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An Innovative Method for Robots Modeling and Simulation

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1. Introduction

The robot dynamics modeling and simulation problem has been studied for the last three decades intensively.

In particular, the forward dynamics problem of a robot is a very relevant issue, which there is still to say about in terms of efficient computation algorithms, that can be also simple to understand, to develop and to implement, above all for practical robots, robots with many links and/or with flexible links (Featherstone, 1987), (Featherstone & Orin, 2000), (Sciavicco & Siciliano, 2000), (Khalil & Dombre, 2002).

Indeed, in these cases the methods based on the classical Lagrange formulation give rise to an analytical model with numerous terms that may be difficult to use. The methods based on the Newton-Euler formulation are not very easy to apply and do not provide easily manageable analytical formulae, even if they are efficient from a computational point of view (Featherstone, 1987).

An important contribution to solve the previous problems is given (Celentano, 2006), (Celentano & Iervolino, 2006), (Celentano & Iervolino, 2007) by a new, simple and efficient methodology of analysis, valid for all of robots, that makes use of a mathematical model containing a lower number of algebraic terms and that allows computing, with a prescribed maximum error, the gradient of the kinetic energy starting from the numerical knowledge of the only inertia matrix rather than using, as usually found in the literature, complex analytical calculations of the closed-form expression of this matrix.

This result is very strong because it allows solving the forward dynamics problem of a robot in a simple and efficient manner, by analytically or numerically computing the inertia matrix and the potential energy gradient only.

Moreover, this method allows students, researchers and professionals, with no particular knowledge of mechanics, to easily model planar and spatial robots with practical links.

From this methodology follows also a simple and efficient algorithm for modeling flexible robots dividing the links into rigid sublinks interconnected by equivalent elastic joints and approximating and/or neglecting some terms related to the deformation variables.

In this chapter some of the main results stated in (Celentano, 2006), (Celentano & Iervolino, 2006), (Celentano & Iervolino, 2007) are reported.

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In details, in Section II the new integration scheme for robots modeling, based on the knowledge of the inertia matrix and of the potential energy only, is reported (Celentano & Iervolino, 2006).

In Section III, for planar robots with revolute joints, theorems can be introduced and demonstrated to provide a sufficiently simple and efficient method of expressing both the inertia matrix and the gradient of the kinetic energy in a closed and elegant analytical form. Moreover, the efficiency of the proposed method is compared to the efficiency of the Articulated-Body method, considered one of the most efficient Newtonian methods in the literature (Celentano & Iervolino, 2006).

In Section IV, for spatial robots with generic shape links and connected, for the sake of brevity, with spherical joints, several theorems are formulated and demonstrated in a simple manner and some algorithms that allow efficiently computing, analytically the inertia matrix, analytically or numerically the gradient of the kinetic and of the gravitational energy are provided. Furthermore, also in this case a comparison of the proposed method in terms of efficiency with the Articulated-Body one is reported (Celentano & Iervolino, 2007).

In Section V some elements of flexible robots modeling, that allow obtaining, quite simply, accurate and efficient, from a computational point of view, finite-dimensional models, are provided. Moreover, a significant example of implementation of the proposed results is presented (Celentano, 2007).

Finally, in Section VI some conclusions and future developments are reported.

2. A new formulation of the Euler-Lagrange equation

In the following a new formulation of the dynamic model of a robot mechanism in a more efficient form (for its analytical and/or numerical study) is presented.

It is well known that the usual form of the Euler-Lagrange equation for a generic robot with n degrees of freedom, which takes on the form (De Wit et al., 1997)

$$\frac{d}{dt} \frac{\partial T}{\partial \dot{q}} - \frac{\partial T}{\partial q} = \gamma + \eta + u_a + u, \quad (1)$$

where: q is the vector of the Lagrangian coordinates; T is the kinetic energy, given by

$T = \frac{1}{2} \dot{q}^T B(q) \dot{q}$, being B the robot inertia matrix; γ is the vector of the gravity forces, η is the vector of the elasticity forces and u_a is the non-conservative generalised forces (e.g. damping torques), which are usually function of q and \dot{q} only, and u is the vector of the actuation generalised forces, is typically rewritten as

$$B(q)\ddot{q} + C(q, \dot{q})\dot{q} = \gamma + \eta + u_a + u, \quad (2)$$

in which an expression of matrix C is the following:

$$C(q, \dot{q}) = \sum_{i=1}^n \frac{\partial B}{\partial q_i} \dot{q}_i - \frac{1}{2} \begin{bmatrix} \dot{q}^T \frac{\partial B}{\partial q_1} \\ \vdots \\ \dot{q}^T \frac{\partial B}{\partial q_n} \end{bmatrix}. \quad (3)$$

By setting $B = \{b_{ij}\}$, $C = \{c_{ij}\}$, it is also

$$c_{ij} = \sum_{k=1}^n \left(\frac{\partial b_{ij}}{\partial q_k} - \frac{1}{2} \frac{\partial b_{ik}}{\partial q_j} \right) \dot{q}_k. \quad (4)$$

Alternatively, an equivalent matrix $C_c = \{c_{cij}\}$, i.e. such that $C_c \dot{q} = C \dot{q}$, that makes uses of the Christoffell's symbols (Sciavicco & Siciliano, 2000), is given by

$$c_{cij} = \sum_{k=1}^n \frac{1}{2} \left(\frac{\partial b_{ij}}{\partial q_k} + \frac{\partial b_{ik}}{\partial q_j} - \frac{\partial b_{jk}}{\partial q_i} \right) \dot{q}_k. \quad (5)$$

There are various methods to calculate both $B(q)$ and $C(q, \dot{q})$. The simplest one, from a conceptual point of view, based on analytical expressions, is extremely complex and onerous, even if symbolic manipulation languages are employed. Other methods do require a more in-depth knowledge of mechanics (De Wit et al., 1997).

The next theorem provides an innovative, efficient and simple method for modeling and simulating a robot.

Theorem 1. If the Euler-Lagrange equation is rearranged as follows:

$$\frac{d}{dt}(B\dot{q}) = c + \gamma + \eta + u_a + u, \quad (6)$$

where c is the gradient of the kinetic energy with respect to q , namely

$$c = \frac{\partial T}{\partial q}, \quad (7)$$

then it can be efficiently integrated according to the innovative integration scheme of Fig. 1, where $\dot{q}_g = B\dot{q}$ and F is a suitable function that allows computing \dot{q} and \ddot{q}_g from u , q and \dot{q}_g .

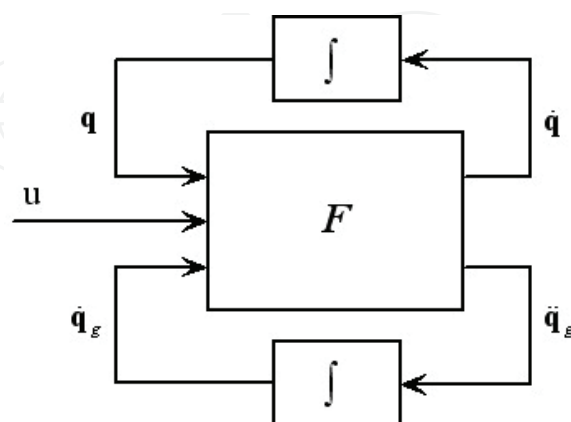


Fig. 1. Block diagram for the robot forward dynamics integration.

Proof. It is easy to verify that the algorithm that describe the function F of the block scheme of Fig. 1 is:

Step 1. Compute $B(q)$ by applying one of the classic methods proposed in the literature (see also Theorems 2 and 3 and (30) in Section 3).

Step 2. Compute \dot{q} through the relationship $\dot{q} = B^{-1}\dot{q}_g$.

Step 3. Compute $c(q, \dot{q})$ and u_a .

Step 4. Compute $\frac{d}{dt}(B\dot{q}) = \ddot{q}_g = c + \gamma + \eta + u_a + u$.

Remark 1. The proposed method allows avoiding the computation of the first term on the right hand-side of (4) or of the first two terms on the right-hand side of (5). Such saving is significant because the computational burden relative to this term is not negligible with respect to the one required by the calculation of $B^{-1}\dot{q}_g$.

This clearly emerges from Fig. 2, where the percentages of saved flops vs. the number of flops required to compute $B^{-1}\dot{q}_g$ are evaluated (in MATLAB™ environment) with reference to the coefficients in (4), for a considerable number of cases of robots with random and anyhow realistic links parameters.

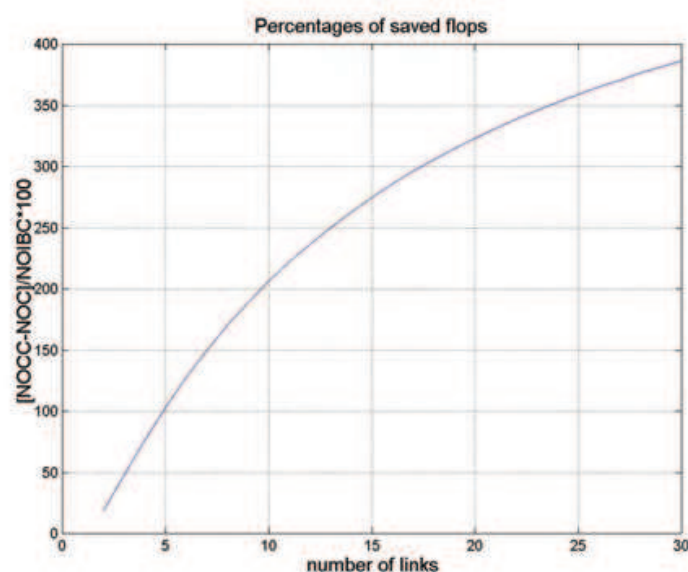


Fig. 2. Percentages of saved flops evaluated for robot models with random links parameters.

Remark 2. Since many efficient algorithms for the numerical computation of the matrix B are available in the literature, the gradient of the kinetic energy c can be computed in a very simple and accurate way numerically, instead of using a symbolic expression.

To this end, since the kinetic energy for the majority of the robots is a quadratic function respect to \dot{q} , whose coefficients are constant with respect to q or depend on q_i according to terms of the type $\sin(aq_i + \varphi)$, the following lemma is useful.

Lemma 1. The derivative of $\sin(aq_i)$ can be numerically computed by a two [three] points formula:

$$\frac{d}{dq_i} \sin(aq_i) = \frac{\sin[a(q_i + \Delta)] - \sin(aq_i)}{\Delta} \quad \left[\frac{d}{dq_i} \sin(aq_i) = \frac{\sin[a(q_i + \Delta)] - \sin[a(q_i - \Delta)]}{2\Delta} \right] \quad (8)$$

with error:

$$e_2 = -\frac{1}{2}a^2 \sin(a\chi)\Delta, \quad |e_2| \leq a^2 \frac{\Delta}{2}, \quad \chi \in (q_i, q_i + \Delta)$$

$$\left[e_3 = -\frac{1}{6}\Delta^2 a^3 \cos(a\chi), \quad |e_3| \leq a^3 \frac{\Delta^2}{6}, \quad \chi \in (q_i - \Delta, q_i + \Delta) \right]. \quad (9)$$

Proof. See any numerical computation book.

In view of Lemma 1 it is possible to compute, with a prescribed maximum error, the gradient of the kinetic energy starting from the numerical knowledge of the inertia matrix B rather than using, as usually found in literature, complex analytical calculations of the analytical expression of B . In details, for practical precision purposes, a good value of $\partial T / \partial q$ can be obtained using a two [three] points formula with $\Delta = 10^{-3} \div 10^{-4}$ [with $\Delta = 10^{-2} \div 10^{-3}$] and evaluating $n-1$ [$2(n-1)$] times the inertia matrix B numerically.

3. Planar robot modeling

In the case of planar robots with revolute joints, theorems can be introduced and demonstrated to provide a sufficiently simple and efficient method of expressing both the inertia matrix and the gradient of the kinetic energy in a closed and elegant analytical form.

3.1 Hypotheses and notations

In the following it is considered the case of a generalised planar robot constituted by n links, each of them with parallel rotation axes C^- , C^+ and center of mass G belonging to the plane containing the two relative rotation axes (see Fig. 3). In Fig. 4 a planar robot with three links and a horizontal work plane is shown.

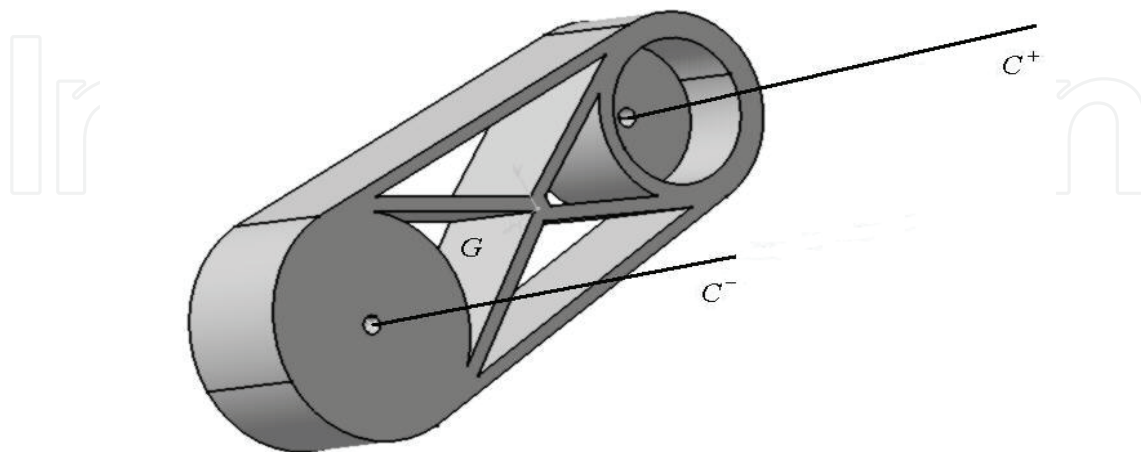


Fig. 3. The generic link of the considered generalised planar robot.

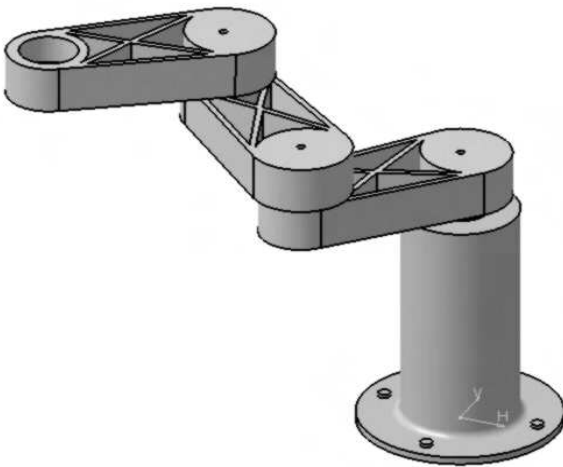


Fig. 4. The considered generalised planar robot with three links and horizontal work plane. Referring to Figs. 5 and 6 the next notations are also employed:

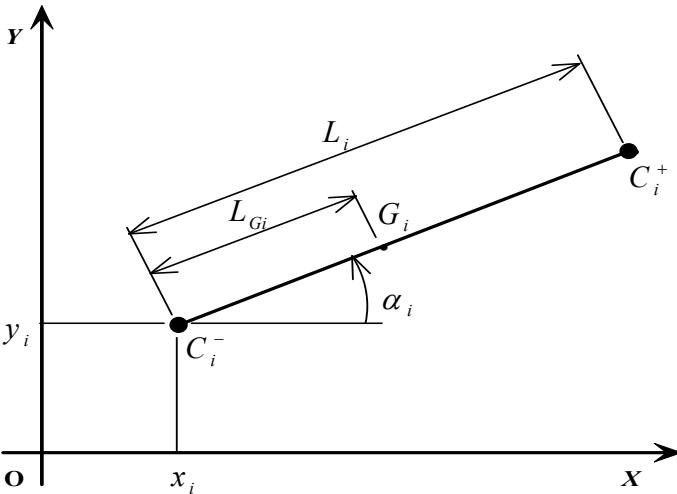


Fig. 5. Geometric characteristics of the i -th link.

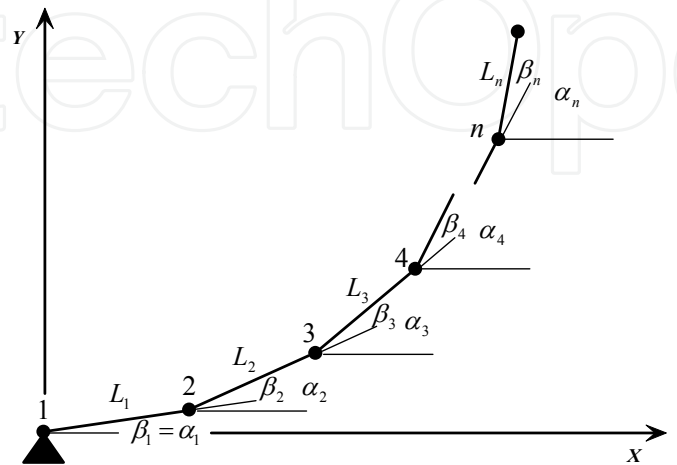


Fig. 6. Schematic representation of a planar robot with n links.

- $q_i = [x_i \ y_i \ \alpha_i]^T$ = absolute coordinates of the i -th link,
 I_i = inertia matrix of the i -th link in terms of the coordinates q_i ,
 M_i = total mass of the i -th link,
 J_i, ρ_i = inertia moment and radius of the i -th link with respect to the rotation axis C_i^- ,
 N_i = static moment of the i -th link with respect to the plane containing C_i^- and orthogonal to the plane passing through C_i^- and C_i^+ ,
 L_{Gi} = distance of the center of mass G_i of the i -th link from the axis C_i^- ,
 $\alpha_{1...i} = [\alpha_1 \ \alpha_2 \ \dots \ \alpha_i]^T$ = vector of the absolute angular coordinates of the chain constituted by the first i links,
 $\beta_{1...i} = [\beta_1 \ \beta_2 \ \dots \ \beta_i]^T$ = vector of the relative angular coordinates of the chain constituted by the first i links,
 A_i = inertia matrix of the i -th link in terms of the coordinates $\alpha_{1...i}$,
 B_i = inertia matrix of the i -th link in terms of the coordinates $\beta_{1...i}$,
 $\alpha = [\alpha_1 \ \alpha_2 \ \dots \ \alpha_n]^T$ = vector of the absolute angular coordinates of the robot,
 $\beta = [\beta_1 \ \beta_2 \ \dots \ \beta_n]^T$ = vector of the relative angular coordinates of the robot,
 A = inertia matrix of the robot in terms of the coordinates α ,
 B = inertia matrix of the robot in terms of the coordinates β .

3.2 Main results

Consider a generic link of a planar robot, with revolute joints. The following theorem can be stated.

Theorem 2. The kinetic energy T_i of the link of a planar robot, whose rotation axes are parallel and center of mass belongs to the plane that contains the respective rotation axes (see Figs. 3 and 5) can be calculated by the relationship

$$T_i = \frac{1}{2} \dot{q}_i^T I_i \dot{q}_i, \quad (10)$$

where

$$I_i = \begin{bmatrix} M_i & 0 & -N_i \sin \alpha_i \\ 0 & M_i & N_i \cos \alpha_i \\ -N_i \sin \alpha_i & N_i \cos \alpha_i & J_i \end{bmatrix} = M_i \begin{bmatrix} 1 & 0 & -L_{Gi} \sin \alpha_i \\ 0 & 1 & L_{Gi} \cos \alpha_i \\ -L_{Gi} \sin \alpha_i & L_{Gi} \cos \alpha_i & \rho_i^2 \end{bmatrix}. \quad (11)$$

Proof. By omitting subscript i , the coordinates of a generic point P of link L are (see Fig. 7)

$$\begin{aligned} x_p &= x + \xi \cos \alpha - \eta \sin \alpha \\ y_p &= y + \xi \sin \alpha + \eta \cos \alpha, \end{aligned} \quad (12)$$

from which it is easy to prove that the square of the velocity of P is

$$\begin{aligned} V_p^2 &= \dot{x}_p^2 + \dot{y}_p^2 = \dot{x}^2 + \dot{y}^2 + \dot{\alpha}^2 (\xi^2 + \eta^2) - 2\dot{x}\dot{\alpha}(\xi \sin \alpha + \eta \cos \alpha) + 2\dot{y}\dot{\alpha}(\xi \cos \alpha - \eta \sin \alpha) \\ &= V_c^2 + \dot{\alpha}^2 d^2 - 2\dot{x}\dot{\alpha}(\xi \sin \alpha + \eta \cos \alpha) + 2\dot{y}\dot{\alpha}(\xi \cos \alpha - \eta \sin \alpha). \end{aligned} \quad (13)$$

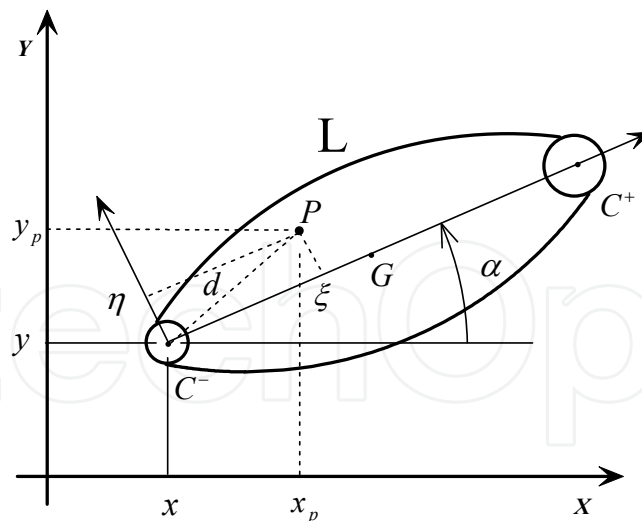


Fig. 7. Determination of the coordinates of a generic point P of the link L .

Since the kinetic energy of link L is $T = \frac{1}{2} \int_L V_p^2 dm$, and for the hypothesis made about the characteristics of the link, the proof of the theorem easily follows.

Consider a planar robot with n links (see Fig. 6). The following theorem holds:

Theorem 3. The kinetic energy of the i -th link of a generic planar robot can be calculated by the relationship

$$T_i = \frac{1}{2} \dot{\alpha}_{12\dots i}^T A_i \dot{\alpha}_{12\dots i}, \quad (14)$$

where the inertia matrix in terms of the angular coordinates $\alpha_{1\dots i}$, expressed by $A_i = \{a_{i,hk}\}$, can be calculated by using the formula:

$$A_i = R_i^T I_i R_i, \quad (15)$$

in which

$$R_i = \begin{bmatrix} -L_1 \sin \alpha_1 & -L_2 \sin \alpha_2 & \dots & -L_{i-1} \sin \alpha_{i-1} & 0 \\ L_1 \cos \alpha_1 & L_2 \cos \alpha_2 & \dots & L_{i-1} \cos \alpha_{i-1} & 0 \\ 0 & 0 & \dots & 0 & 1 \end{bmatrix} \quad (16)$$

or more directly by the formulae:

$$a_{i,hh} = \begin{cases} M_i L_h^2, & \text{if } h < i \\ J_i = M_i \rho_i^2, & \text{if } h = i, \end{cases} \quad (17a)$$

$$a_{i,hk} = \begin{cases} M_i L_h L_k \cos(\alpha_k - \alpha_h), & \text{if } h < k < i \\ M_i L_h L_{Gi} \cos(\alpha_i - \alpha_h), & \text{if } h < k = i. \end{cases} \quad (17b)$$

An alternative equivalent and more mnemonic expression for (17a)-(17b) is

1) $\forall h \leq i$,

$$a_{i,hh} = M_i \tilde{L}_h^2, \quad \tilde{L}_h = \begin{cases} L_h, & \text{if } h < i \\ \rho_i, & \text{if } h = i, \end{cases} \quad (18a)$$

2) $\forall h < k \leq i$,

$$a_{i,hk} = M_i \tilde{L}_h \tilde{L}_k \cos(\alpha_k - \alpha_h), \quad \tilde{L}_{h,k} = \begin{cases} L_{h,k}, & \text{if } h, k < i \\ L_{Gi}, & \text{if } h, k = i. \end{cases} \quad (18b)$$

Proof. For the generic i -th link, in the hypothesis that $x_1 = y_1 = 0$, it is

$$\begin{aligned} x_i &= L_1 \cos \alpha_1 + L_2 \cos \alpha_2 + \dots + L_i \cos \alpha_i \\ y_i &= L_1 \sin \alpha_1 + L_2 \sin \alpha_2 + \dots + L_i \sin \alpha_i \\ \alpha_i &= \alpha_i, \end{aligned} \quad (19)$$

from which

$$\begin{aligned} \dot{x}_i &= -\dot{\alpha}_1 L_1 \sin \alpha_1 - \dot{\alpha}_2 L_2 \sin \alpha_2 - \dots - \dot{\alpha}_i L_i \sin \alpha_i \\ \dot{y}_i &= \dot{\alpha}_1 L_1 \cos \alpha_1 + \dot{\alpha}_2 L_2 \cos \alpha_2 + \dots + \dot{\alpha}_i L_i \cos \alpha_i \\ \dot{\alpha}_i &= \dot{\alpha}_i, \end{aligned} \quad (20)$$

or

$$\dot{q}_i = R_i \dot{\alpha}_{1\dots i}. \quad (21)$$

It follows that

$$T_i = \frac{1}{2} \dot{q}_i^T I_i \dot{q}_i = \frac{1}{2} \dot{\alpha}_{12\dots i}^T R_i^T I_i R_i \dot{\alpha}_{1\dots i} \quad (22)$$

and hence (15).

From (15) it is:

1) $\forall h \leq k < i$

$$a_{i,hk} = \begin{bmatrix} -L_h \sin \alpha_h & L_h \cos \alpha_h & 0 \end{bmatrix} \cdot \begin{bmatrix} M_i & 0 & -N_i \sin \alpha_i \\ 0 & M_i & N_i \cos \alpha_i \\ -N_i \sin \alpha_i & N_i \cos \alpha_i & J_i \end{bmatrix} \cdot \begin{bmatrix} -L_k \sin \alpha_k \\ L_k \cos \alpha_k \\ 0 \end{bmatrix} = M_i L_h L_k \cos(\alpha_k - \alpha_h); \quad (23)$$

2) $\forall h < k = i$

$$a_{i,hi} = \begin{bmatrix} -L_h \sin \alpha_h & L_h \cos \alpha_h & 0 \end{bmatrix} \cdot \begin{bmatrix} M_i & 0 & -N_i \sin \alpha_i \\ 0 & M_i & N_i \cos \alpha_i \\ -N_i \sin \alpha_i & N_i \cos \alpha_i & J_i \end{bmatrix} \cdot \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} = L_h N_i \cos(\alpha_i - \alpha_h); \quad (24)$$

3) $h = k = i$

$$a_{i,iii} = \begin{bmatrix} 0 & 0 & 1 \end{bmatrix} \cdot \begin{bmatrix} M_i & 0 & -N_i \sin \alpha_i \\ 0 & M_i & N_i \cos \alpha_i \\ -N_i \sin \alpha_i & N_i \cos \alpha_i & J_i \end{bmatrix} \cdot \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} = J_i \quad (25)$$

and hence (17a), (17b).

In view of Theorem 3, the kinetic energy of the robot is

$$T = T_1 + T_2 + \dots + T_n = \frac{1}{2} \dot{\alpha}^T A \dot{\alpha}, \quad (26)$$

where

$$\alpha = \alpha_{1...n} = [\alpha_1 \ \alpha_2 \ \dots \ \alpha_n]^T, \quad (27)$$

and the inertia matrix $A = A(\alpha)$ of the robot in terms of the angular coordinates α is obtained by assembling the matrices A_i as illustrated in Fig. 8 (namely, by summing to all the principal minors of A_n of order $i < n$ the respective A_i 's).

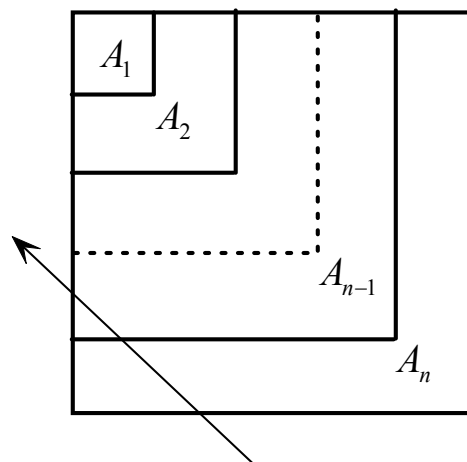


Fig. 8. Graphical representation of the algorithm that computes the robot inertia matrix.

In MATLAB-like instructions the algorithm to calculate the inertia matrix is:

```
A=An;
for i=n-1:-1:1
A(1:i,1:i)=A(1:i,1:i)+Ai;
end.
```

Moreover, since $\alpha = U\beta$, where

$$U = \begin{bmatrix} 1 & 0 & \cdot & 0 \\ 1 & 1 & \cdot & 0 \\ \cdot & \cdot & \cdot & \cdot \\ 1 & 1 & 1 & 1 \end{bmatrix}, \quad (28)$$

the kinetic energy of the robot can be calculated by using the relationship alternatively

$$T = \frac{1}{2} \dot{\beta}^T B \dot{\beta}, \quad (29)$$

where the inertia matrix $B = B(\beta)$ of the robot in terms of the angular coordinates β is obtained by

$$B = U^T A(\alpha) U = U^T A(U^{-1}(\beta)) U. \quad (30)$$

Note that the computation of the inertia matrix is performed by jointly using the absolute and relative coordinates. This allows simplifying, in a natural manner, the algebraic expression of the analytical model.

Clearly, if for a given β and hence for a given $\alpha = U\beta$ the interest is only about the numerical value of B , it is possible to numerically compute the matrix A preliminarily (from α by using Theorem 3 and the algorithm illustrated in Fig. 7) and then the matrix B through the relationship $B = U^T A U$ or through the following algorithm described in MATLAB-like instructions:

```

for i=n-1:-1:1
    A(:,i)=A(:,i)+A(:,i+1);
end
for i=n-1:-1:1
    A(1:i,i+1)=A(i+1,1:i)';
    A(i,1:i)=A(i,1:i)+A(i+1,1:i);
end
B=A.
```

Remark 3. It is important to note that, although by virtue of Theorem 3 the analytical expression of A_i is relatively simple, the analytical expression of A is quite complex and, as a result, the expression of B is extremely complex (e.g. the entry $B(1,1)$ of B is the sum of all the elements of all the matrices A_i !).

Once the inertia matrix B is obtained, the dynamic model of the robot can be derived by the Euler-Lagrange equation in the usual form (1) or in the alternative form (6). In the case of a planar robot, with link characteristics as required by Theorem 2, the efficiency of the method illustrated in Theorem 1 can be further improved by using the following result.

Theorem 4. The gradient of the kinetic energy of the i -th link in terms of $\alpha_{12\dots i}$ $[\beta_{12\dots i}]$ can be calculated by the relationship

$$\frac{\partial T_i}{\partial \beta_{12\dots i}} = c_i = - \begin{bmatrix} 0 \\ M_i \sum_{h < k: \text{var } \beta_2} \tilde{L}_h \dot{\alpha}_h \tilde{L}_k \dot{\alpha}_k \sin(\alpha_k - \alpha_h) \\ M_i \sum_{h < k: \text{var } \beta_3} \tilde{L}_h \dot{\alpha}_h \tilde{L}_k \dot{\alpha}_k \sin(\alpha_k - \alpha_h) \\ \vdots \\ M_i \sum_{h < k: \text{var } \beta_i} \tilde{L}_h \dot{\alpha}_h \tilde{L}_k \dot{\alpha}_k \sin(\alpha_k - \alpha_h) \end{bmatrix},$$

$$\frac{\partial T_i}{\partial \beta_{12\dots i}} = c_i = - \begin{bmatrix} 0 \\ M_i \sum_{h < k: \text{var } \beta_2} \tilde{L}_h \sum_{p=1}^h \dot{\beta}_p \tilde{L}_k \sum_{q=1}^k \dot{\beta}_q \sin(\sum_{r=h+1}^k \beta_r) \\ M_i \sum_{h < k: \text{var } \beta_3} \tilde{L}_h \sum_{p=1}^h \dot{\beta}_p \tilde{L}_k \sum_{q=1}^k \dot{\beta}_q \sin(\sum_{r=h+1}^k \beta_r) \\ \vdots \\ M_i \sum_{h < k: \text{var } \beta_i} \tilde{L}_h \sum_{p=1}^h \dot{\beta}_p \tilde{L}_k \sum_{q=1}^k \dot{\beta}_q \sin(\sum_{r=h+1}^k \beta_r) \end{bmatrix}, \quad (31a)$$

where

$$\tilde{L}_m = \begin{cases} L_m, & \text{if } m < i \\ L_{Gi}, & \text{if } m = i \end{cases}, \quad m = \{h, k\}, \quad (31b)$$

and the notation $h < k : \text{var } \beta_l$, $l = 2, \dots, i$, indicates the set of the pairs h, k , with $h < k$, such that the angle $\alpha_{hk} = \alpha_k - \alpha_h = \beta_{h+1} + \beta_{h+2} + \dots + \beta_k$ between the links h and k varies with β_l .

Proof. In view of Theorem 3, the kinetic energy of the generic i -th link is

$$\begin{aligned} T_i &= \frac{1}{2} M_i \sum_{h=1, \dots, i-1} L_h^2 \dot{\alpha}_h^2 + \frac{1}{2} M_i \dot{\alpha}_i^2 + M_i \sum_{\substack{h=1, \dots, i-1 \\ k=h+1, \dots, i}} \tilde{L}_h \tilde{L}_k \dot{\alpha}_h \dot{\alpha}_k \cos(\alpha_k - \alpha_h) \\ &= T_c + M_i \sum_{\substack{h=1, \dots, i-1 \\ k=h+1, \dots, i}} \tilde{L}_h \tilde{L}_k \dot{\alpha}_h \dot{\alpha}_k \cos(\alpha_k - \alpha_h), \end{aligned} \quad (32)$$

where T_c is the portion of kinetic energy that is independent of $\alpha_{12\dots i}$.

From (32) the expression (31a) of the gradient of T_i in terms of $\alpha_{12\dots i}$ [$\beta_{12\dots i}$] easily follow.

The vector c can be simply obtained by suitably assembling the various vectors c_i (i.e. by progressively summing to the first i components of c_n , $i < n$, the respective c_i 's).

In MATLAB-like instructions:

```
c=cn;
for i=n-1:-1:2
c(1:i)=c(1:i)+ci;
end.
```

Finally, it is important to note that if the work plane of the robot is vertical, the gravitational torques can be calculated by using the following result.

Theorem 5. If the work plane of the robot is vertical, the weight of the i -th link originates a torque on the k -th joint, $k \leq i$:

$$\gamma_k = M_i g \sum_{m=k}^i \tilde{L}_m \cos \alpha_m, \quad (33)$$

where g is the gravitational acceleration and the lengths \tilde{L}_m are given by (31b).

Proof. The proof easily follows by considering Fig. 6.

Remark 4. The proposed method efficiency has been explicitly compared to the Articulated Body approach efficiency in terms of number of flops (in MatlabTM environment), for executing the same elementary integration step by using the absolute Lagrangian coordinates (see Fig. 9). From this figure it results that, for industrial robots with a limited number of links, the proposed method is more efficient.

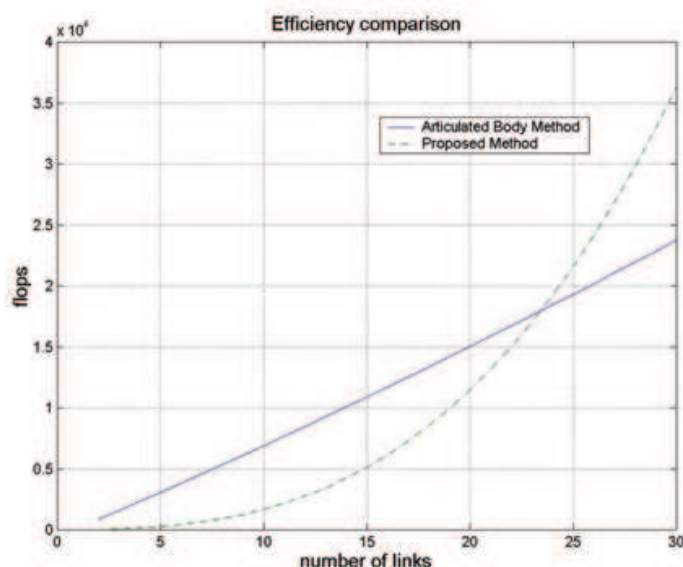


Fig. 9. Efficiency comparison between the Articulated Body method and the proposed method.

Remark 5. The above-derived analytical expressions, in explicit form, of B , c and γ are useful for speeding up the simulation, for comparing them to the proposed numerical approach and for finally evaluating which terms can be simplified or neglected in the case of robots with many and/or flexible links. Thus, for instance, for flexible link robots models, obtained via discretization of the links, the terms relative to the deformation angles can be simplified by substituting the *sine* function with the respective argument or even be completely neglected, considering the terms relative to the motion only.

4. Spatial robot modeling

For spatial robots with generic shape links and connected, for the sake of brevity, with spherical joints, several theorems are formulated and demonstrated in a simple manner and some algorithms that allow computing in an efficient way analytically or numerically both the inertia matrix and the gradient of the potential energy are provided.

4.1 Hypotheses and notations

In the following, for brevity, only spatial robots constituted by n links of generic shape, connected by spherical joints are considered.

Moreover, the following notations will be used (see Fig. 10):

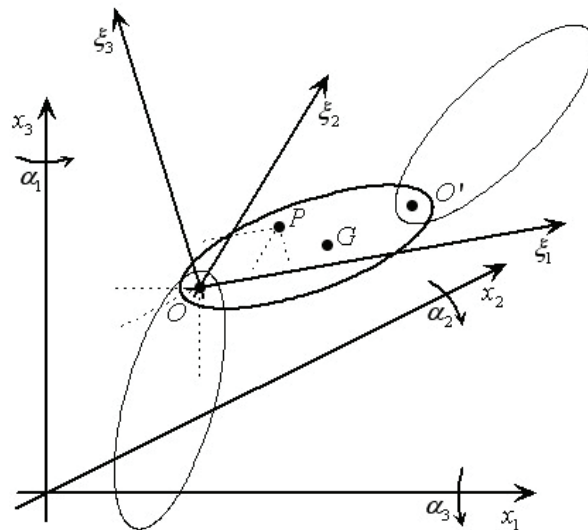


Fig. 10. Coordinates of the generic point P of a link with respect to the body frame and to the inertial frame.

$\mathbf{q}_i = [\mathbf{x}_{O_i}^T \quad \boldsymbol{\alpha}_i^T]^T = [x_{O1i} \quad x_{O2i} \quad x_{O3i} \quad \alpha_{1i} \quad \alpha_{2i} \quad \alpha_{3i}]^T$ = "absolute coordinates" of the i -th link,

$I_i = \{I_{ihk}\}$ = inertia matrix of the i -th link with respect to the axes ξ_1, ξ_2, ξ_3 ,

$\boldsymbol{\xi}_{G_i} = [\xi_{G1i} \quad \xi_{G2i} \quad \xi_{G3i}]^T$ = coordinates of the center of mass G_i in the body frame,

$\boldsymbol{\xi}_{O'_i} = [\xi_{O'1i} \quad \xi_{O'2i} \quad \xi_{O'3i}]^T$ = coordinates of the rotation center O'_i in the body frame,

M_i = total mass of the i -th link,

$\boldsymbol{\alpha}_{1\dots i} = [\boldsymbol{\alpha}_1^T \quad \boldsymbol{\alpha}_2^T \quad \dots \quad \boldsymbol{\alpha}_i^T]^T$ = "angular absolute coordinates" of the chain constituted by the first i links,

$\mathfrak{R}_i(\boldsymbol{\alpha}_i)$ = rotation matrix of the i -th link,

$$R_{ih} = \frac{\partial \mathfrak{R}_i(\boldsymbol{\alpha}_i)}{\partial \alpha_{ih}},$$

A_i = inertia matrix of the i -th link with respect to the coordinates $\boldsymbol{\alpha}_{1\dots i}$,

$\boldsymbol{\alpha} = \boldsymbol{\alpha}_{1\dots n}$ = "absolute angular coordinates" of the robot,

A = inertia matrix of the robot with respect to the coordinates $\boldsymbol{\alpha}$,

I = identity matrix of order 3,

$s_{hi} = \sin(\alpha_{hi})$,

$c_{hi} = \cos(\alpha_{hi})$.

4.2 Main results

Let us consider a generic link of a generic shape of a spatial robot with spherical joints.

The following theorems and algorithms for the computation of the kinetic energy T_i of the i -th link first as a function of the absolute coordinates \mathbf{q}_i of the i -th link and then as a function of the absolute angular coordinates $\boldsymbol{\alpha}_{1\dots i}$ of the chain constituted by the first i links are provided.

Theorem 6. The kinetic energy of the i -th link may be calculated by means of the relationship

$$T_i = \frac{1}{2} \dot{\mathbf{q}}_i^T \mathfrak{S}_i \dot{\mathbf{q}}_i, \quad (34)$$

where

$$\mathfrak{S}_i = \begin{bmatrix} M_i I & M_i \sum_h R_{hi} \xi_{Ghi} \\ M_i \sum_h R_{hi}^T \xi_{Ghi} & \sum_{h,k} R_{hi}^T R_{ki} I_{hki} \end{bmatrix}. \quad (35)$$

Proof. By omitting subscript i , the coordinates of the generic point P of the link in the fixed frame (see Fig. 10) are given by

$$\mathbf{x}_P = \mathbf{x}_O + \mathfrak{R}(\boldsymbol{\alpha}) \boldsymbol{\xi}_P, \quad (36)$$

where \mathfrak{R} is the rotation matrix of the body frame with respect to the fixed frame and $\boldsymbol{\xi}_P$ is the vector of the coordinates of the generic point P of the link in the body frame.

From (36) it follows

$$\dot{\mathbf{x}}_P = \dot{\mathbf{x}}_O + \sum_h R_h \xi_{Ph} \dot{\boldsymbol{\alpha}}, \quad (37)$$

where

$$R_h = \frac{\partial \mathfrak{R}}{\partial \alpha_h}, h = 1, 2, 3. \quad (38)$$

From (37) it is

$$V^2 = \dot{\mathbf{x}}_P^T \dot{\mathbf{x}}_P = \dot{\mathbf{x}}_O^T \dot{\mathbf{x}}_O + \dot{\boldsymbol{\alpha}}^T \sum_{h,k} R_h^T R_k \xi_{Ph} \xi_{Pk} \dot{\boldsymbol{\alpha}} + 2 \dot{\mathbf{x}}_O^T \sum_h R_h \xi_{Ph} \dot{\boldsymbol{\alpha}}, \quad (39)$$

from which

$$\begin{aligned} T &= \frac{1}{2} \int V^2 dm = \\ &= \frac{1}{2} M \dot{\mathbf{x}}_O^T \dot{\mathbf{x}}_O + \frac{1}{2} \dot{\boldsymbol{\alpha}}^T \sum_{h,k} R_h^T R_k I_{hk} \dot{\boldsymbol{\alpha}} + M \dot{\mathbf{x}}_O^T \sum_h R_h \xi_{Gh} \dot{\boldsymbol{\alpha}} = \\ &= \frac{1}{2} \begin{pmatrix} \dot{\mathbf{x}}_O^T & \dot{\boldsymbol{\alpha}}^T \end{pmatrix} \begin{bmatrix} MI & M \sum_h R_h \xi_{Gh} \\ M \sum_h R_h^T \xi_{Gh} & \sum_{h,k} R_h^T R_k I_{hk} \end{bmatrix} \begin{pmatrix} \dot{\mathbf{x}}_O \\ \dot{\boldsymbol{\alpha}} \end{pmatrix} \end{aligned} \quad (40)$$

and hence the proof.

Remark 6. The vector of the coordinates of the center of mass $\boldsymbol{\xi}_{G_i}$ and the inertia matrix $\{I_{hk}\}$ of the i -th link, in the case of a complex structure, may be easily evaluated by making use of software packages, such as CATIA™.

Theorem 7. If the Euler angles of the i -th link $\boldsymbol{\alpha}_i = [\alpha_{1i} \ \alpha_{2i} \ \alpha_{3i}]^T$ are the ZYX angles, also called RPY angles (Roll-Pitch-Yaw), (see Fig. 10), by omitting subscript i , the matrices R_h result:

$$R_1 = \begin{bmatrix} -s_1c_2 & -s_1s_2s_3 - c_1c_3 & -s_1s_2c_3 + c_1s_3 \\ c_1c_2 & c_1s_2s_3 - s_1c_3 & c_1s_2c_3 + s_1s_3 \\ 0 & 0 & 0 \end{bmatrix} \quad (41)$$

$$R_2 = \begin{bmatrix} -c_1s_2 & c_1c_2s_3 & c_1c_2c_3 \\ -s_1s_2 & s_1c_2s_3 & s_1c_2c_3 \\ -c_2 & -s_2s_3 & -s_2c_3 \end{bmatrix} \quad (42)$$

$$R_3 = \begin{bmatrix} 0 & c_1s_2c_3 + s_1s_3 & -c_1s_2s_3 + s_1c_3 \\ 0 & s_1s_2c_3 - c_1s_3 & -s_1s_2s_3 - c_1c_3 \\ 0 & c_2c_3 & -c_2s_3 \end{bmatrix}. \quad (43)$$

Proof. The rotation matrix $\mathfrak{R}_i(\alpha_i)$ of the i -th link, by omitting subscript i , results (43)

$$\mathfrak{R}(\alpha) = \begin{bmatrix} c_1c_2 & c_1s_2s_3 - s_1c_3 & c_1s_2c_3 + s_1s_3 \\ s_1c_2 & s_1s_2s_3 + c_1c_3 & s_1s_2c_3 - c_1s_3 \\ -s_2 & c_2s_3 & c_2c_3 \end{bmatrix}, \quad (44)$$

from which, for (38), the proof.

The computation of the matrices \mathfrak{I}_i may be sped up by the following algorithm, where subscript i has been omitted.

Algorithm 1.

Step 1. Compute

$$s_1, s_2, s_3; \quad c_1, c_2, c_3. \quad (45)$$

Step 2. Compute

$$\begin{aligned} & s_1s_2, s_1s_3, s_2s_3; \quad c_1c_2, c_1c_3, c_2c_3; \\ & s_1c_2, s_1c_3, s_2c_3; \quad c_1s_2, c_1s_3, c_2s_3; \\ & (s_1s_2)s_3, (s_1s_2)c_3; \quad (s_1c_2)s_3, (s_1c_2)c_3; \\ & (c_1c_2)s_3, (c_1c_2)c_3; \quad (c_1s_2)s_3, (c_1s_2)c_3. \end{aligned} \quad (46)$$

Step 3. Compute

$$R_1^T R_1, R_1^T R_2, R_1^T R_3; \quad R_2^T R_2, R_2^T R_3; \quad R_3^T R_3. \quad (47)$$

If the axis ξ_{1i} of the body frame is chosen such as to contain the center of gravity G_i and it results principal of inertia then, by choosing the other two axes parallel to the other inertia principal axes, the computation of the matrix \mathfrak{I}_i may be further sped up, as it results from the following theorem.

Theorem 8. If the body reference axes of the i -th link are parallel to the inertia principal axes and the center of gravity G_i belong to the axis ξ_{1i} , the matrix \mathfrak{I}_i results

$$\mathfrak{I}_i = \begin{bmatrix} M_i I & M_i R_{1i} \xi_{G1i} \\ M_i R_{1i}^T \xi_{G1i} & \sum_h R_{hi}^T R_{hi} I_{hhi} \end{bmatrix}. \quad (48)$$

Proof. The proof easily follows from (35) and from

$$\xi_{G2i}=0, i=2,3; \quad I_{hki}=0, h=1,2, k=h+1,3. \quad (49)$$

Let us consider now a robot with n links (see Fig. 11). The following theorems and algorithms for the computation of the inertia matrix A of the whole robot, as a function of the absolute angular coordinates α of the robot, hold.

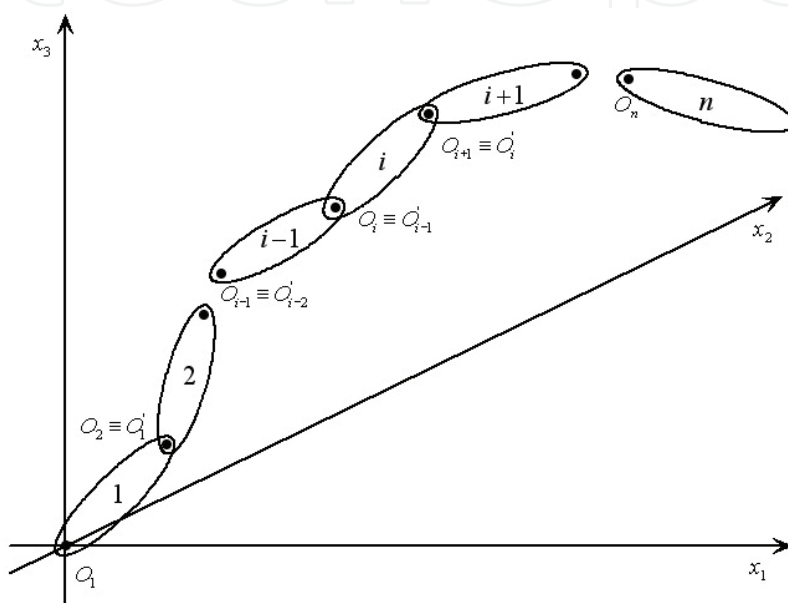


Fig. 11. Schematic representation of a spatial robot with n links.

Theorem 9. The kinetic energy of the i -th link of a generic spatial robot may be computed by means of the relationship

$$T_i = \frac{1}{2} \dot{\mathbf{a}}_{1\dots i}^T A_i \dot{\mathbf{a}}_{1\dots i}, \quad (50)$$

where

$$A_i = \begin{cases} \mathfrak{I}_{22i}, & \text{if } i = 1 \\ \begin{bmatrix} M_i V_{1...i-1}^T V_{1...i-1} & V_{1...i-1}^T \mathfrak{I}_{22i} \\ \mathfrak{I}_{22i}^T V_{1...i-1} & \mathfrak{I}_{22i} \end{bmatrix}, & \text{if } i = 2, 3, \dots, n \end{cases} \quad (51)$$

in which:

$$\begin{aligned} \mathfrak{T}_{12i} &= M_i \sum_h R_{hi} \xi_{Ghi}, \quad \mathfrak{T}_{22i} = \sum_{h,k} R_{hi}^T R_{ki} I_{hki} \\ V_{1...i-1} &= [V_1 \quad V_2 \quad \cdots \quad V_{i-1}], \quad V_j = \sum_h R_{hj} \xi_{O'hj}, \quad j = 1, 2, \dots, i-1. \end{aligned} \quad (52)$$

Proof. From the following relationships (see Fig. 11 and (37))

$$\dot{\mathbf{x}}_{O_i} = \dot{\mathbf{x}}_{O_{i-1}} = \dot{\mathbf{x}}_{O_{i-1}} + \sum_h R_{hi-1} \xi_{O'hi-1} \dot{\alpha}_{i-1}, \quad i = 2, \dots, n, \quad \dot{\mathbf{x}}_{O_1} = 0, \quad (53)$$

it is

$$\dot{\mathbf{x}}_{O_i} = \begin{bmatrix} V_1 & V_2 & \dots & V_{i-1} \end{bmatrix} \begin{bmatrix} \dot{\alpha}_1 \\ \dot{\alpha}_2 \\ \vdots \\ \dot{\alpha}_{i-1} \end{bmatrix} = V_{1\dots i-1} \dot{\alpha}_{1\dots i-1}, \quad (54)$$

from which

$$\dot{\mathbf{q}}_i = \begin{bmatrix} \dot{\mathbf{x}}_{O_i} \\ \dot{\alpha}_i \end{bmatrix} = \begin{bmatrix} V_{1\dots i-1} \dot{\alpha}_{1\dots i-1} \\ \dot{\alpha}_i \end{bmatrix} = \begin{bmatrix} V_{1\dots i-1} & 0 \\ 0 & I \end{bmatrix} \dot{\alpha}_{1\dots i}. \quad (55)$$

By substituting (55) in (34), it is

$$\begin{aligned} T_i &= \frac{1}{2} \dot{\mathbf{q}}_i^T \mathfrak{I}_i \dot{\mathbf{q}}_i = \\ &= \frac{1}{2} \dot{\alpha}_{1\dots i}^T \begin{bmatrix} V_{1\dots i-1}^T & 0 \\ 0 & I \end{bmatrix} \begin{bmatrix} M_i I & \mathfrak{I}_{12i} \\ \mathfrak{I}_{12i}^T & \mathfrak{I}_{22i} \end{bmatrix} \begin{bmatrix} V_{1\dots i-1} & 0 \\ 0 & I \end{bmatrix} \dot{\alpha}_{1\dots i} = \\ &= \frac{1}{2} \dot{\alpha}_{1\dots i}^T \begin{bmatrix} M_i V_{1\dots i-1}^T V_{1\dots i-1} & V_{1\dots i-1}^T \mathfrak{I}_{12i} \\ \mathfrak{I}_{12i}^T V_{1\dots i-1} & \mathfrak{I}_{22i} \end{bmatrix} \dot{\alpha}_{1\dots i} \end{aligned} \quad (56)$$

from which the proof follows.

The next theorem, similar to Theorem 8, is useful to simplify the computation of the matrices A_i .

Theorem 10. If the axes of the body frame of link j , $j = 1, 2, \dots, i-1$, are parallel to the inertia principal axes and the centers of gravity G_j and the rotation centers O_j belong to the axis ξ_{1j} , matrices (52) may be more easily calculated by the relations:

$$\mathfrak{I}_{12i} = M_i R_{1i} \xi_{G1i}, \quad \mathfrak{I}_{22i} = \sum_h R_{hi}^T R_{hi} I_{hhi}, \quad (57)$$

$$V_{1\dots i-1} = [V_1 \quad V_2 \quad \dots \quad V_{i-1}], \quad V_j = R_{1j} \xi_{O'1j}, \quad j = 1, 2, \dots, i-1. \quad (58)$$

Proof. The proof easily follows from Theorem 8 and from the fact that

$$\xi_{O'hj} = 0, \quad h = 2, 3, \quad j = 1, 2, \dots, i-1. \quad (59)$$

The following theorems may be effectively used to compute the gradient of the gravitational potential energy.

Theorem 11. The gradient of the gravitational potential energy of the i -th link U_i with respect to the angular coordinates $\alpha_{1\dots i}$ may be calculated by the formula

$$\frac{\partial U_i}{\partial \alpha_{1...i}} = \mathbf{Y}_i = -M_i \begin{bmatrix} R_{11}\xi_{0_1} & R_{21}\xi_{0_1} & R_{31}\xi_{0_1} & R_{12}\xi_{0_2} & R_{22}\xi_{0_2} & R_{32}\xi_{0_2} & \cdots \\ \cdots & R_{1i}\xi_{G_i} & R_{2i}\xi_{G_i} & R_{3i}\xi_{G_i} \end{bmatrix}^T \mathbf{g}, \quad (60)$$

where \mathbf{g} is the gravity acceleration vector in the inertial frame.

Proof. It is easy to verify that the gravitational potential energy of the generic i -th link results

$$U_i = -M_i \mathbf{g}^T (\mathfrak{R}_1(\alpha_1)\xi_{O_1} + \mathfrak{R}_2(\alpha_2)\xi_{O_2} + \cdots + \mathfrak{R}_i(\alpha_i)\xi_{G_i}). \quad (61)$$

Relation (60) easily follows from (61).

The next theorem is useful to simplify the computation of $\frac{\partial U_i}{\partial \alpha_{1...i}}$.

Theorem 12. If the x_3 -axis of the inertial frame is vertical and upwards oriented, the rotation centers O_j of the j -th links, $j=1,2,\dots,i-1$, belong to the axes ξ_{1j} of the body frames, the center of gravity G_i of the i -th link belongs to the axis ξ_{1i} and the Euler angles of the j -th links, $j=1,2,\dots,i$, are the ZYX angles, the gradient $\frac{\partial U_i}{\partial \alpha_{1...i}}$ may be calculated much more easily by the relationship

$$\begin{aligned} \frac{\partial U_i}{\partial \alpha_{1...i}} &= \mathbf{Y}_i = \\ &= M_i \mathbf{g} \begin{bmatrix} 0 & c_{21}\xi_{O_1 1} & 0 & 0 & c_{22}\xi_{O_2 1} & 0 & \cdots & 0 & c_{2i}\xi_{G_i 1} & 0 \end{bmatrix}^T. \end{aligned} \quad (62)$$

Proof. Since the x_3 -axis is vertical and upwards oriented it is $\mathbf{g} = [0 \ 0 \ -g]^T$, where g is the gravity acceleration. Moreover, since $O_j \in \xi_{1j}$, $j=1,2,\dots,i-1$, $G_i \in \xi_{1i}$:

$$\xi_{O_j} = [\xi_{O_j 1} \ 0 \ 0]^T, \quad \xi_{G_i} = [\xi_{G_i 1} \ 0 \ 0]^T, \quad i=1\dots n. \quad (63)$$

The proof easily follows from (41)-(43), (60), (63).

4.3 Application of the main results

On the base of Theorem 9 the kinetic energy of the whole robot is

$$T = T_1 + T_2 + \dots + T_n = \frac{1}{2} \dot{\alpha}^T A \dot{\alpha}, \quad (64)$$

where the *inertia matrix* $A = A(\alpha)$ is obtained by assembling the matrices A_i according to the following efficient algorithm written in Matlab-like language:

```
A=An;  
for i=n-1:-1:1  
A(1:3*i,1:3*i)=A(1:3*i,1:3*i)+Ai;  
end.
```

Once the vectors $\gamma_i, i = 1, 2, \dots, n$, have been computed by the (60) or the (62), the gradient of the gravity energy γ is obtained by assembling the γ_i according to the following efficient algorithm written in Matlab-like language:

```
 $\gamma = \gamma_n$ ;  
for i=n-1:-1:1  
 $\gamma(1:i) = \gamma(1:i) + \gamma_i$ ;  
end.
```

Finally, by using the efficient integration scheme shown in Fig. 1 (Celentano & Iervolino, 2006), it is possible to get the dynamic model of the robot in a simple way.

4.4 Comparison in terms of efficiency with the articulated-body method

The efficiency of the method of modeling and simulation proposed for spatial robots has been explicitly compared to the efficiency of a Newtonian method of order N considered one of the most efficient methods in the literature, i.e. the Articulated-Body Method (Featherstone, 1987), (Featherstone & Orin, 2000). More in details, such comparison has been executed in terms of number of flops (in MATLAB™ environment), for the execution of the same integration step (see Table 1).

Number of links	Flops of the proposed method	Flops of the Articulated-Body method
1	362	1269
2	2459	7168
3	8400	13758
4	21451	20804
5	45868	28252

Table 1. Efficiency comparison between the proposed method and the Articulated-Body method.

From Table 1 it results that for practical spatial robots (which of course have a limited number of links), the proposed method is more efficient when $n < 4$, although spatial robots with three degrees of freedom spherical joints have been considered. Moreover, it is important to note that the great easiness of the proposed method with respect to the Articulated-Body one is an important index of its efficiency. However, the authors have the plan to continue and to improve this comparison.

5. Elements of flexible robots modeling

5.1 Methodology

The results stated in the previous sections allow obtaining, quite simply, accurate and efficient, from a computational point of view, finite-dimensional models. These models, for rectilinear links, can be obtained approximating the i -th link as follows:

1. in order to compute the kinetic energy and, therefore, the inertia matrix and the gradient of the kinetic energy, with $\nu_i + 1$ rigid sublinks with lengths $L_{1i}, \dots, L_{\nu_i+1,i}$, interconnected by relative angles $\beta_{1i}, \dots, \beta_{\nu_i}$ (see Fig. 12);
2. in order to compute the elastic energy, with ν_i flexible sublinks joined and deformable such that their axes make a suitable spline tangent to the axes of the intermediate rigid sublinks in the extreme points of the link and to the intermediate rigid sublinks in suitable interior points (see Fig. 12).

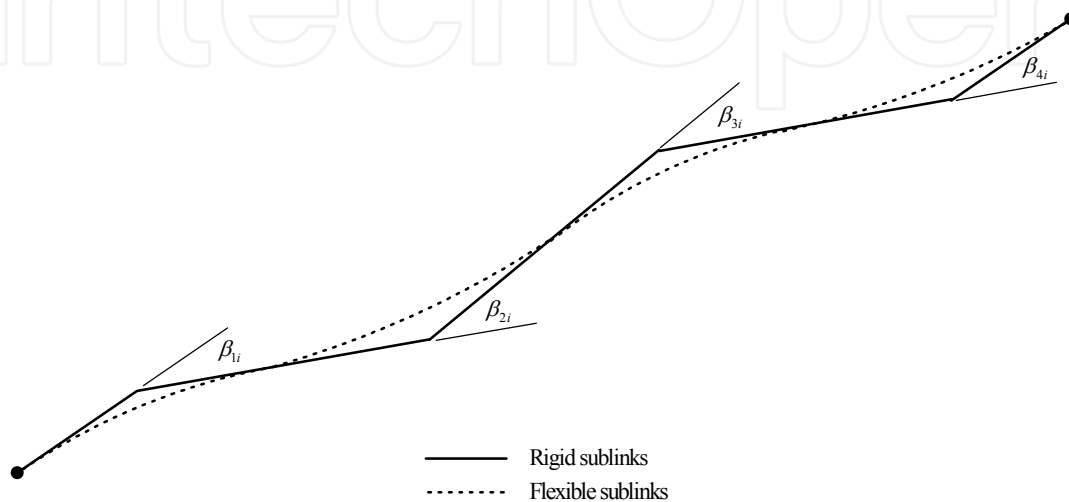


Fig. 12. Approximation of a flexible link with $\nu_i + 1 = 5$ rigid sublinks and with $\nu_i = 4$ flexible sublinks congruent to each other at the end.

In this way, neglecting, about the computation of the kinetic energy, the deformation angles respect to the angles of motion, it is easy to verify, from the results of Sections 2 and 3, that the inertia matrix can be computed as follows:

$$B = B(\beta_2, \dots, \beta_n) = B_1 + B_2 \cos \beta_2 + \dots + B_n \cos \beta_n + B_{23} \cos(\beta_2 + \beta_3) + B_{24} \cos(\beta_2 + \beta_3 + \beta_4) + \dots + B_{n-1,n} \cos(\beta_{n-1} + \beta_n), \quad (65)$$

where β_i is the angle of motion between the $(i-1)$ -th link and the i -th link, $i = 2, \dots, n$ (see Fig. 13).

From (65) it is easy to compute also the gradient of the kinetic energy respect to $\beta = [\beta_1 \ \beta_{11} \dots \beta_{\nu_1 1} \ \beta_2 \ \beta_{12} \dots \beta_{\nu_2 2} \ \beta_n \ \beta_{1n} \dots \beta_{\nu_n n}]^T$ using the following relation

$$c = \frac{1}{2} \begin{bmatrix} 0_{1, \nu_1+1} & \dot{\beta}^T \frac{\partial B}{\partial \beta_2} \dot{\beta} & \dots & 0_{1, \nu_{n-1}} & \dot{\beta}^T \frac{\partial B}{\partial \beta_n} \dot{\beta} & 0_{1, \nu_n} \end{bmatrix}^T. \quad (66)$$

There are different ways of approximating a flexible link using a chain of rigid sublinks and flexible sublinks congruent to each other at the end (Celentano, 2007).

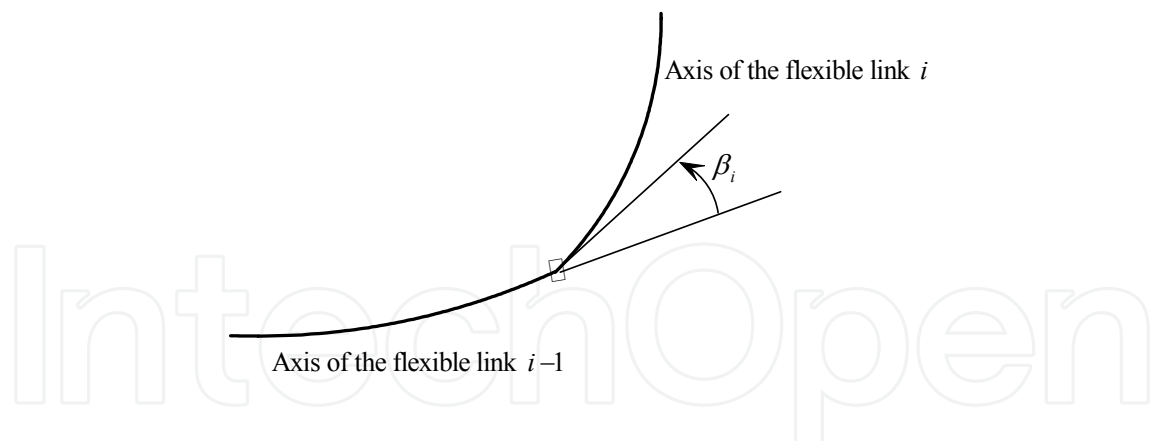


Fig. 13. Angle of motion between two flexible consecutive links.

In the hypothesis that the stiffness of the i -th link is uniform, a very efficient way for the choice of the sublinks is to consider the intermediate rigid sublinks with a length of L_i/v_i , the extreme ones with an half length, the flexible sublinks with a length of L_i/v_i and deformable as parabolic arcs. Consequently the elastic energy of the i -th link results

$$U_i = \frac{1}{2} K_i (\beta_{1i}^2 + \dots + \beta_{v_i}^2), \quad (67)$$

where

$$K_i = \frac{E_i I_i}{L_i/v_i} \quad (68)$$

is the stiffness of the flexible sublinks.

Therefore the vector of the elastic torques is

$$\gamma = [0 \ K_1 \beta_{11} \ K_1 \beta_{21} \dots K_1 \beta_{v_1} \ 0 \ K_2 \beta_{12} \ K_2 \beta_{22} \dots K_2 \beta_{v_2} \dots 0 \ K_n \beta_{1n} \ K_n \beta_{2n} \dots K_n \beta_{v_n}]^T. \quad (69)$$

Remark 7. The proposed method is simple to apply, is very efficient from a computational point of view and it provides a model with null error when the control and disturbance torques are constant. For control actions and disturbances, also in forces, with a fixed and limited band-pass, this model provides an error strongly decreasing when the number of sublinks increase. This fact clearly emerges from the following example.

5.2 Example: Modeling of a Single-Link Flexible Arm

Consider a flexible robot with a single flexible link (see Fig. 14).

If the link is a steel bar with a length of $L = 2m$ and with a squared section of side $l = 0.01m = 1cm$, the first five frequencies computed with the theoretic formulae:

$$f_i = \sqrt{\frac{EI}{m}} \frac{\pi}{32L^2} (4i+1)^2, \quad i = 1, 2, \dots, 5 \quad (70)$$

where E is the modulus of elasticity of Joung ($\cong 10^{11} N/m^2$), $I = l^4/12$ is the moment of inertia of the section respect to the deflection axis and $m \cong 7876l^2 Kg/m$ is the mass per unit length of the link, are:

$$f = 6.31, 20.4, 42.7, 73.0, 111.3 \text{ Hz} . \quad (71)$$

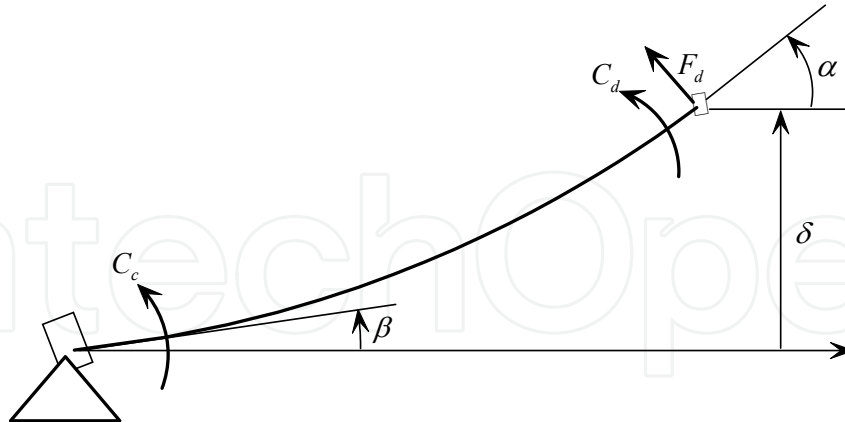


Fig. 14. Schematic representation of a single-link flexible arm.

The frequencies obtained with the proposed method, dividing the link into five and ten flexible sublinks, respectively result:

$$f = 6.31, 20.4, 41.9, 65.6, 90.8 \text{ Hz} \quad (72)$$

$$f = 6.31, 20.4, 42.6, 72.77, 110.5 \text{ Hz} . \quad (73)$$

Stressing the link with $C_c = -1 \text{ Nm}$, $C_d = 1 \text{ Nm}$, $F_d = 0 \text{ N}$, if $\beta = 0$, the theoretic values of α and δ are:

$$\alpha = \frac{L}{EI} C_d = 1.375^\circ, \quad \delta = \frac{L^2}{2EI} C_d = 2.40 \text{ cm} . \quad (74)$$

It is worth noting that these values and the ones obtained using the proposed method are coincident $\forall \nu \geq 1$ (Celentano, 2007).

Applying to the link $C_c = -2 \text{ Nm}$, $C_d = 0 \text{ Nm}$, $F_d = 1 \text{ N}$, if $\beta = 0$, the theoretic values of α and δ result:

$$\alpha = \frac{L^2}{2EI} F_d = 1.375^\circ, \quad \delta = \frac{L^3}{3EI} F_d = 3.20 \text{ cm} . \quad (75)$$

The first value (the orientation angle of payload due to the arm deflection) and the ones computed with the proposed method $\forall \nu \geq 1$ are coincident, while the second value (the motion of the payload due to the arm deflection) obtained using the proposed method has a relative error of (Celentano, 2007)

$$\varepsilon_\delta = -\frac{1}{4\nu^2} = \begin{cases} -0.01 = 1\%, & \text{if } \nu = 5 \\ -0.0025 = 0.25\%, & \text{if } \nu = 10. \end{cases} \quad (76)$$

6. Conclusion

In this chapter an innovative method for robots modeling and simulation, based on an appropriate mathematical formulation of the relative equations of motion and on a new

integration scheme, has been illustrated. The proposed approach does require the calculation of the inertia matrix and of the gradient of the kinetic energy only. It provides a new analytical-numerical methodology, that has been shown to be simpler and numerically more efficient than the classical approaches, requires no a priori specialized knowledge of the dynamics of mechanical systems and is formulated in order to allow students, researchers and professionals to easily employ it for the analysis of manipulators with the complex-shaped links commonly used in industry.

In the case of planar robots with revolute joints, theorems have been stated and proved that offer a particularly simple and efficient method of computation for both the inertia matrix and the gradient of the kinetic energy. Then a comparison has been made in terms of efficiency between the proposed method and the Articulated-Body one.

Moreover, for spatial robots with generic shape links and connected, for the sake of brevity, with spherical joints, several theorems have been formulated and demonstrated in a simple manner and some algorithms that allow efficiently computing, analytically the inertia matrix, analytically or numerically the gradient of the kinetic and of the gravitational energy have been provided. Furthermore, also in this case a comparison of the proposed method in terms of efficiency with the Articulated-Body one has been reported.

Finally, a methodology for flexible robots modeling, that allow obtaining, quite simply, accurate and efficient, from a computational point of view, finite-dimensional models, has been provided. This method is illustrated with a very significant example.

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The book *New Approaches in Automation and Robotics* offers in 22 chapters a collection of recent developments in automation, robotics as well as control theory. It is dedicated to researchers in science and industry, students, and practicing engineers, who wish to update and enhance their knowledge on modern methods and innovative applications. The authors and editor of this book wish to motivate people, especially under-graduate students, to get involved with the interesting field of robotics and mechatronics. We hope that the ideas and concepts presented in this book are useful for your own work and could contribute to problem solving in similar applications as well. It is clear, however, that the wide area of automation and robotics can only be highlighted at several spots but not completely covered by a single book.

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