , calculated by Eq.(4) was used for each elevation.

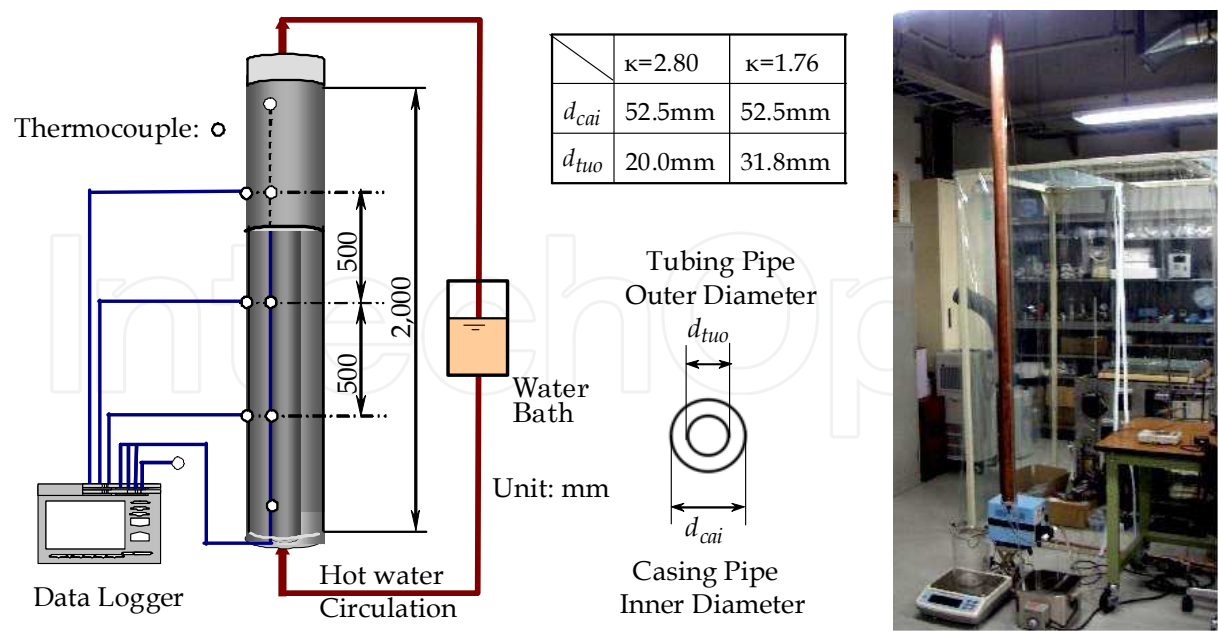


Fig. 6. Experimental setup to verify natural convection heat transfer coefficient in the annulus (Yasunami et al., 2010)

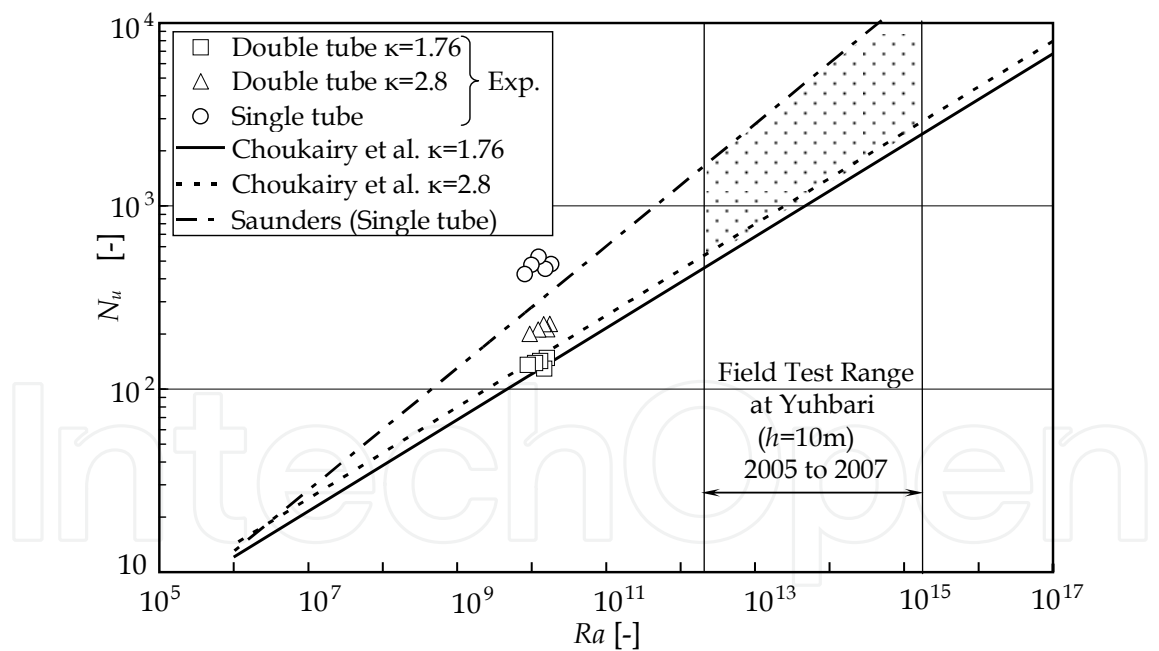


Fig. 7. Experimental results of Nusselt number for convective heat transfer in annulus (Yasunami et al., 2010)

Laboratory experiments were carried out to verify the reliability of Choukairy’s equation and to investigate the heat transfer rate using the well models consisting of two copper pipes with different diameters as shown in Fig. 6. Hot water at 40 to 60 °C was circulated through the inner pipe instead of CO₂. Pipe temperatures were measured by T-thermocouples that were placed on the pipe surfaces. Figure 7 shows experimental results

obtained for N_u and compared with those from Choukairy's equation. In addition, measured values of N_u on the outer surface of the single tubing determined using the equation proposed by Saunders [after Rohsenow et al. (1998)] were also compared in Fig. 7. Based on these results, we have found that Choukairy's equation is able to evaluate the heat transfer rate in the annulus.

2.7 Unsteady casing temperature

On the other hand, the temperature of the formation outside the casing pipe (outer surface) increases gradually after the injection. Assume T^0 is the initial strata formation temperature and T^{an} is the temperature in the annulus, the temperature at outer surface of the casing T^w , can be given by the solution for the unsteady heat conduction equation; Eq. (3). It has been presented by Starfield and Bleloch (1983):

$$T^w = T^{an} + \frac{\eta_t}{\eta_t + B_i} (T^0 - T^{an}) \quad (11)$$

where η_t is defined as the elapsed time factor and B_i is the non-dimensional Biot number, and B_i is defined by following equation:

$$B_i = \frac{\alpha_{ca} \cdot r_{cao}}{\lambda_r} \quad (12)$$

where α_{ca} is the apparent heat transfer rate at the inner casing and λ_r is the heat conductivity of rock. Starfield and Bleloch (1983) reported equations for the elapsed time factor η_t , which is a function of the Fourier number, τ , and Sasaki & Dindiwe (2002) revised it for $\tau \leq 1.5$ as:

$$\tau \leq 1.5 \quad ; \quad \eta_t = \frac{1}{[0.9879 + 0.3281(\ln \tau) + 0.03064(\ln \tau)^2]} \quad (13)$$

$$1.5 \leq \tau \leq 10 \quad ; \quad \eta_t = \frac{1}{[0.979813 + 0.383760(\ln \tau)]} \quad (14)$$

$$10 \leq \tau \leq 100 \quad ; \quad \eta_t = \frac{1}{[0.839337 + 0.444718(\ln \tau)]} \quad (15)$$

$$1000 \leq \tau \quad ; \quad \eta_t = \frac{2\Lambda(1 - \Lambda - \Lambda^2 - \Lambda^3)}{0.57722} \quad (16)$$

$$\Lambda = \frac{0.57722}{[\ln(4\tau) - 1.15444]} \quad (17)$$

$$\tau = \frac{a_r t}{r_{cao}^2} \quad (18)$$

The elapsed time factor η_t vs. Fourier number τ calculated by equations (13) to (18), is presented in Fig. 8.

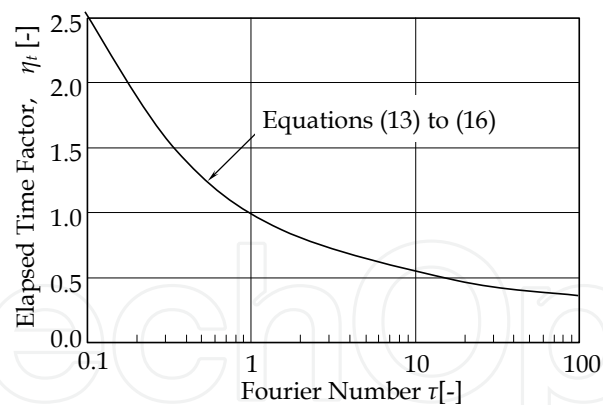


Fig. 8. Elapsed time factor vs. Fourier number

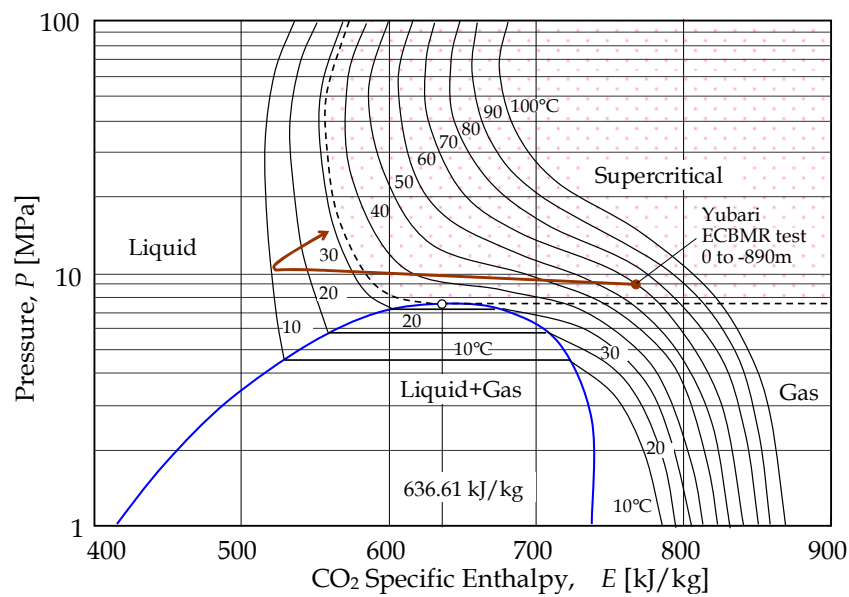


Fig. 9. CO₂ pressure-specific enthalpy and phase diagram calculated with PROPATH(2008) for 10 to 100 °C and 1 to 100 MPa (○; Critical point; 31.1°C and 7.38MPa)

2.8 Numerical equations for the determination of the CO₂ specific enthalpy

Changes in CO₂ temperature and phase (gas, liquid and supercritical) are accompanied by a specific enthalpy change. CO₂ specific enthalpy, $E(P,T)$ may be expressed by:

$$E(P,T) = \int_{T_0}^T C_p dT^f + \int_{P_0}^P V(1 - T_i \beta) dP \tag{19}$$

where V is the specific volume, β is the coefficient of thermal expansion, T_0 and P_0 are triple point temperature ($\approx -56.57^\circ\text{C}$) and pressure ($\approx 0.5185\text{MPa}$). The diagram CO₂ pressure-specific enthalpy for temperature range 10 to 100 °C and pressure range 1 to 100 MPa, that is calculated by PROPATH(2008), is shown in Fig. 9.

The specific enthalpy of CO₂ decreases with depth x by heat loss from CO₂ flow to the formation around the injection well.

$$\Delta q = 2\pi\lambda(T^w - T^f)\Delta x \tag{20}$$

In these numerical simulations, $E(P,T)$, $\theta(P,T)$ and $\rho(P,T)$ and other fluids properties are calculated using a corresponding software sub-routines, such as PROPATH(Propath Group, 2008) and NIST (2007). Calculation step $\Delta x_i = 1.0\text{m}$ can be used to get enough accuracy (Yasunami et al., 2010).

2.9 Required values in the numerical calculations

For these numerical calculations, three values for each depth are required.

1. Heat diffusivity of formation.
In Yubari ECBMR test project introduced in this book, no rock core drilling was carried out from 0 m to -800 m, thus we had to estimate rock properties (Fujioka et al., 2010). The heat conductivity λ_r and the heat diffusivity a_f of the rock formation outer casing have not been measured previously, so values of $a_f=1.30\times 10^{-6}\text{m}^2/\text{s}$ and $\lambda_r=1.30\text{ W/mK}$ were assumed and this was based on standard heat properties of sedimentary rocks (Yasunami et al., 2010).
2. Circulation height of natural convection flow in the annulus.
It was difficult to measure the circulation height h of natural convection in the annulus at the Yubari site. However, the bottom hole temperature was not sensitive to h , even when h changed from 5 to 20 m. Thus $h=10\text{m}$ was assumed as an appropriate value since natural convection was not observed at lower than 2m in the experiments described in the previous section.
3. Heat capacity of the tubing or casing.
We assumed that temperature changes of tubing and casing pipes were quasi-steady and thus the heat capacity of these pipes was not included in the equations.

3. Results of Yubari ECBMR test project

3.1 Injection well formation

Figure 11 shows a well structure and formation used at the Yubari CO₂-ECBMR test project in 2005. CO₂ heat loss occurs during flow down to the bottom and propagates through various cylindrical combinations of steels and fluids with various thermal properties in the well configuration. Table 1 shows conditions used for the models from 2005 to 2007 carried out in the project denoted as;

- a. Model 2005:
The well was drilled in 2005 (hereafter denoted as Model 2005) and consisted of thermally insulated tubing 180 m in length from the well head.
- b. Model 2006:
In 2006, thermally insulated tubing of 180m in length was used at the head (0 to 180m) and the bottom (650 to 890m) while the annulus was filled with liquid CO₂.
- c. Model 2007:
In 2007 all the injection pipe tubing was replaced with thermally insulated tubing of 890 m in length and H₂O was used to fill the annulus. This was done to minimize heat loss from the tubing and thus keep CO₂ in its supercritical condition.
- d. Heater Model 2007:
To overcome the difficulty of low temperature and low injection rate, numerical predictions were done considering the use of an electric line heater to heat up CO₂ flow at the position of 180m from the surface. The heater capacity of $W = 1.43\text{kW}$ was chosen because of the cable strength and restrictions of materials against corrosion of supercritical CO₂.

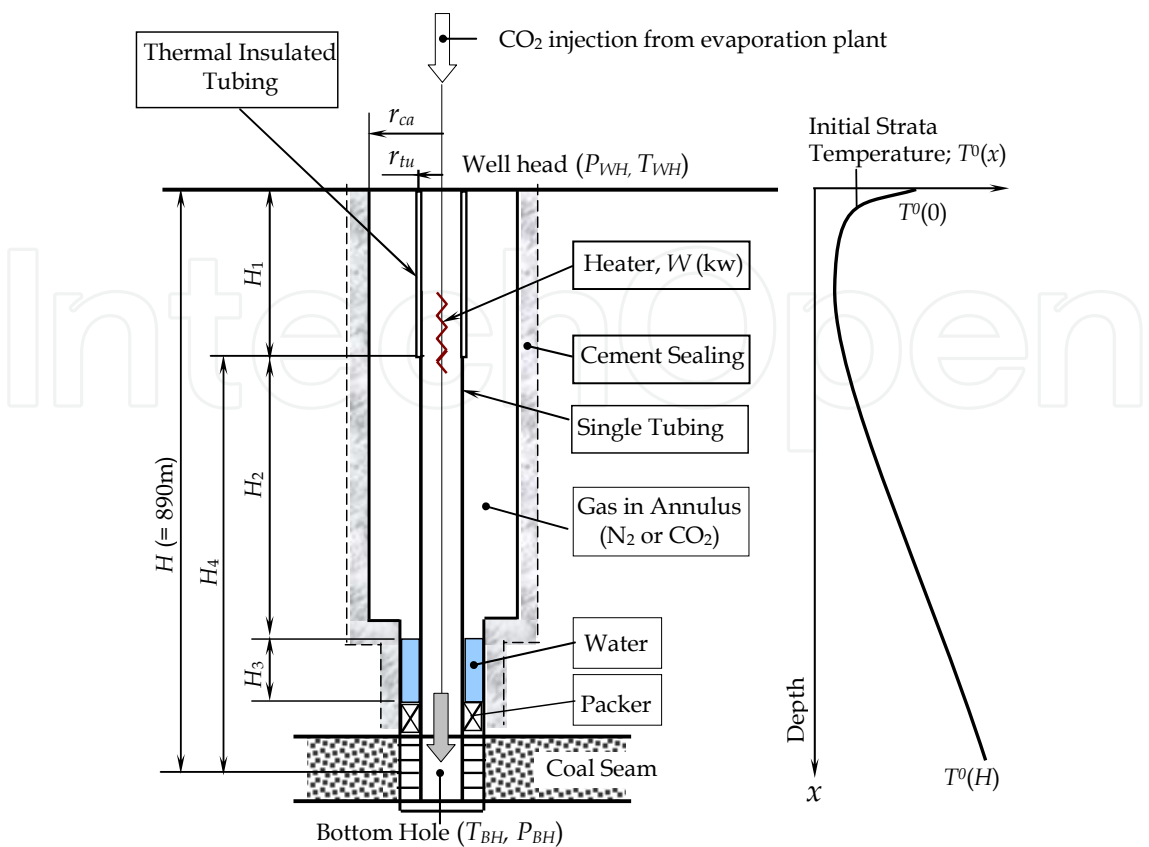


Fig. 11. An example of CO₂ injection well formation, initial strata temperature and schematic annulus formation (Yubari CO₂-ECBMR test project, Model 2005)

Term/Model		Model 2005	Model 2006	Model 2007	Model 2007+ Heater
Temp. and Press. At well head (T_{WH} , P_{WH})		70 °C 9.0 MPa	70 °C 9.0 MPa	70 °C 8.6 MPa	70 °C 8.6 MPa
S. Enthalpy at well head; E_{WH} (kJ/kg)		766.34	766.34	771.27	771.27
Injection rate M (kg/ day)		3000	3000	3000	11000
Tubing	H_1 ; $x=0\sim180$ m	Insulated tube	Insulated tube	Insulated tube	Insulated Tube
	H_2 ; $x=180\sim667$ m	Single tube	Single tube	Insulated tube	Insulated Tube
	H_3 ; $x=690\sim890$ m	Single tube	Insulated tube	Insulated tube	Insulated Tube
Fluid in annulus	for H_1	N ₂ Gas	Liquid CO ₂	Water	Water
	for H_2	N ₂ Gas	Liquid CO ₂	Water	Water
	for H_3	Water	Water	Water	Water
Heater output; W (kW) ($x=0\sim255$ m)		0	0	0	1,430

Table 1. Parameters of injection well formation for CO₂ injection

3.2 Effect of natural convection in the annulus

The results of temperature prediction against depth after one day by the heat conduction model ($N_{u}=1$ and $n=1$) and the heat convection model for an injection temperature of 70°C, an injection pressure of 9MPa and an injection rate of 3.0ton/day is shown in Fig. 12.

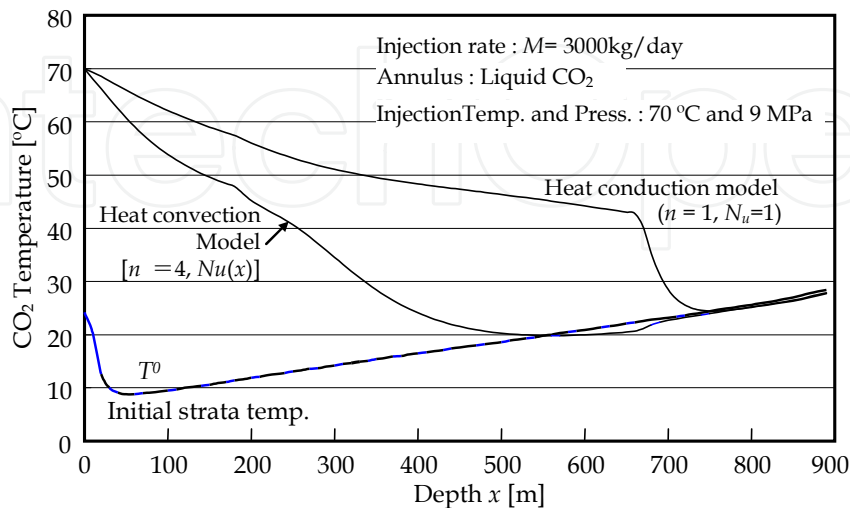


Fig. 12. A comparison of CO₂ temperatures after 1 day between heat conduction and heat convection models (Model 2005, 180m insulated tubing was partly used from well head)

Figure 13 shows comparisons of CO₂ temperature and pressure at the bottom hole. Numerical simulation results for the data, obtained in 2005 at the Yubari field, show that the heat convection model is better than the conduction model. This is because the temperature of the CO₂ decreased by heat loss caused by natural convection in the annulus. The bottom pressure increased because of the increase in CO₂ density that resulted from the CO₂ supercritical to liquid phase change.

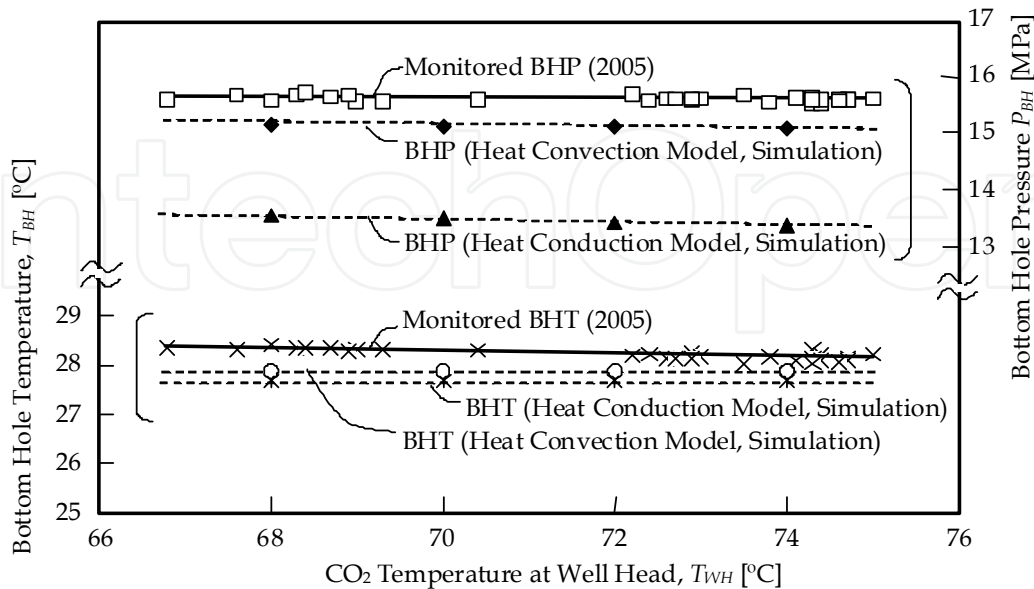


Fig. 13. Comparisons of bottom hole temperature and pressure (BHT and BHP) that were simulated by heat conduction and heat convection models with monitored values (Model 2005)

Figures 14 and 15 show comparisons of CO₂ temperature and pressure between simulations and the well logging data for injection conditions of 68.54°C, 9MPa and 4.5ton/day at the well head. Since logging from the surface to a level of -890m took 2.4hours (Prensky,1992), simulated CO₂ temperatures against depth were plotted for each time segment. The reason for the rise in the measured pressure near the well head of about of 0.3 MPa is unknown at present.

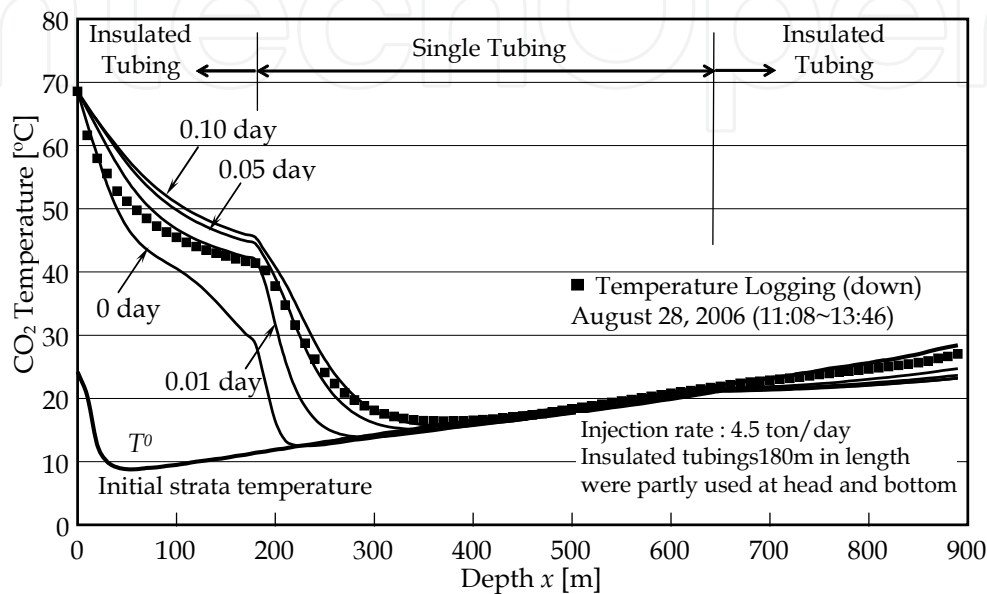


Fig. 14. Comparison between logged temperature and simulation results for Model 2006 (Logging data was obtained on August 28, 2006 (11:08 to 13:46) at the Yubari ECBMR test site)

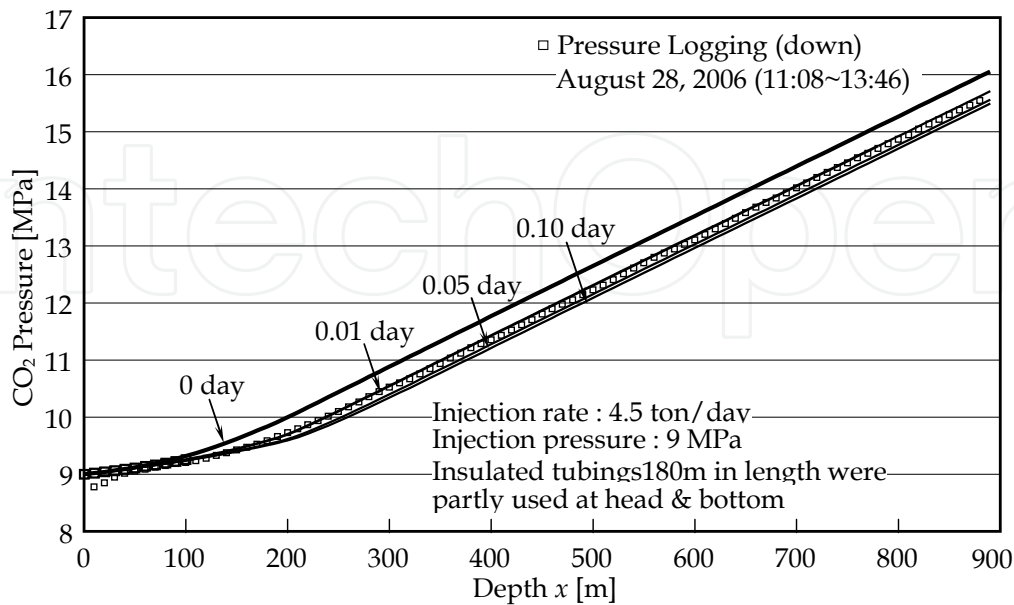


Fig. 15. Comparison between logged pressure and simulation results (Model 2006) (Logging data was obtained on August 28,2006 (11:08 to 13:46) at the Yubari test site).

3.3 Thermal insulated tubing partly used at the well head and bottom (Model 2006)

For the case of Model 2006 of Yubari ECBMR test, numerical simulations at 0, 22 and 68 days are shown in Figure 16. Figures 17 and 18 show the Nusselt number, the heat conductivity, the density and the specific enthalpy versus depth after 1 day.

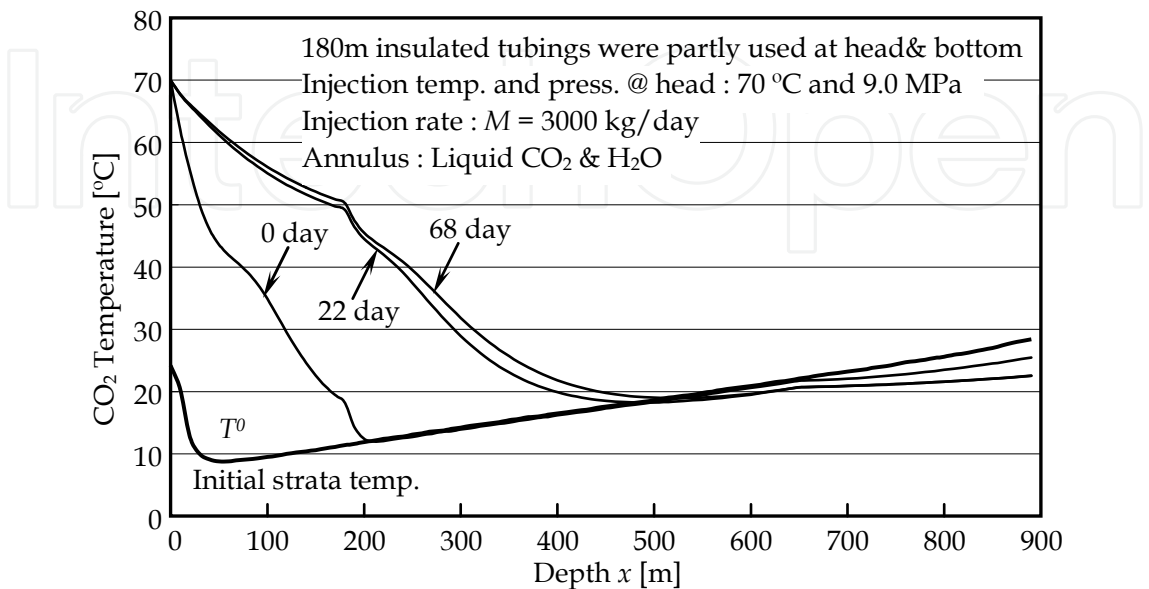


Fig. 16. CO₂ temperature distribution vs. depth (Model 2006).

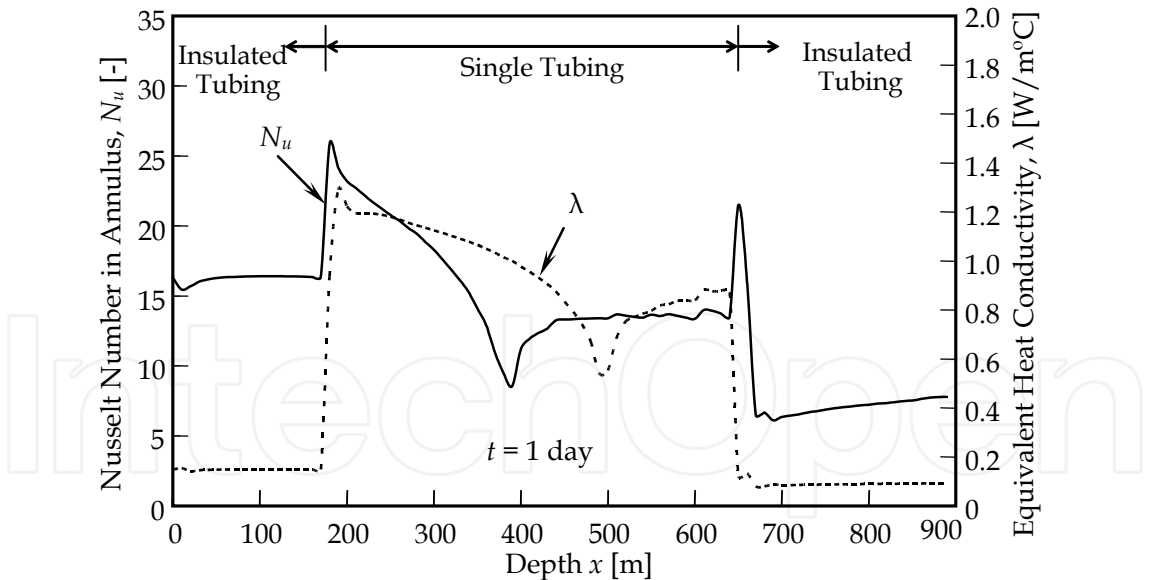


Fig. 17. Nusselt number of convective heat transfer in the annulus vs. depth (Model 2006).

The temperature was still lower than the supercritical temperature (=31.4°C) at the bottom hole because liquid CO₂ filled the annulus and cold CO₂ flow was maintained from 650 to 890m for the insulated tubing despite the formation temperature increasing with depth. The line in Fig. 9 shows s typical phase changes in the injection tubing on the CO₂ pressure-specific enthalpy diagram.

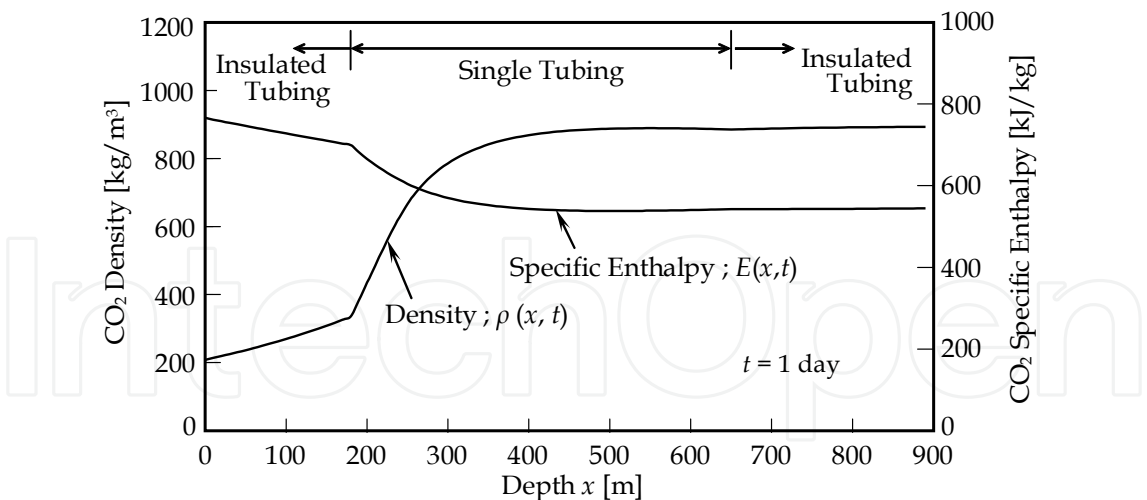


Fig. 18. CO₂ density and specific enthalpy vs. depth (Model 2006).

3.4 All usage of thermal insulated tubing (Model 2007)

In 2007, the all injection tubing pipe was replaced with thermally insulated tubing of 890 m in length and the annulus was filled with water. Figure 13 shows numerical calculation results for Model 2007. The predicted temperature for the bottom hole at an injection rate of 3.0ton/day is 26.0°C, which is lower than the observed temperature at the outer surface of the annulus of 27.5°C. This was influenced by the formation temperature in the annulus.

3.5 The effect of injection rate on the bottom hole temperature

It was expected that the temperature of CO₂ at the bottom hole would increase as the injection rate was increased, since heat loss is not sensitive to flow rate. Figure 19 shows a sensitivity analysis for temperature versus the injection rate at the bottom hole for Model 2007. The CO₂ phase was supercritical at the bottom hole after 1 day when the injection rate was over 12ton/day as shown in this figure. An operation like a hydraulic fracture is required to improve permeability, since the injection rate depends on the permeability around the injection well.

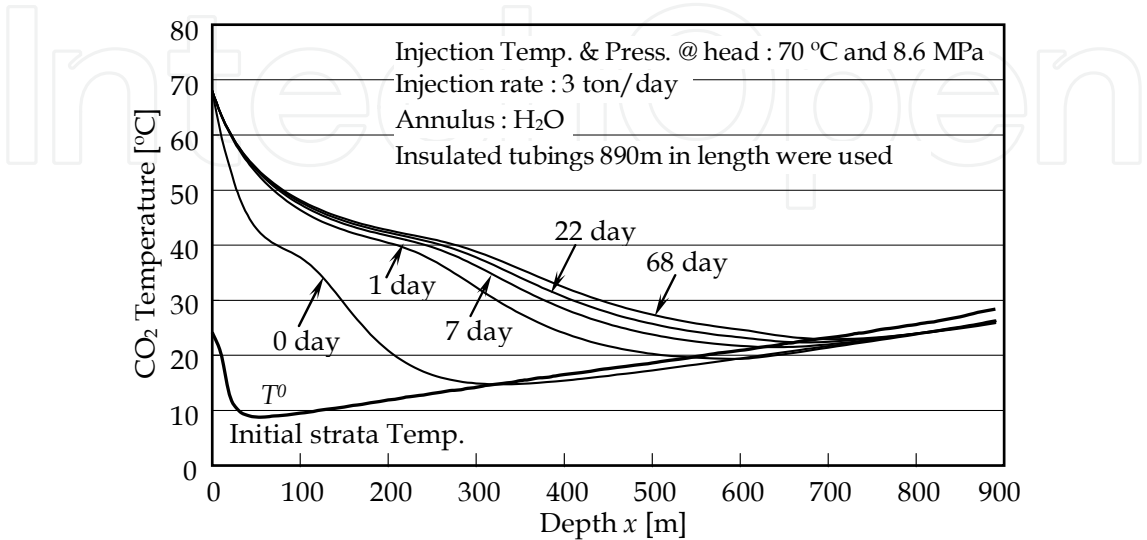


Fig. 19. CO₂ temperature distribution vs. depth for H=H1 (Model 2007)

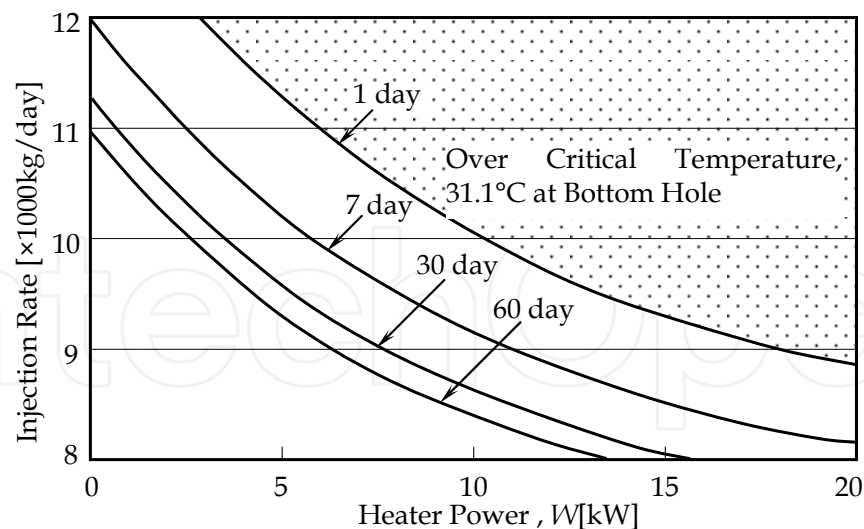


Fig. 20. Control map for maintaining the supercritical condition at the bottom hole (Model 2007+Heater, Injection temperature and pressure at well head; 70 °C and 8.6 MPa, Annulus; H₂O, Insulated tubing 890m in length was used)

3.6 Prediction of CO₂ temperature using a line heater

All tubing was replaced with thermally insulated tubing but the bottom hole temperature was still not adequate to maintain CO₂ in its critical condition. To overcome the difficulty of injection, numerical predictions were done considering the use of an electric line heater with 1.43kW of heating from the surface to 180m (denoted as Heater Model 2007). The heater capacity of 1.43kW was chosen because of the cable strength and because of restrictions of materials for supercritical CO₂. Table 1 shows conditions used in the calculation. The temperature at the bottom hole from Heater Model 2007 is 5°C higher than that from Model 2007. Even if the energy efficiency of CO₂ injection becomes lower by heating in the injector, it is better that CO₂ temperature is in a supercritical condition at the bottom hole to keep larger CO₂ injection rate into the coal seam. Figure 15 shows a control map for the CO₂ injection rate to maintain the supercritical condition at the bottom. This model shows that the critical temperature increases with the heater power and the elapsed time from the start of CO₂ injection.

4. Summary

In this chapter, a numerical model of heat transfer and calculation procedure for flow and heat transfer phenomena, related to CO₂ flow in a vertical deep injector, has been focused in order to predict CO₂ temperature, pressure and phase change (supercritical or liquid CO₂) in the well. Especially, it was considered that the heat loss from the injector to surrounding casing pipes and rock formation including natural convection heat transfer in annulus and insulated tubing pipes. Furthermore, numerical simulations have been presented for the Yubari CO₂-ECBMR test project carried out from 2005 to 2007.

The results are summarized as follows:

1. The bottom hole pressure and temperature in the injector at Yubari CO₂-ECBMR test field were successfully simulated by considering heat loss accelerated by natural convection flow in the annulus.

2. The thermal equivalent conductivity of the insulated tubing was determined to be $0.21\text{W/m}^\circ\text{C}$ based on the well logging temperature carried out at the Yubari test site.
3. A control map showing targeted injection rates against the heater power for elapsed time as a parameter was compiled to maintain the supercritical condition at the bottom hole of the injector.
4. CO_2 at the bottom hole is expected to be supercritical at a CO_2 injection rate over 12 ton/day without any heating or 11 ton/day using the 1.43kW line heater in the injector.

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6. Nomenclature

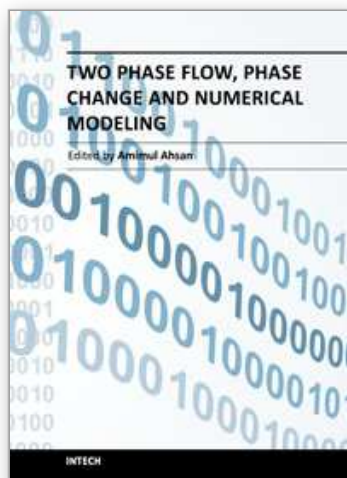
a_f	=	heat diffusivity of fluid [m^2/s]
a_r	=	heat diffusivity of rock [m^2/s]
A	=	aspect ratio [-]
B_i	=	Biot number [-]
C_p	=	isobaric specific heat [$\text{kJ}/(\text{kg}^\circ\text{C})$]
E	=	specific enthalpy [kJ/kg]
E_{WH}	=	specific enthalpy at well head ($x=0$) [kJ/kg]
f	=	friction factor of tubing pipe [-]
g	=	acceleration of gravity [m/s^2]
h_c	=	circulation height of natural convection flow [m]
h_R	=	reservoir height [m]
K_w	=	permeability of reservoir [m^2]
L	=	width of annulus [m]
m	=	constant number defined by Choukairy et al. (2004) [-]
M	=	CO_2 mass flow rate [kg/s]
n	=	correction to adjust heat conductivity in thermally insulated tubing [-]
N_u	=	Nusselt number [-]
P_{BH}	=	pressure at bottom hole [MPa]
P_R	=	pressure at reservoir outer boundary [MPa]
P_r	=	Prandtl number ($=\nu_f/a_f$) [-]
P_{WH}	=	pressure at well head ($x=0$) [MPa]
Q	=	CO_2 flow rate [m^3/s]
r_{cao}	=	outer radius of casing [m]
r_{cai}	=	internal radius of casing [m]
r_{tuo}	=	outer radius of single tubing [m]
r_{tui}	=	internal radius of single tubing [m]
r_{thco}	=	outer radius of outer thermal insulated tubing [m]
r_{thci}	=	internal radius of outer thermal insulated tubing [m]
r_{tho}	=	outer radius of inner thermal insulated tubing [m]
r_{thi}	=	internal radius of inner thermal insulated tubing [m]
R_a	=	Rayleigh number [-]

T^0	=	initial strata temperature [°C]
T_{an}	=	temperature of fluid in the annulus [°C]
T_{BH}	=	temperature at bottom hole [°C]
T^f	=	CO ₂ temperature in tubing [°C]
T_m	=	dimensionless temperature [-]
T_w	=	temperature at outer surface of casing [°C]
T_{WH}	=	temperature at well head ($x=0$) [°C]
v	=	average velocity of CO ₂ flow in tubing pipe [m/s]
V	=	specific volume [m ³ /kg]
x	=	length from surface (depth) [m]
α_f	=	heat transfer rate at inner surface of casing [W/(m ² °C)]
α_{thi}	=	heat transfer rate at inner surface of tubing [W/(m ² °C)]
α_{ca}	=	equivalent heat transfer rate at inner casing [W/(m ² °C)]
β	=	coefficient of thermal expansion of CO ₂ [1/K]
β_T	=	coefficient of thermal expansion of the fluid [1/K]
η_t	=	elapsed time factor [-]
κ	=	radius ratio [-]
λ	=	overall heat conductivity [W/(m°°C)]
λ_{Ar}	=	heat conductivity of fluid (Ar) in thermal insulated tubing [W/(m°°C)]
λ_f	=	heat conductivity of fluid (N ₂ , CO ₂ or water) in annulus [W/(m°°C)]
λ_r	=	heat conductivity of rock [W/(m°°C)]
λ_{Steal}	=	heat conductivity of casing and tubing [W/(m°°C)]
μ_f	=	CO ₂ viscosity [Pas]
μ_w	=	water viscosity [Pas]
ν_f	=	dynamic viscosity [m ² /s]
τ	=	Fourier number [-]
Δq	=	heat flow rate at tubing element [W]
ΔW	=	heat generated by a heater during x_i to x_{i+1} [W]
Δx	=	length of tubing element [m]

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The heat transfer and analysis on laser beam, evaporator coils, shell-and-tube condenser, two phase flow, nanofluids, complex fluids, and on phase change are significant issues in a design of wide range of industrial processes and devices. This book includes 25 advanced and revised contributions, and it covers mainly (1) numerical modeling of heat transfer, (2) two phase flow, (3) nanofluids, and (4) phase change. The first section introduces numerical modeling of heat transfer on particles in binary gas-solid fluidization bed, solidification phenomena, thermal approaches to laser damage, and temperature and velocity distribution. The second section covers density wave instability phenomena, gas and spray-water quenching, spray cooling, wettability effect, liquid film thickness, and thermosyphon loop. The third section includes nanofluids for heat transfer, nanofluids in minichannels, potential and engineering strategies on nanofluids, and heat transfer at nanoscale. The fourth section presents time-dependent melting and deformation processes of phase change material (PCM), thermal energy storage tanks using PCM, phase change in deep CO₂ injector, and thermal storage device of solar hot water system. The advanced idea and information described here will be fruitful for the readers to find a sustainable solution in an industrialized society.

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